

Observational Heliophysics I: Particles, Plasma, and Fields

Jasper S Halekas, University of Iowa Department of Physics and Astronomy

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Introduction

Why Not?



Some Random Numbers

- Temperature of solar corona = 1-2 million degrees K
- Solar wind speed = 1 million mph
- Solar wind ram pressure at 1 AU = 1 nPa
- Magnetospheric plasma density $\sim 0.01\text{-}1000\text{ cm}^{-3}$
 - Compared to Earth surface atmospheric density of $2.5 \times 10^{19}\text{ cm}^{-3}$
- Typical solar wind motional electric field at 1 AU = 1 mV/m
- Typical solar wind parallel electric field = $<1\text{ nV/m}$
- Solar wind magnetic field at 1 AU = 10 nT

Space Plasma Physics: How Hard Could it Be?

Maxwell's Equations for the Fields

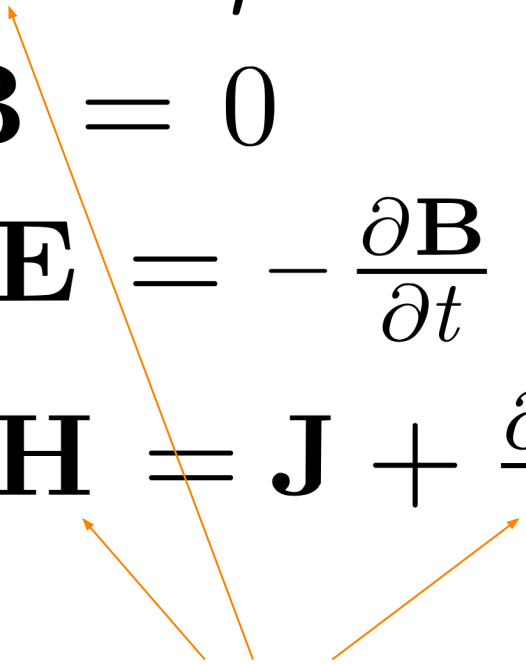
$$\nabla \cdot \mathbf{D} = \rho$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}$$

Here be dragons



Charge and Current Density

$$\rho = \sum_s q_s n_s = e(Zn_i - n_e)$$

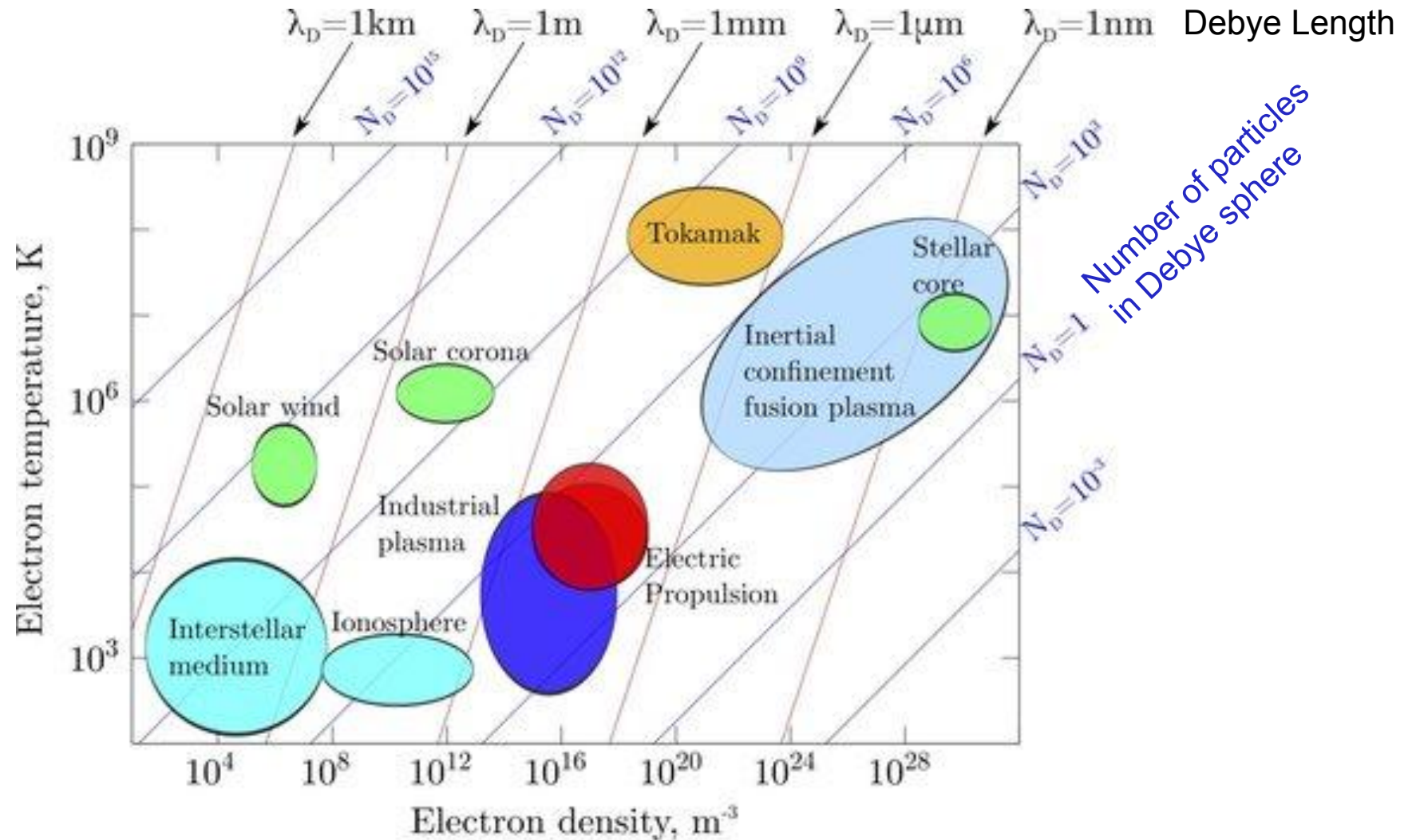
$$\mathbf{J} = \sum_s q_s n_s \mathbf{v}_s = e(Zn_i \mathbf{v}_i - n_e \mathbf{v}_e)$$

Equation of Motion (Lorentz Force)

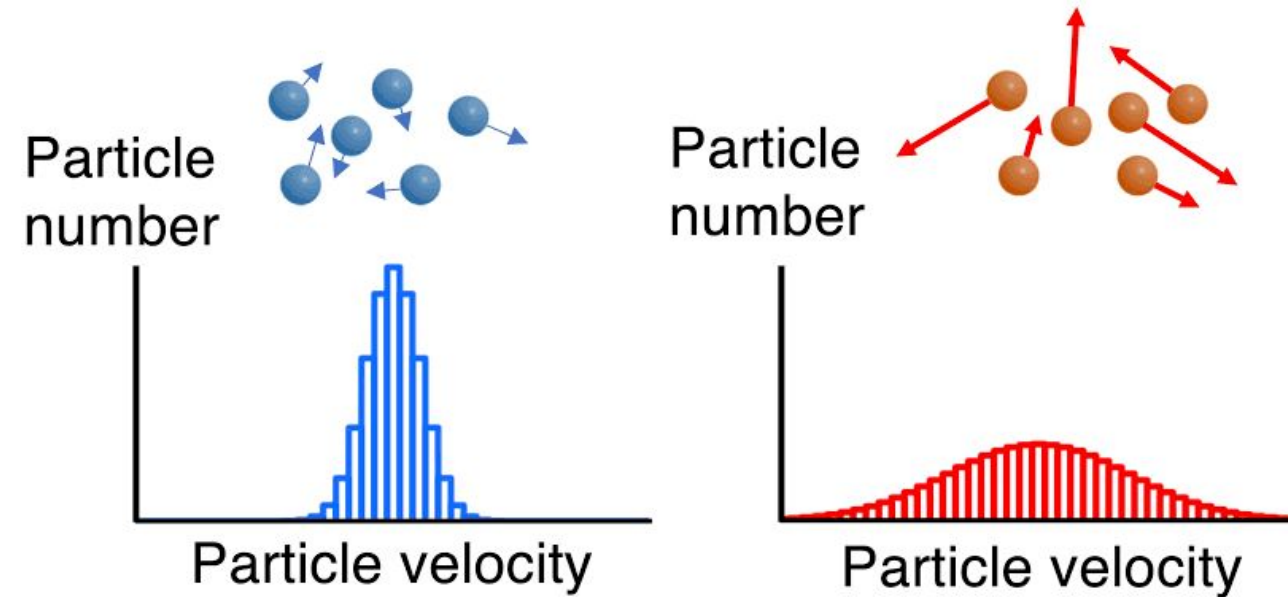
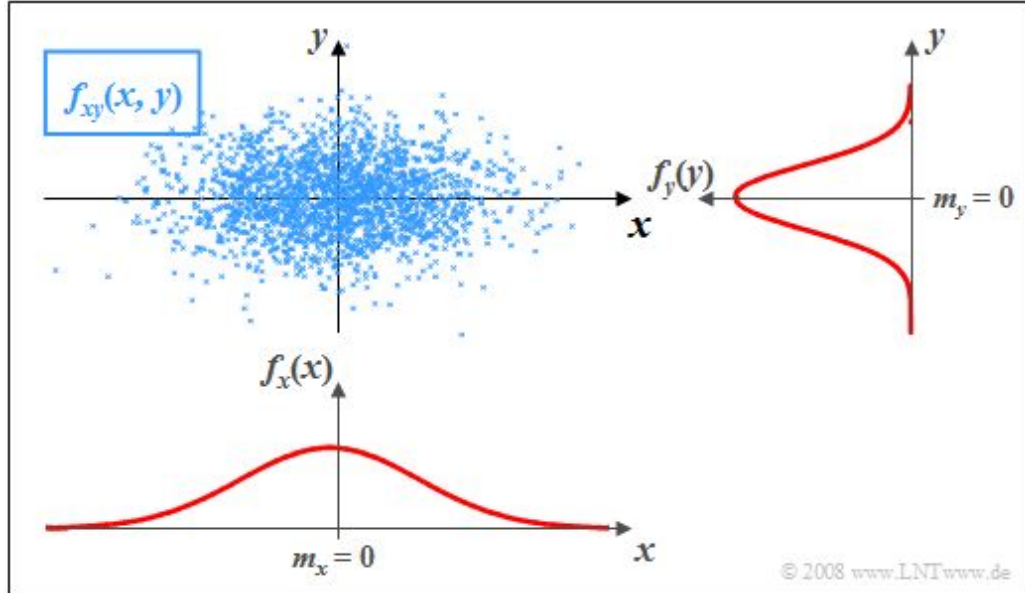
$$\mathbf{F} = m \frac{d\mathbf{v}}{dt} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B})$$

+ sprinkle in collisions, other forces as needed

So Many Particles...



The Distribution Function



$$f(\mathbf{r}, \mathbf{v}, t) = \frac{dN}{d^3\mathbf{r} d^3\mathbf{v}}$$

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla_{\mathbf{x}} f + \frac{\mathbf{F}}{m} \cdot \nabla_{\mathbf{v}} f = \left(\frac{\partial f}{\partial t} \right)_{\text{collisions}}$$

Boltzmann Equation

Moments of the Distribution Function

(a) Particle density,

$$n = \int f(\vec{v}) d\vec{v}$$

(b) Particle flux density,

$$\vec{\Gamma} = \int \vec{v} f(\vec{v}) d\vec{v}$$

(c) Flow velocity,

$$\vec{u} = \vec{\Gamma}/n$$

(d) Current density,

$$\vec{j} = Ze\vec{\Gamma}$$

(e) Stress tensor (diagonal components only),

$$P_{xx} = m \int v_x^2 f(\vec{v}) d\vec{v}$$

$$P_{yy} = m \int v_y^2 f(\vec{v}) d\vec{v}$$

$$P_{zz} = m \int v_z^2 f(\vec{v}) d\vec{v}$$

(f) Pressure tensor (diagonal components only),

$$p_{xx} = m \int (v_x - u_x)^2 f(\vec{v}) d\vec{v}$$

$$p_{yy} = m \int (v_y - u_y)^2 f(\vec{v}) d\vec{v}$$

$$p_{zz} = m \int (v_z - u_z)^2 f(\vec{v}) d\vec{v}$$

(g) Scalar pressure (from the trace of the pressure tensor),

$$p = \frac{1}{3}(p_{xx} + p_{yy} + p_{zz})$$

(h) Energy flux density

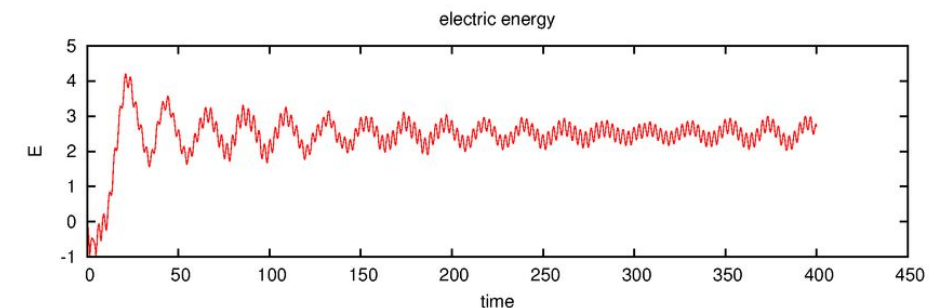
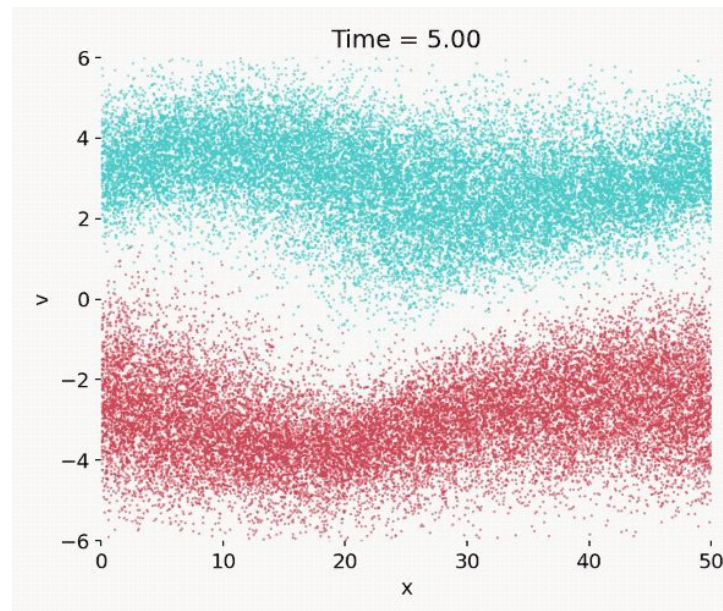
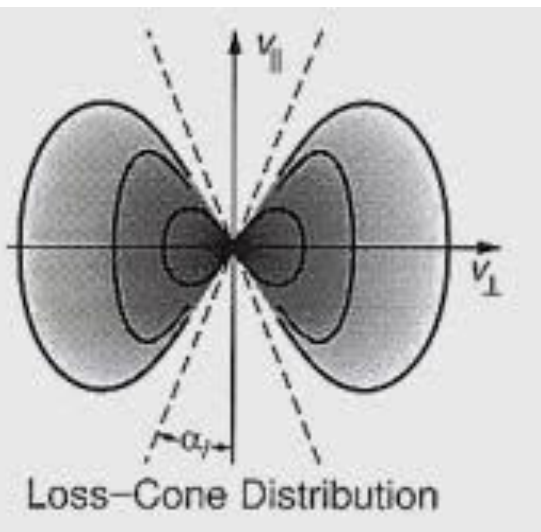
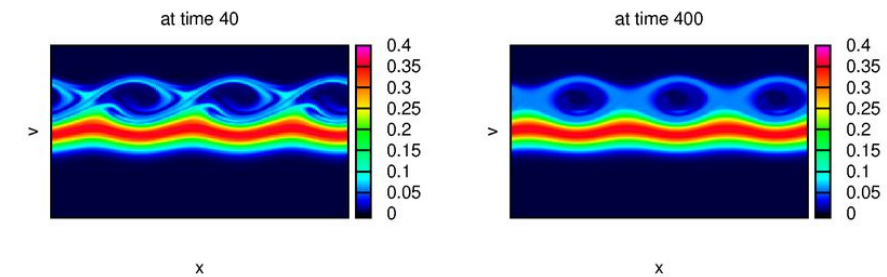
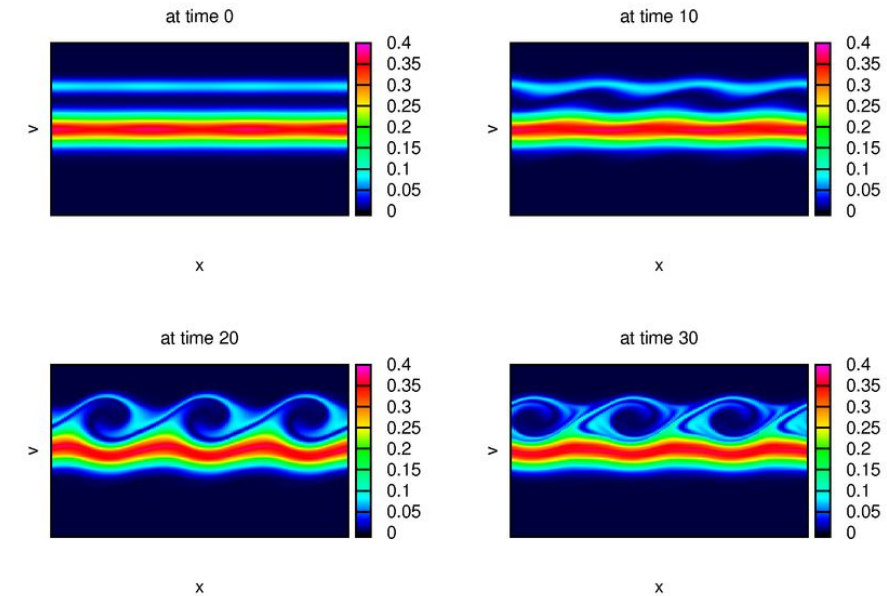
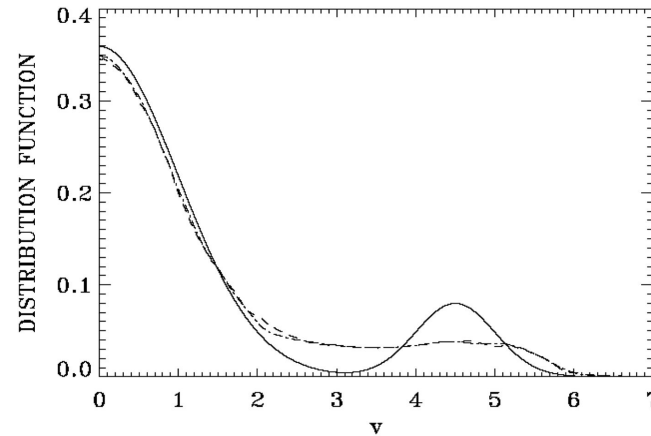
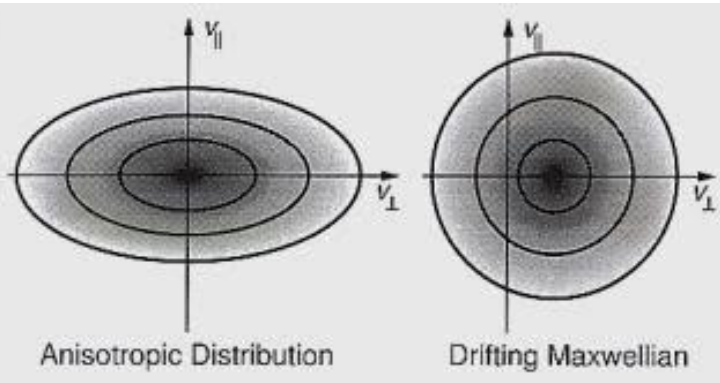
$$\vec{Q} = \frac{1}{2}m \int v^2 \vec{v} f(\vec{v}) d\vec{v}$$

(i) Heat flow

$$\vec{q} = \frac{1}{2}m \int w^2 (\vec{v} - \vec{u}) f(\vec{v}) d\vec{v}$$

$$w^2 = (v_x - u_x)^2 + (v_y - u_y)^2 + (v_z - u_z)^2$$

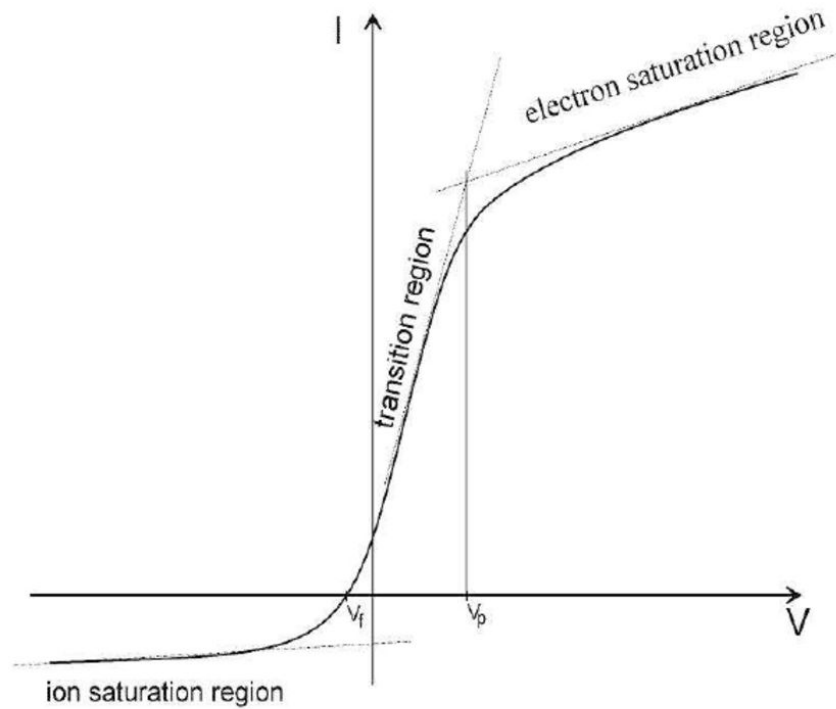
Beyond Moments



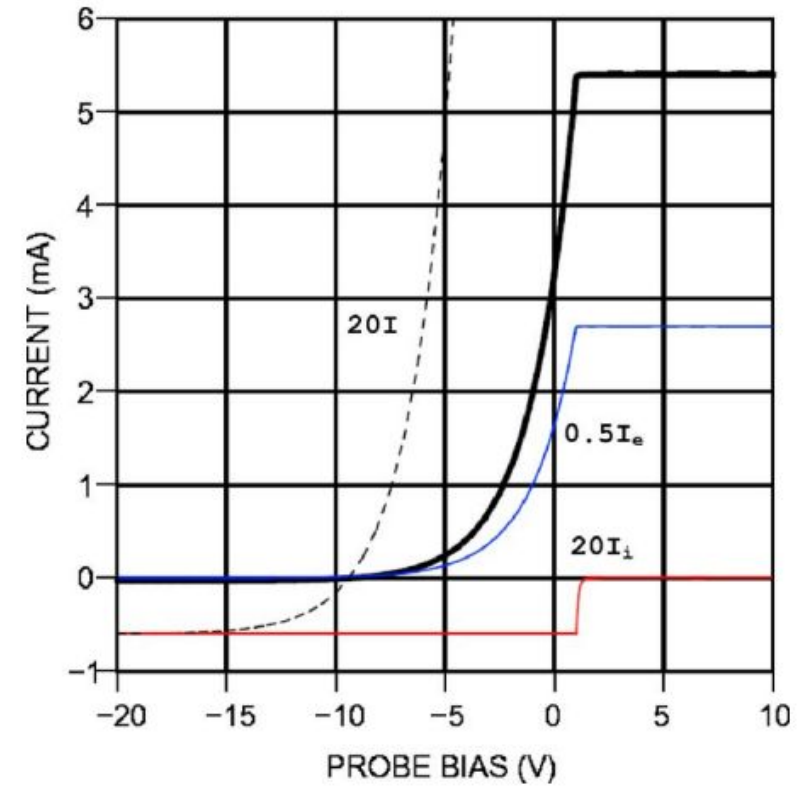
Plasma and Particles

In Situ Techniques: Measuring Current – Langmuir Probe

Thick Sheath

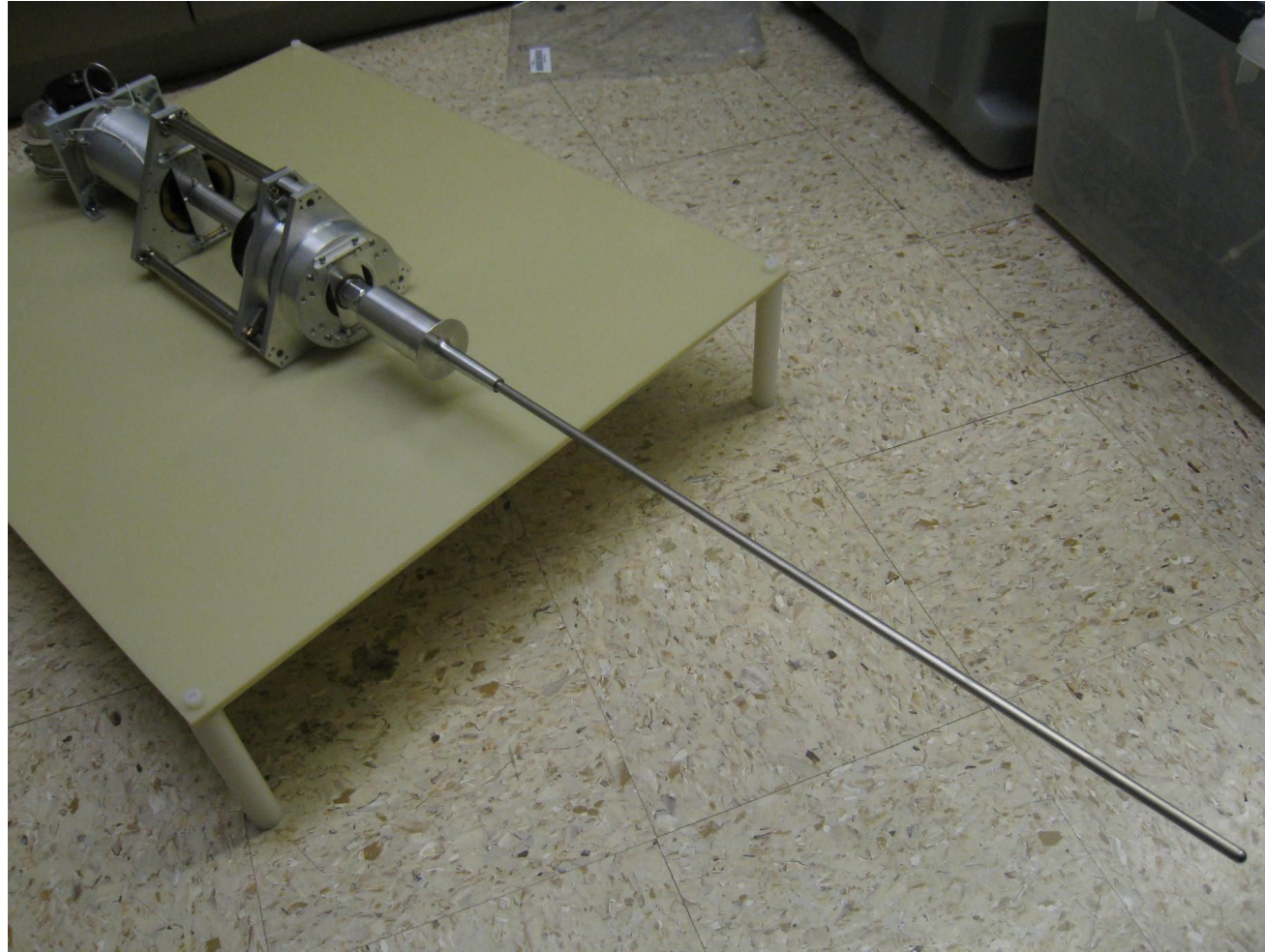


Thin Sheath

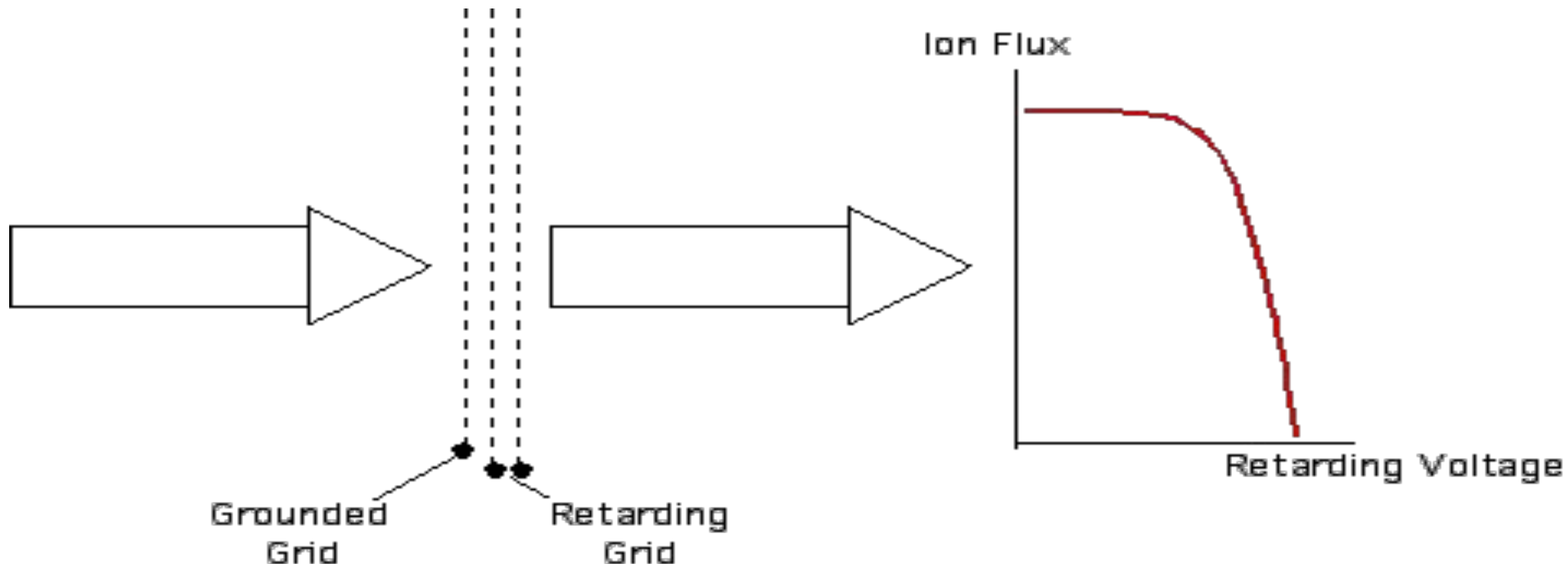


V_f = Probe Floating Potential, V_p = Plasma Potential

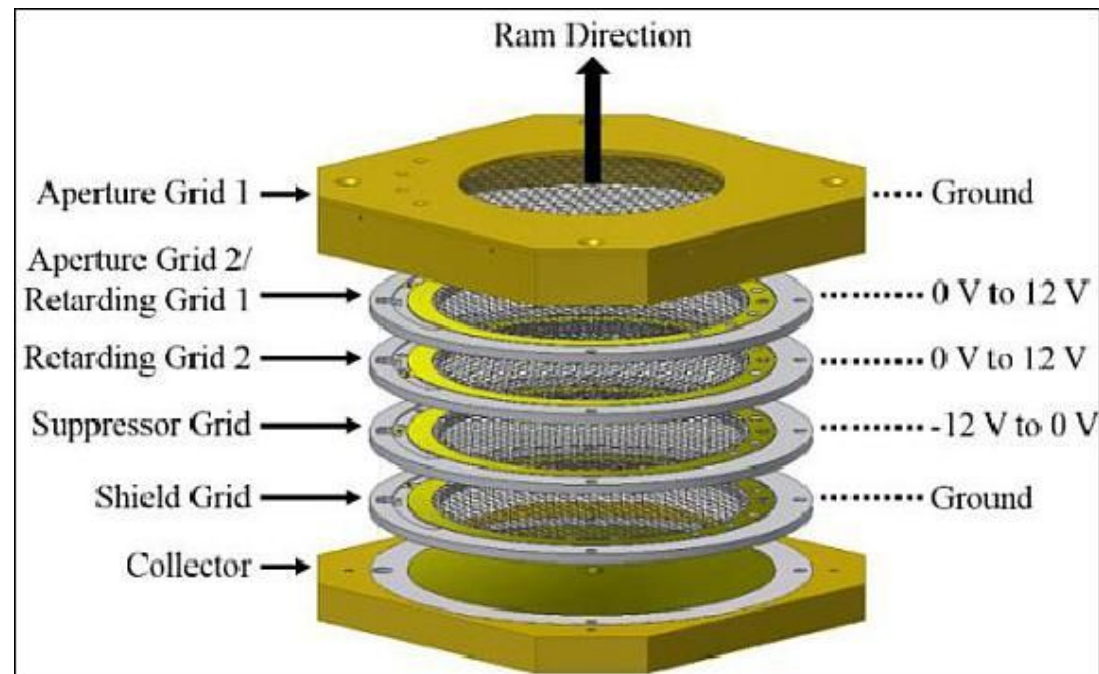
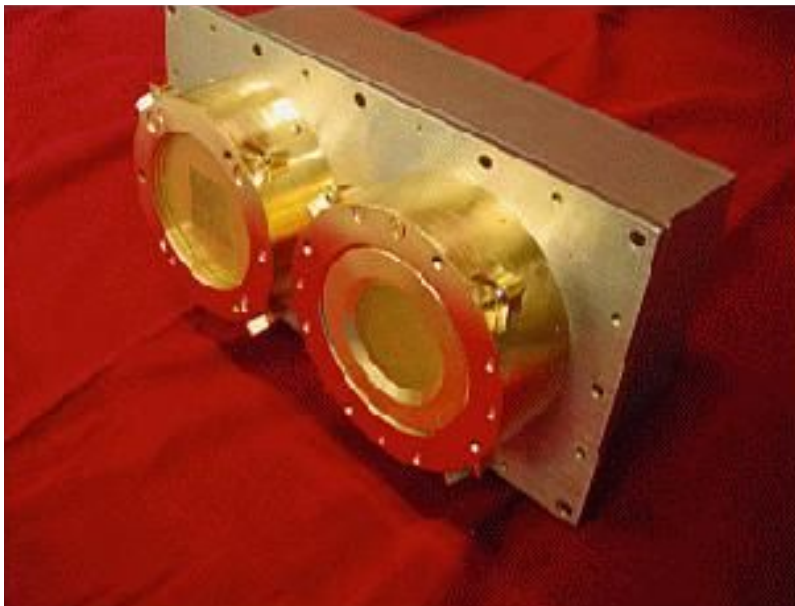
In Situ Techniques: Measuring Current – Langmuir Probe



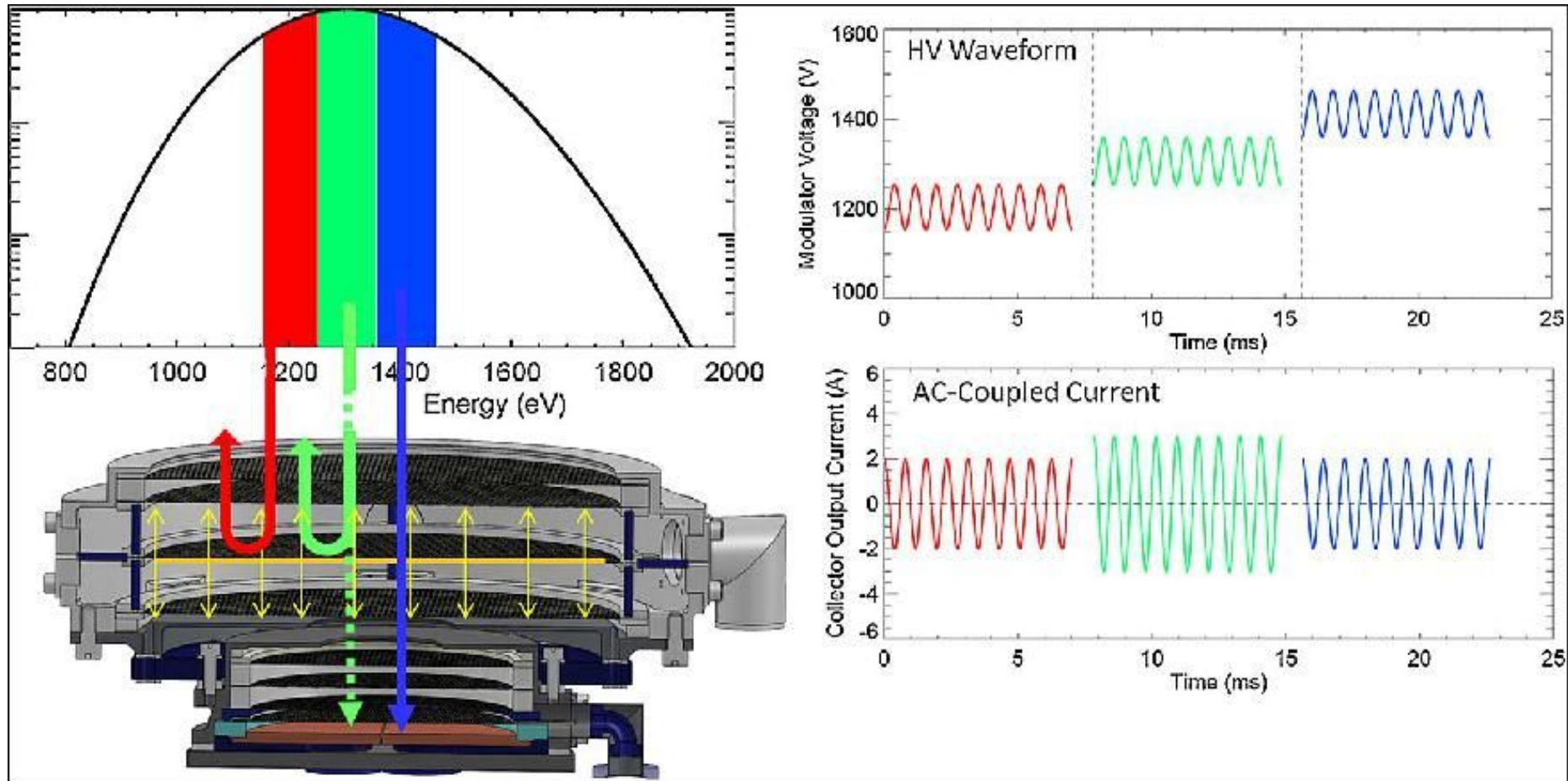
In Situ Techniques: Measuring Current - Retarding Potential Analyzer



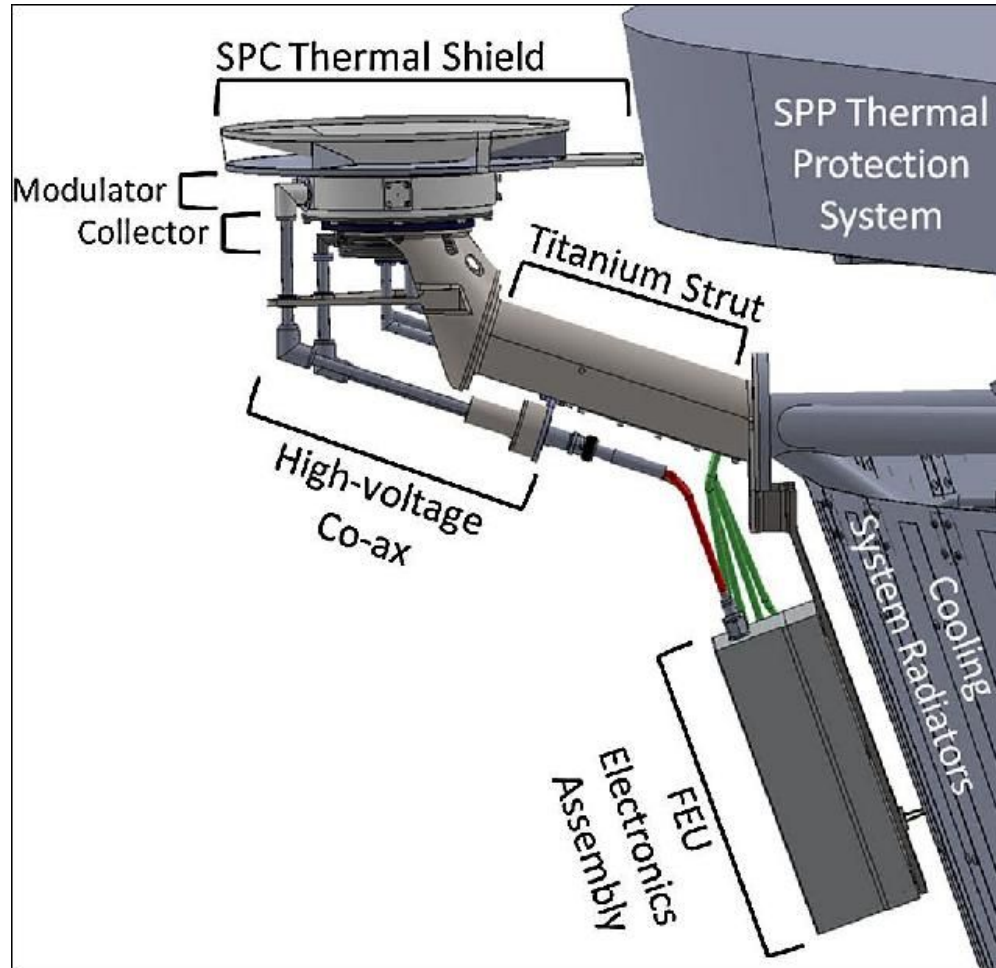
In Situ Techniques: Measuring Current – Retarding Potential Analyzer



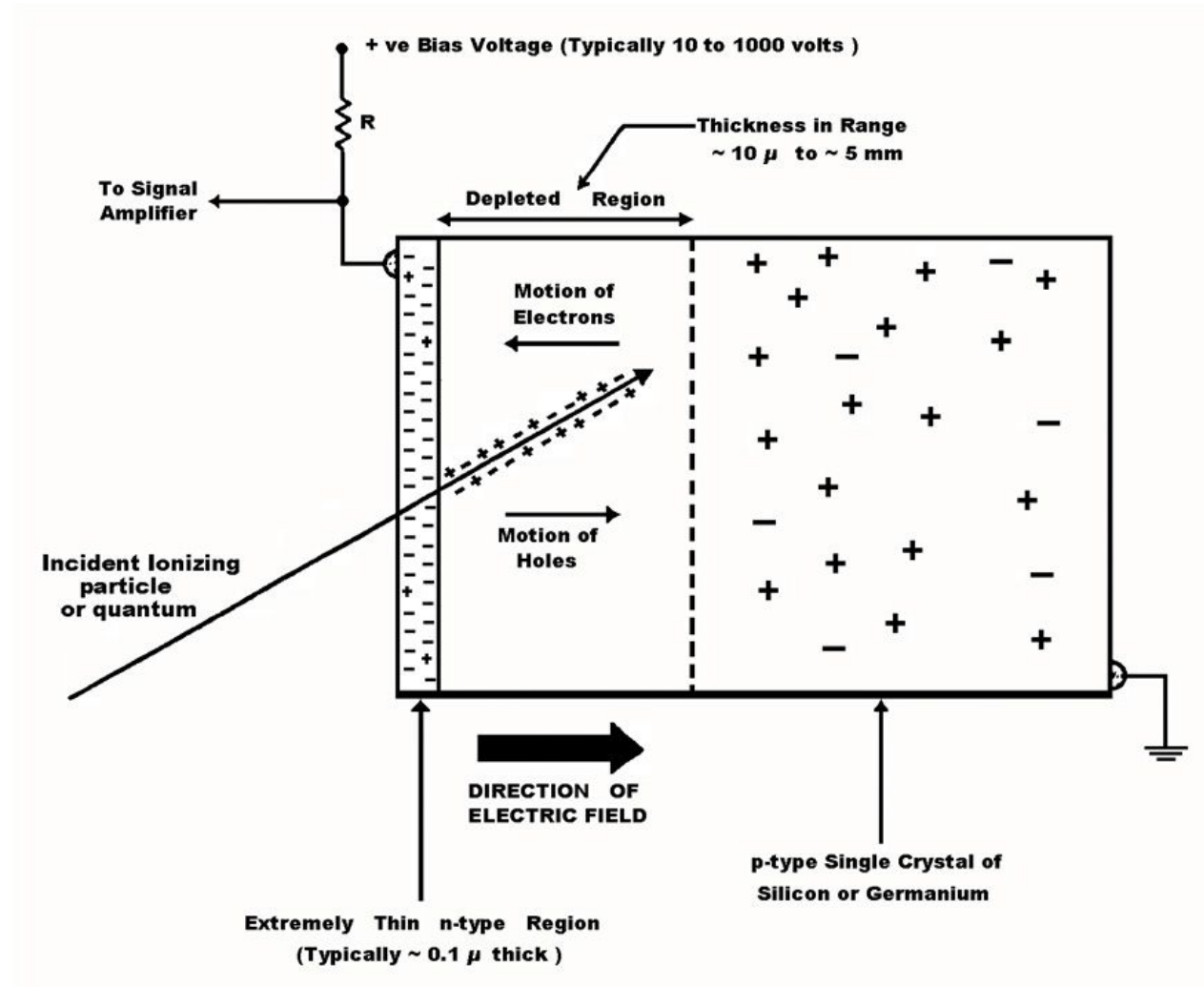
In Situ Techniques: Measuring Current – Faraday Cup



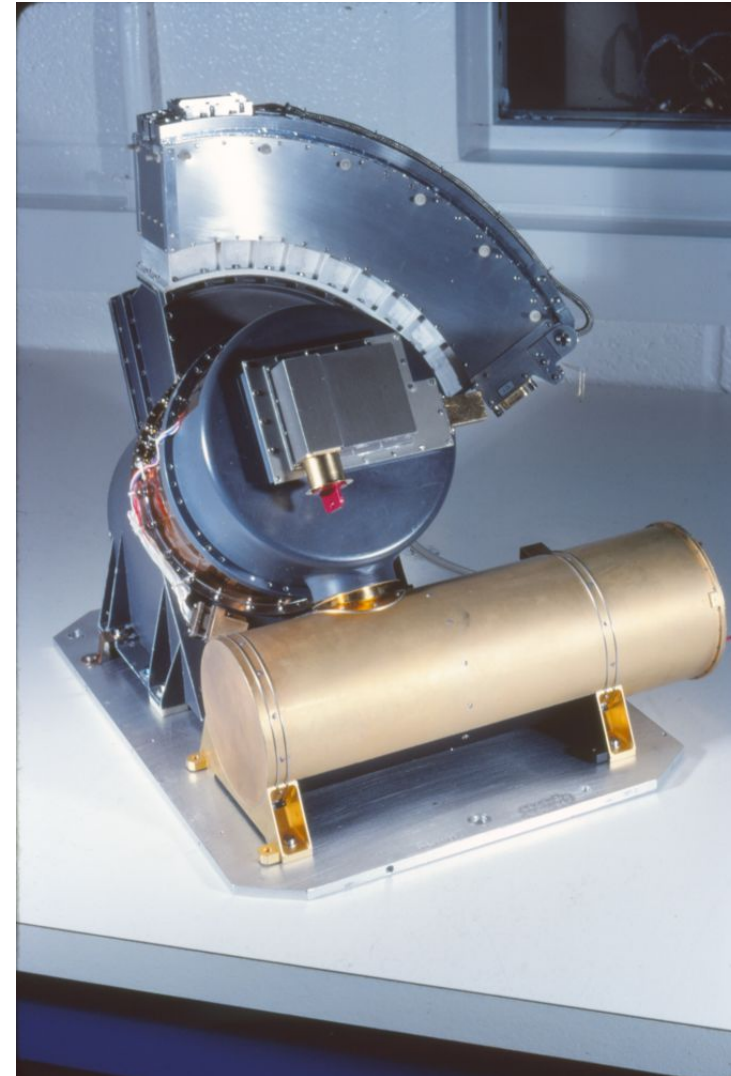
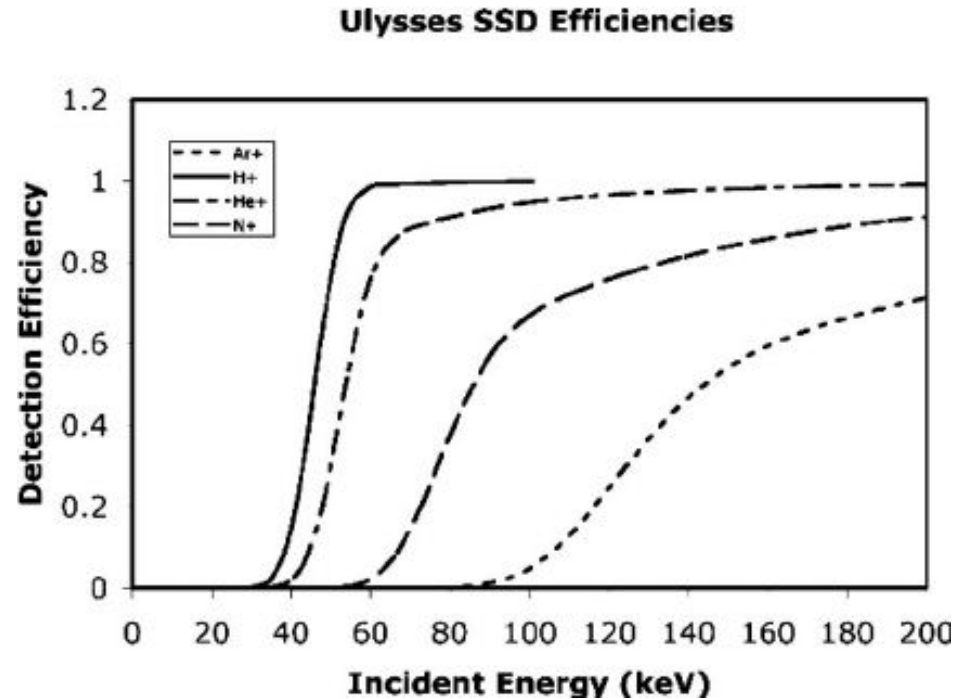
In Situ Techniques: Measuring Current – Faraday Cup



In-Situ Techniques: Counting Particles – Solid-State Detectors



In-Situ Techniques: Counting Particles – Solid-State E Deposition



Bethe-Bloch Formula

$$-\frac{dE}{dx} = 2\pi N_a r_e^2 m_e c^2 \rho \frac{Z}{A} \frac{z^2}{\beta^2} \left[\ln \left(\frac{2m_e \gamma^2 v^2 W_{\max}}{I^2} \right) - 2\beta^2 - \delta - 2 \frac{C}{Z} \right],$$

with

$$2\pi N_a r_e^2 m_e c^2 = 0.1535 \text{ MeV cm}^2/\text{g}$$

r_e : classical electron
radius = $2.817 \times 10^{-13} \text{ cm}$

m_e : electron mass

N_a : Avogadro's
number = $6.022 \times 10^{23} \text{ mol}^{-1}$

I : mean excitation potential

Z : atomic number of absorbing
material

A : atomic weight of absorbing material

ρ : density of absorbing material

z : charge of incident particle in
units of e

$\beta = v/c$ of the incident particle

$\gamma = 1/\sqrt{1-\beta^2}$

δ : density correction

C : shell correction

$W_{\max}$: maximum energy transfer in a
single collision.

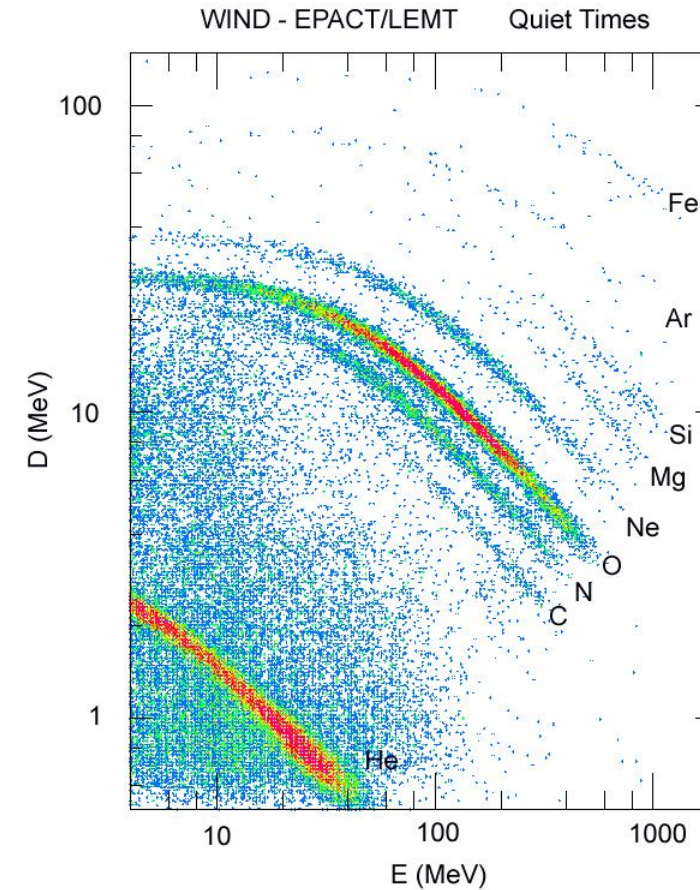
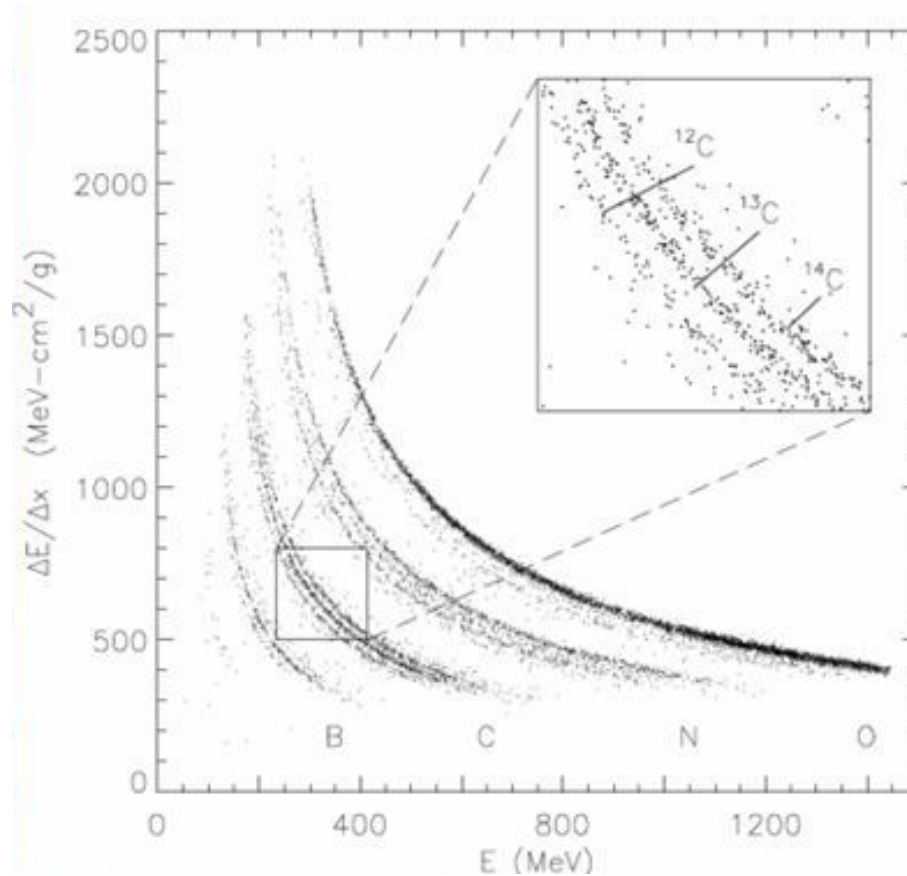
The maximum energy transfer is that produced by a head-on or *knock-on* collision. For an incident particle of mass M , kinematics gives

$$W_{\max} = \frac{2m_e c^2 \eta^2}{1 + 2s\sqrt{1 + \eta^2 + s^2}},$$

$$\begin{aligned} \frac{I}{Z} &= 12 + \frac{7}{Z} \text{ eV} & Z < 13 \\ \frac{I}{Z} &= 9.76 + 58.8 Z^{-1.19} \text{ eV} & Z \geq 13 \end{aligned}$$

where $s = m_e/M$ and $\eta = \beta\gamma$. Moreover, if $M \gg m_e$, then $W_{\max} \simeq 2m_e c^2 \eta^2$

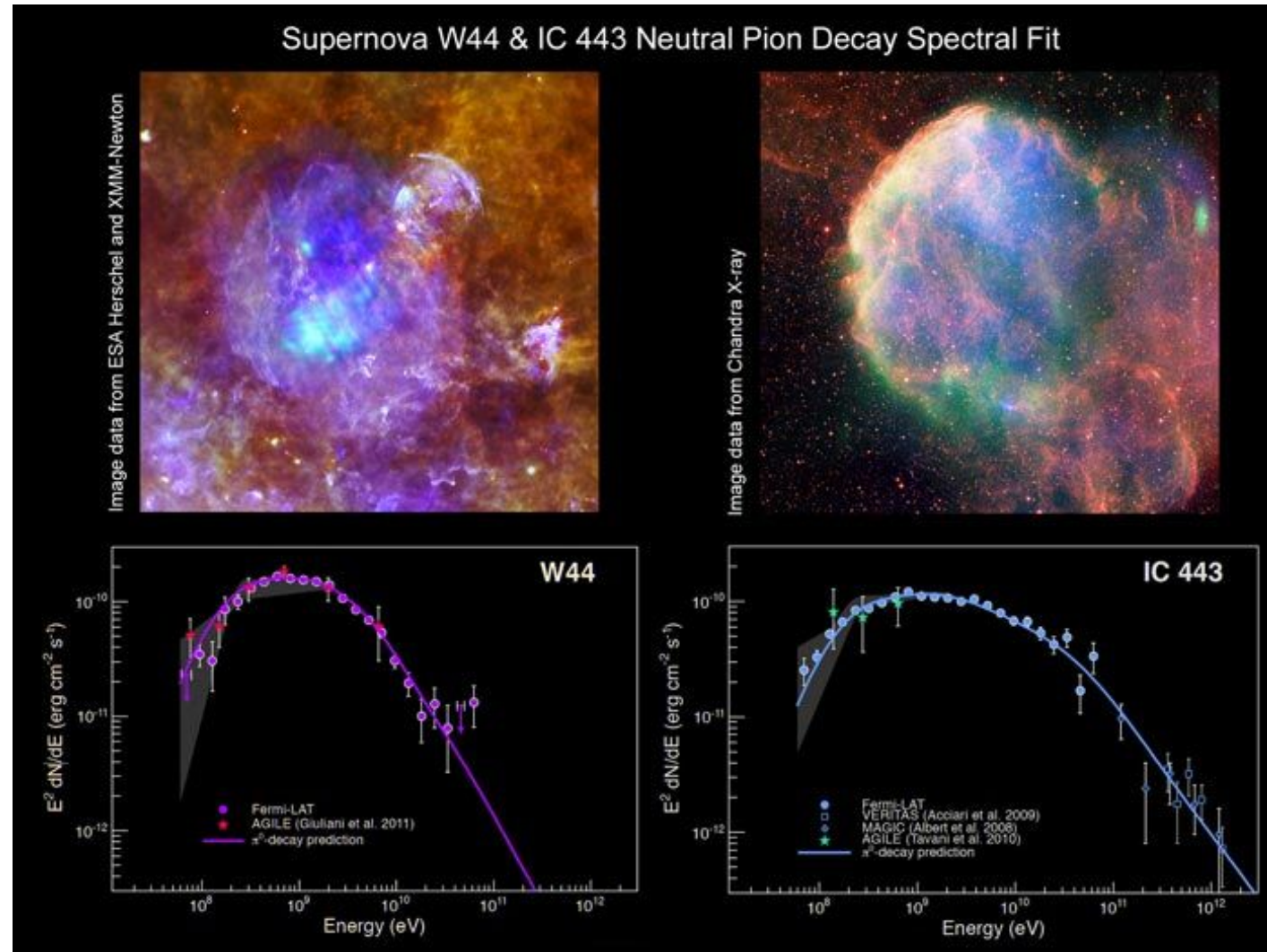
In Situ Techniques: Counting Particles – Solid State dE/dx



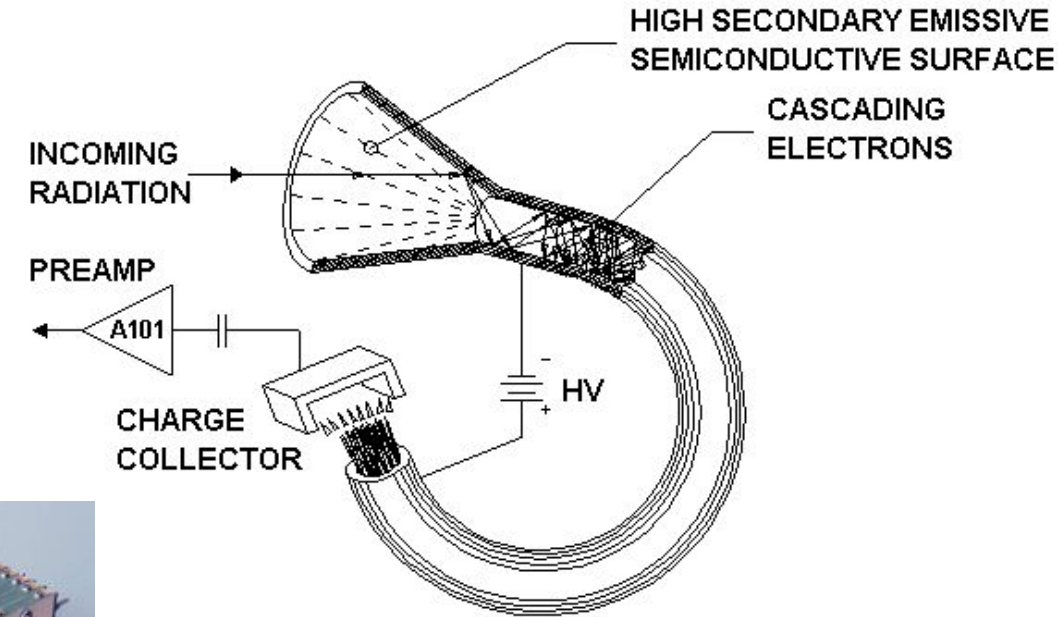
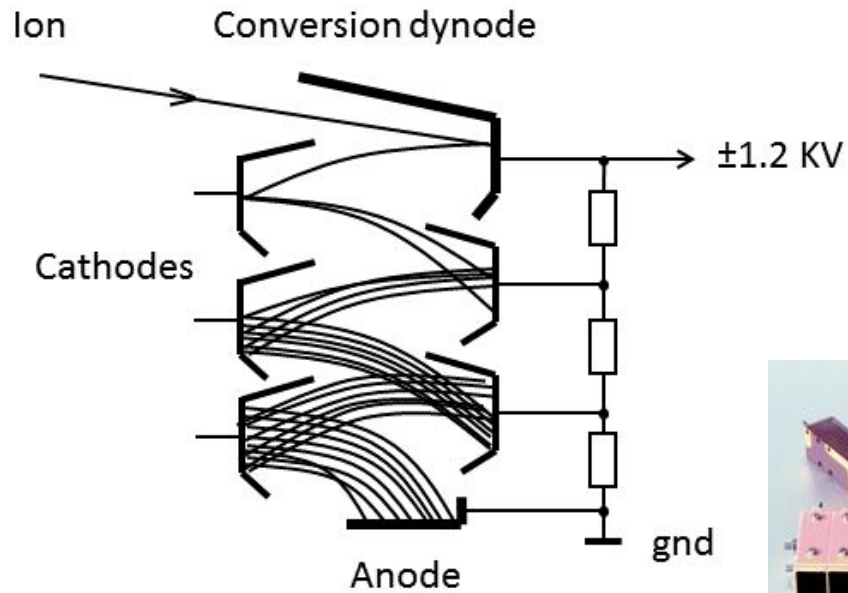
In Situ Techniques: Counting Particles – Solid State dE/dx



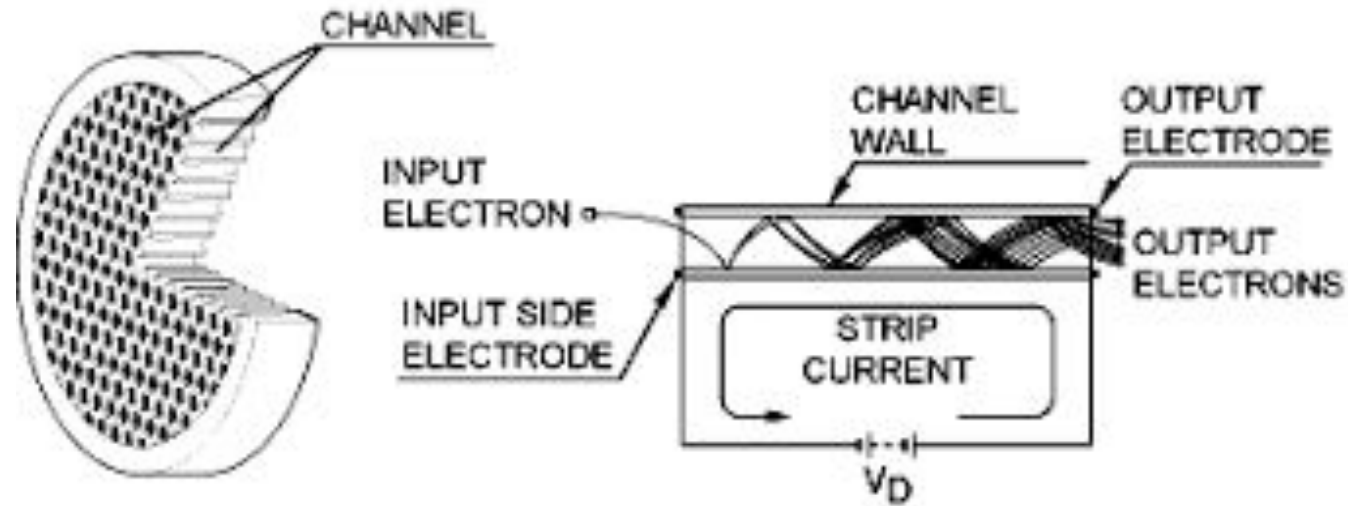
In Situ → Remote Techniques: Cosmic Rays



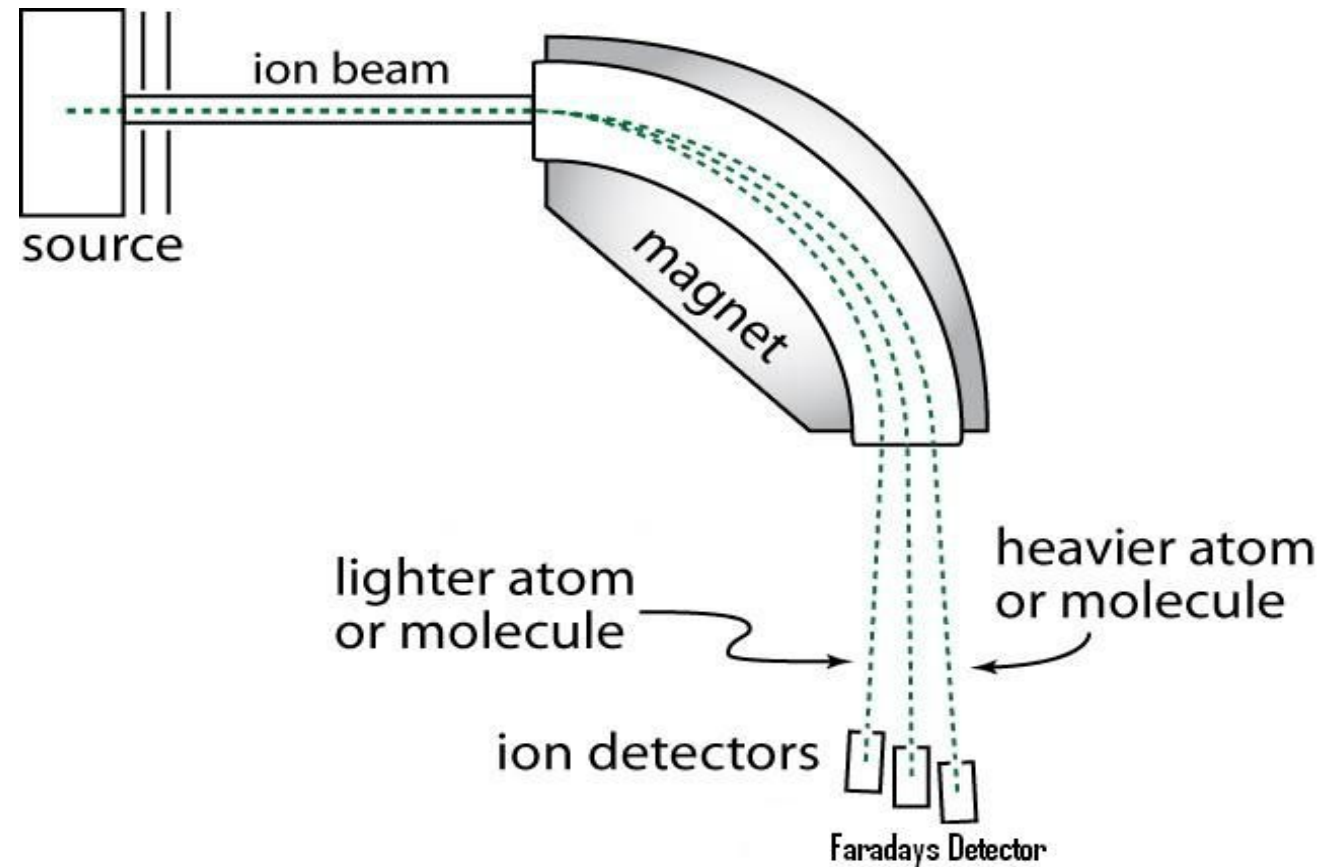
In Situ Techniques: Counting Particles – Dynodes & CEMs



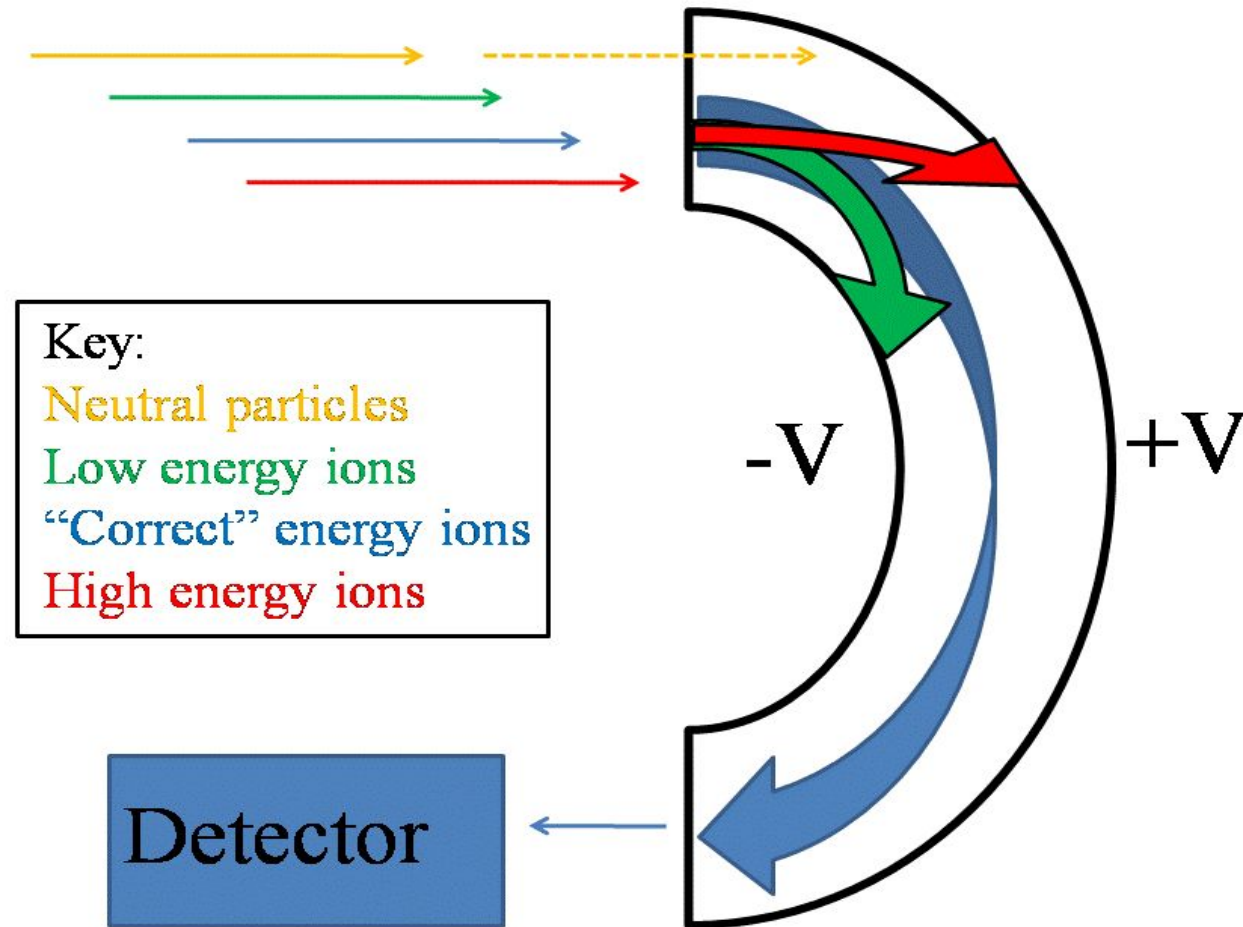
In Situ Techniques: Counting Particles – Microchannel Plates



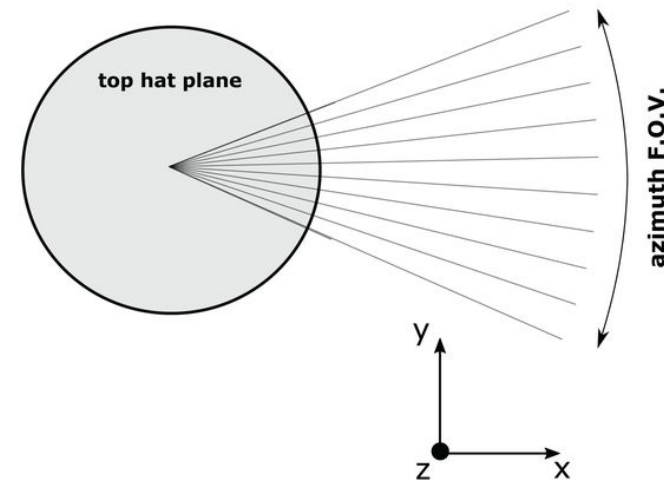
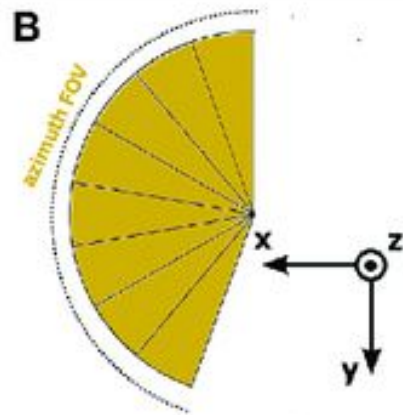
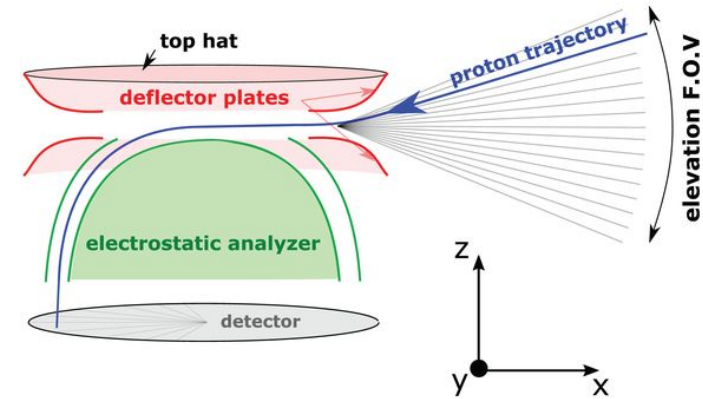
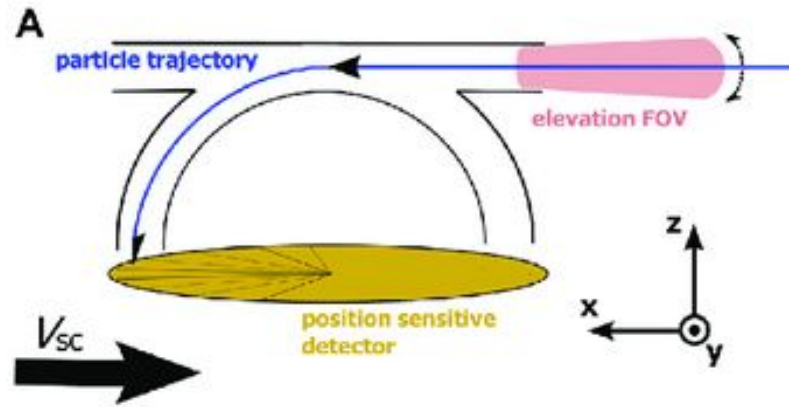
In Situ Techniques: Magnetic Optics



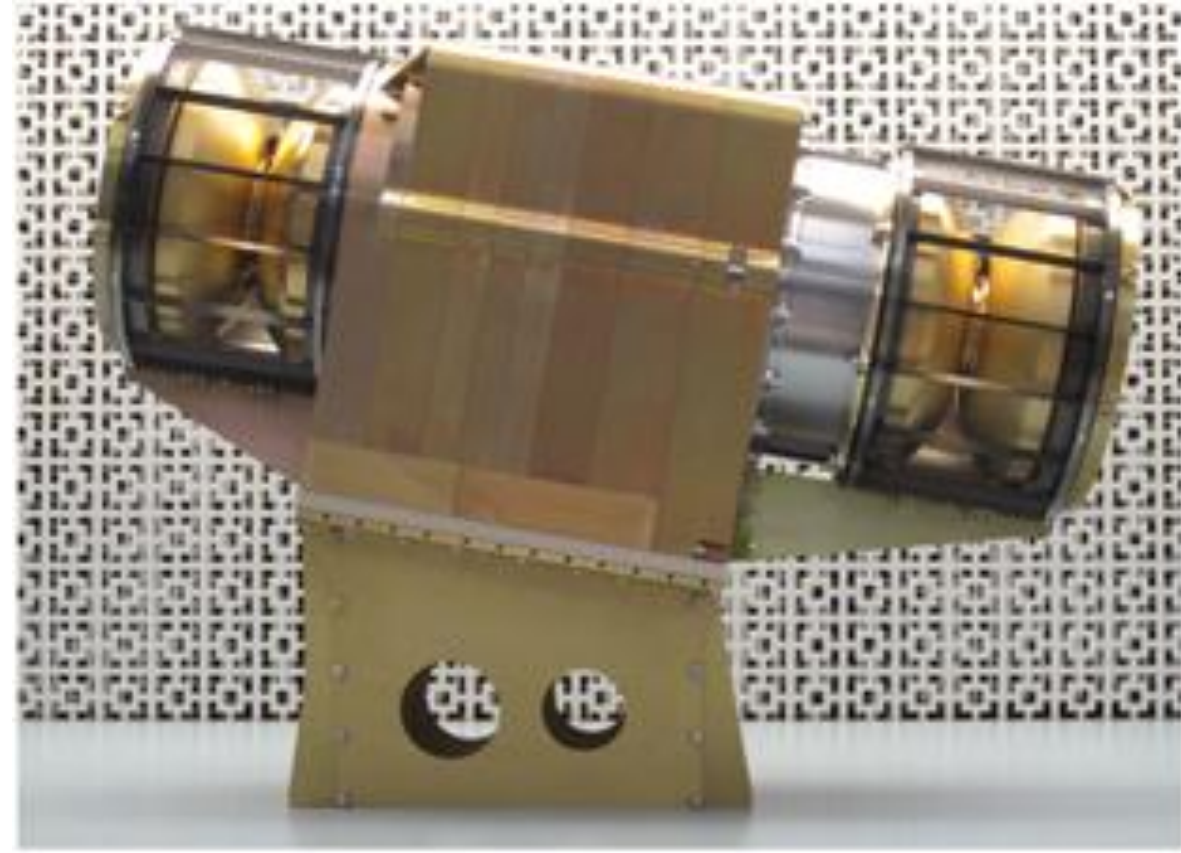
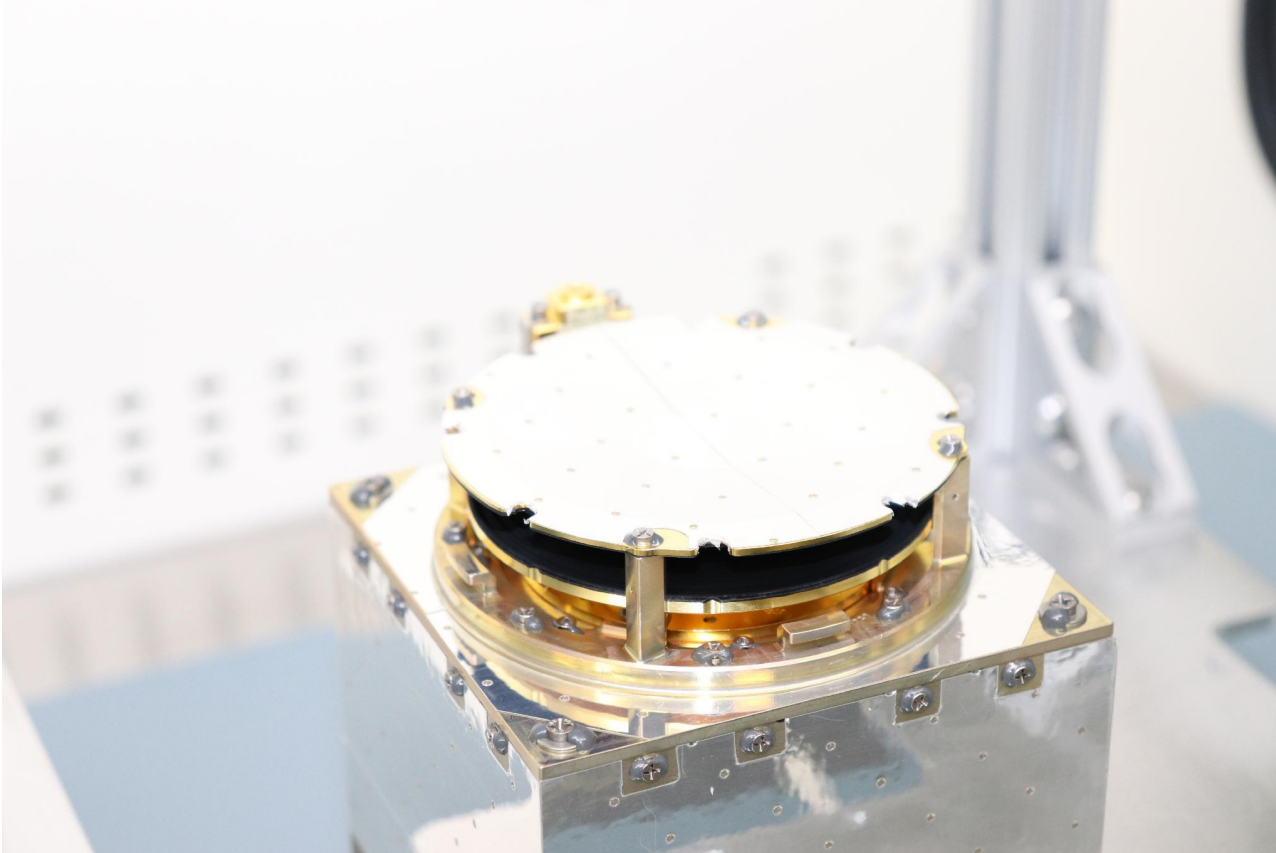
In Situ Techniques: Electrostatic Optics



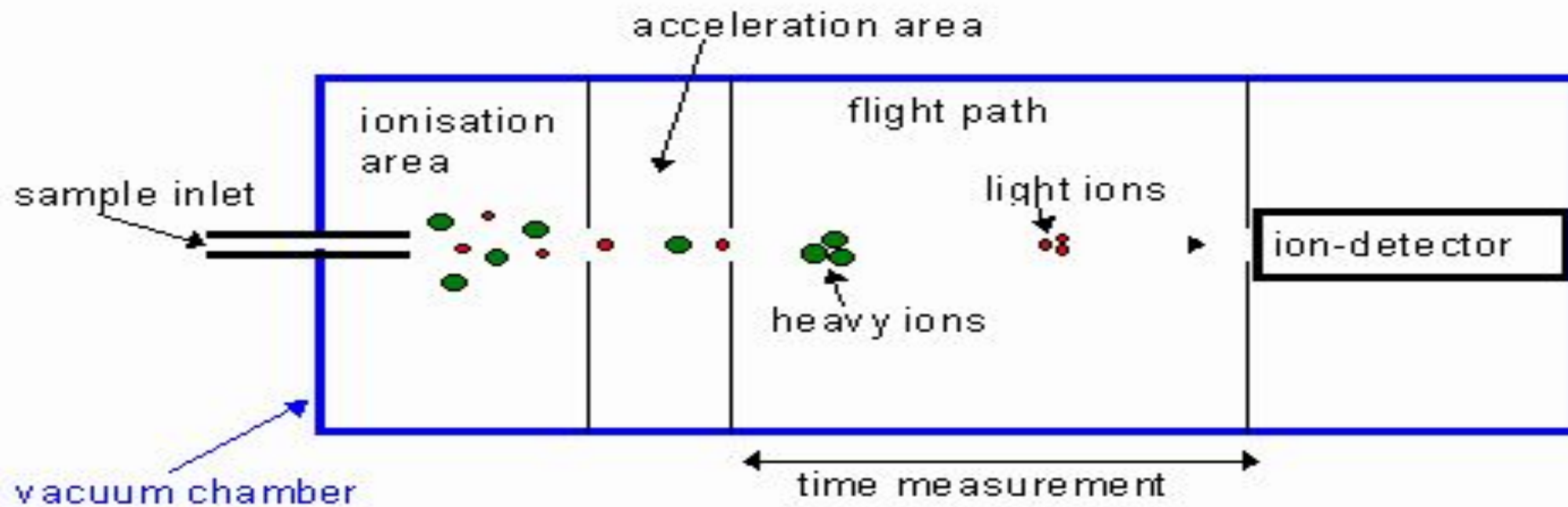
In Situ Techniques: Top-Hat Electrostatic Optics



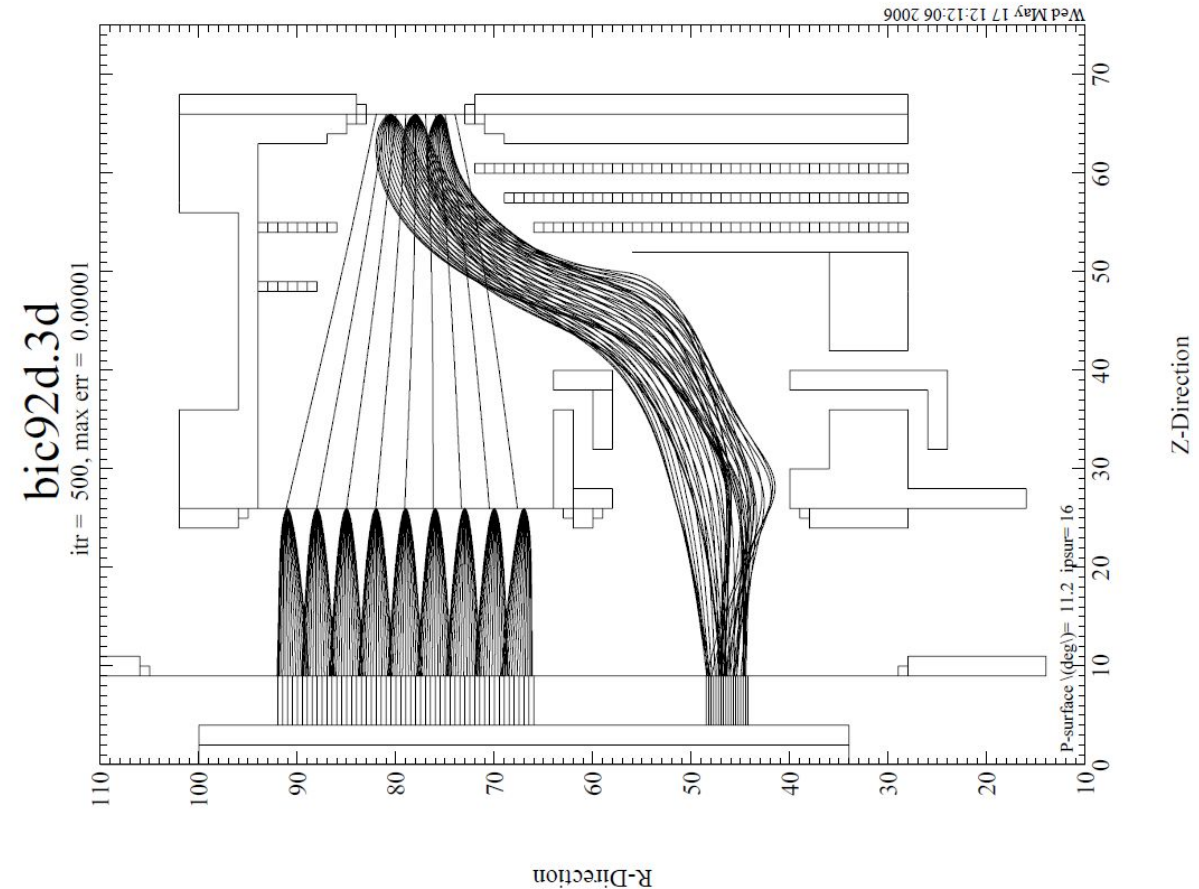
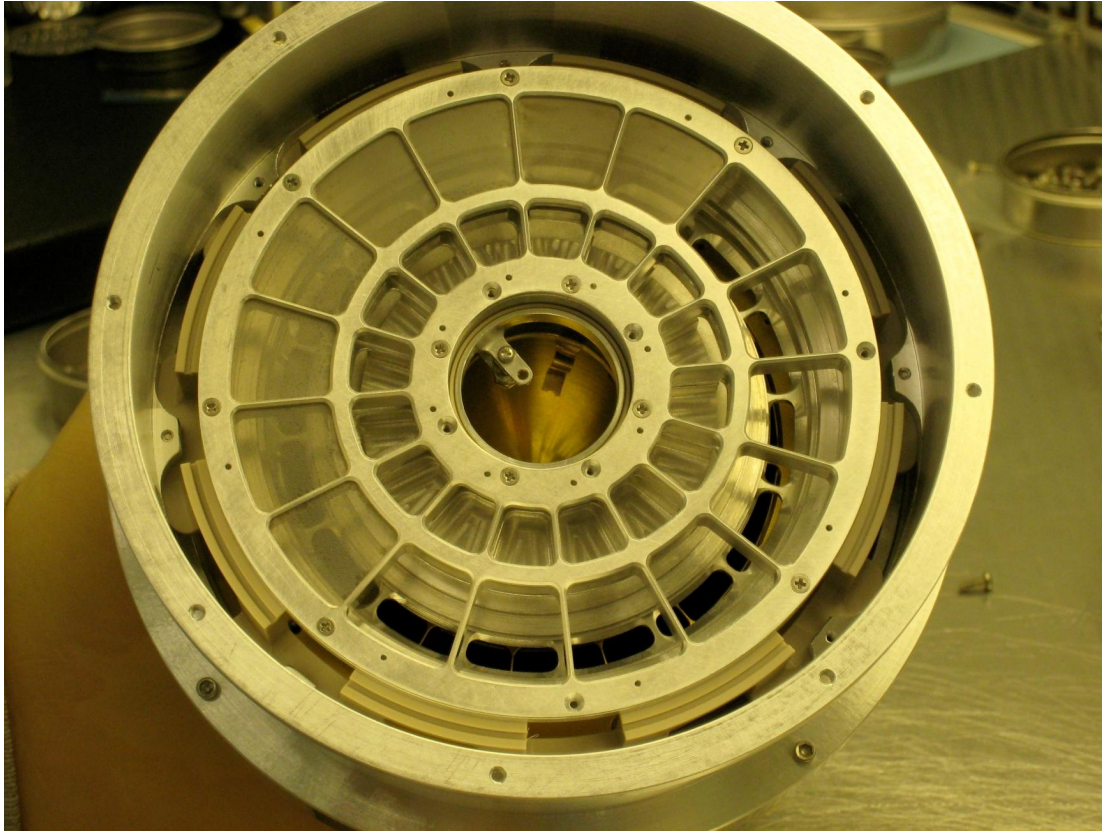
In Situ Techniques: Top-Hat Electrostatic Optics



In Situ Techniques: Counting Particles - Time-Of-Flight

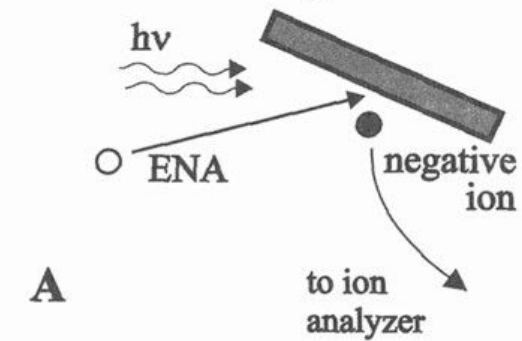


In Situ Techniques: Carbon Foils

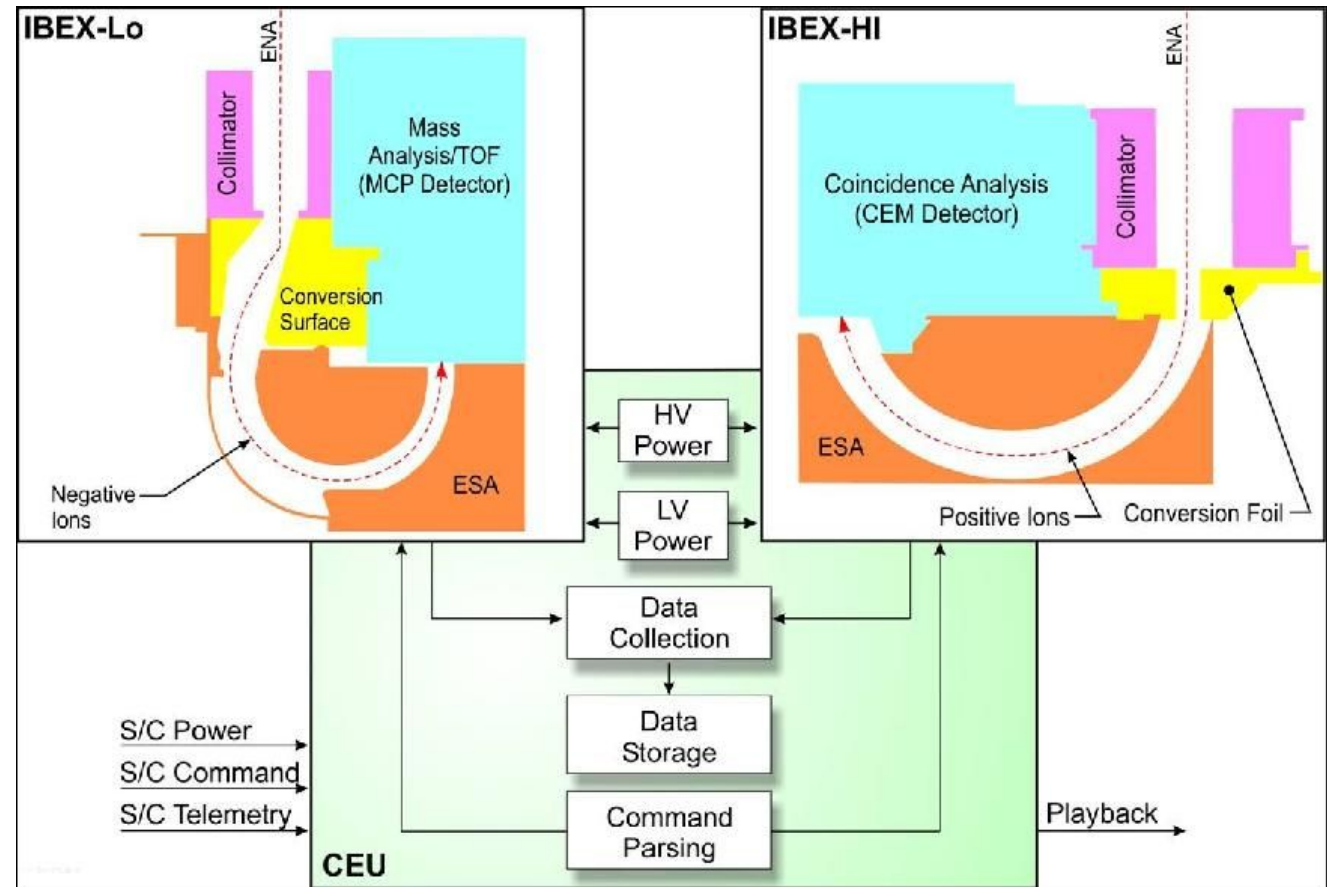
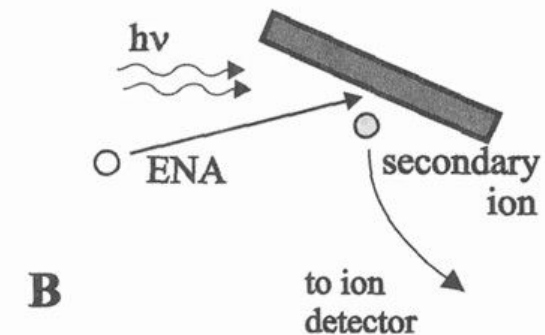


In Situ Techniques: Conversion Surfaces

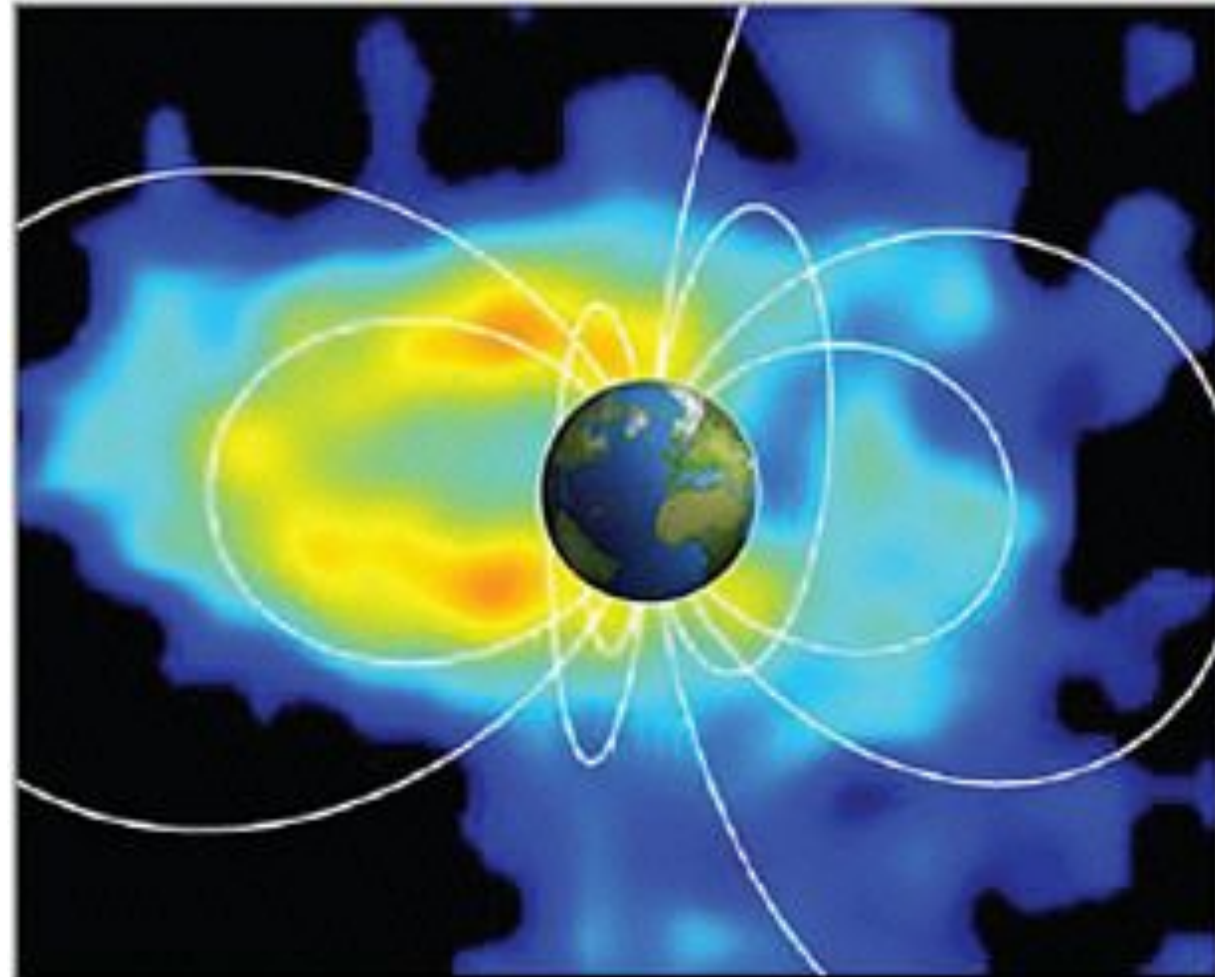
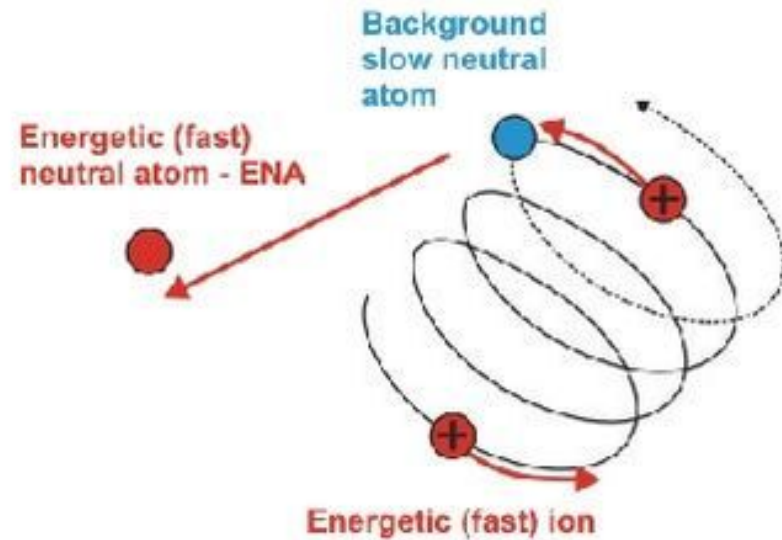
Conversion to negative ions



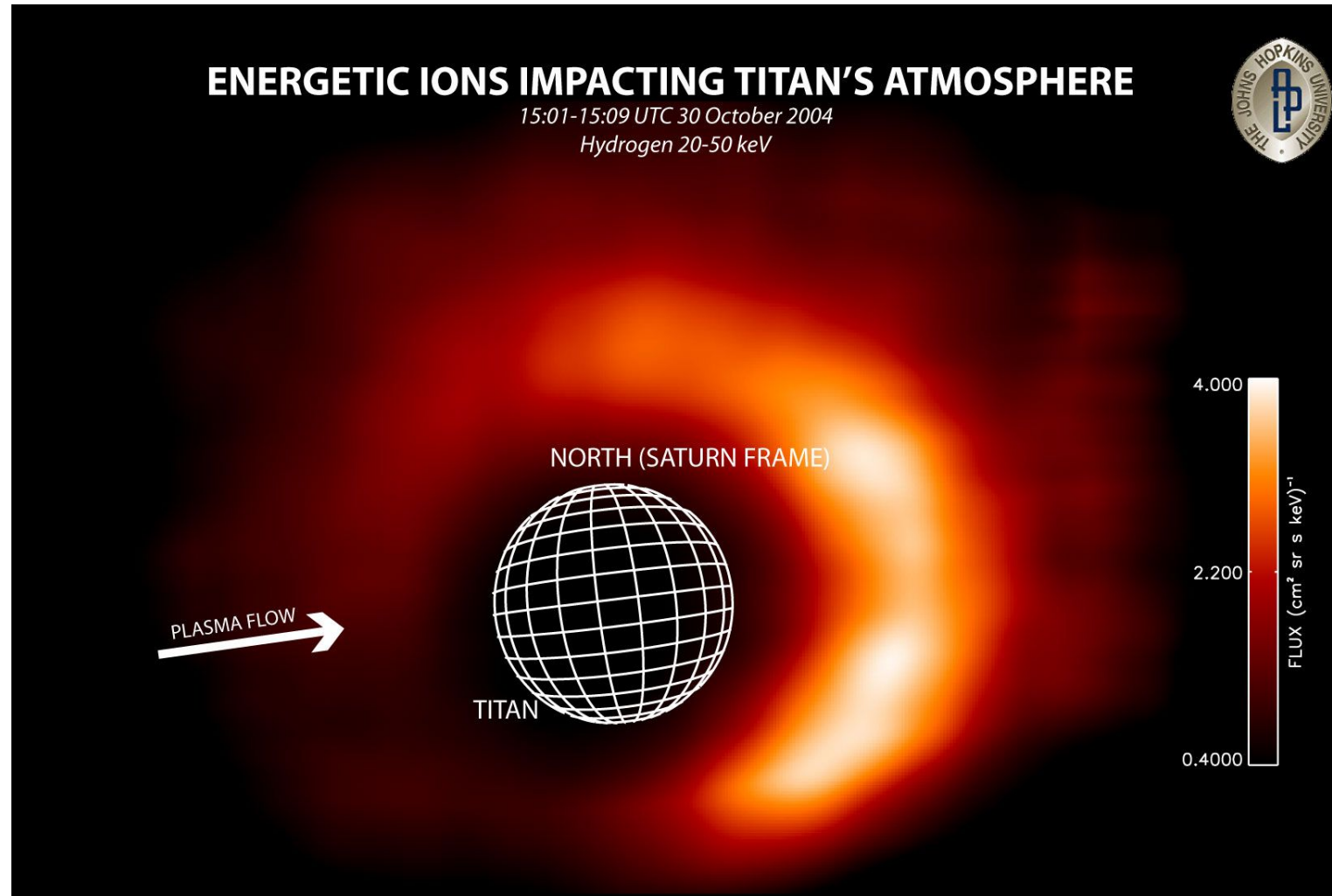
Secondary ion emission



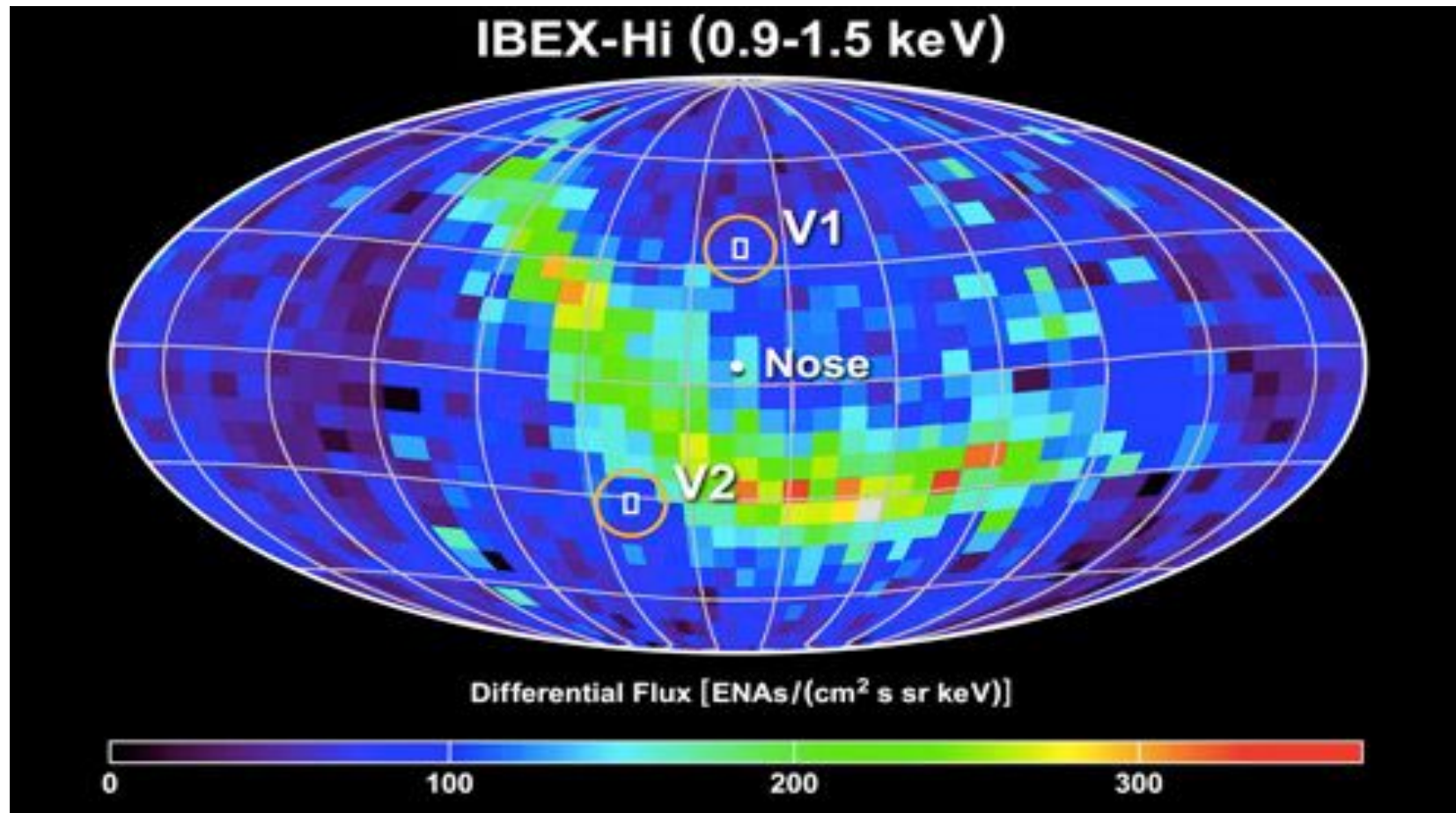
In Situ → Remote Techniques: ENA Imaging



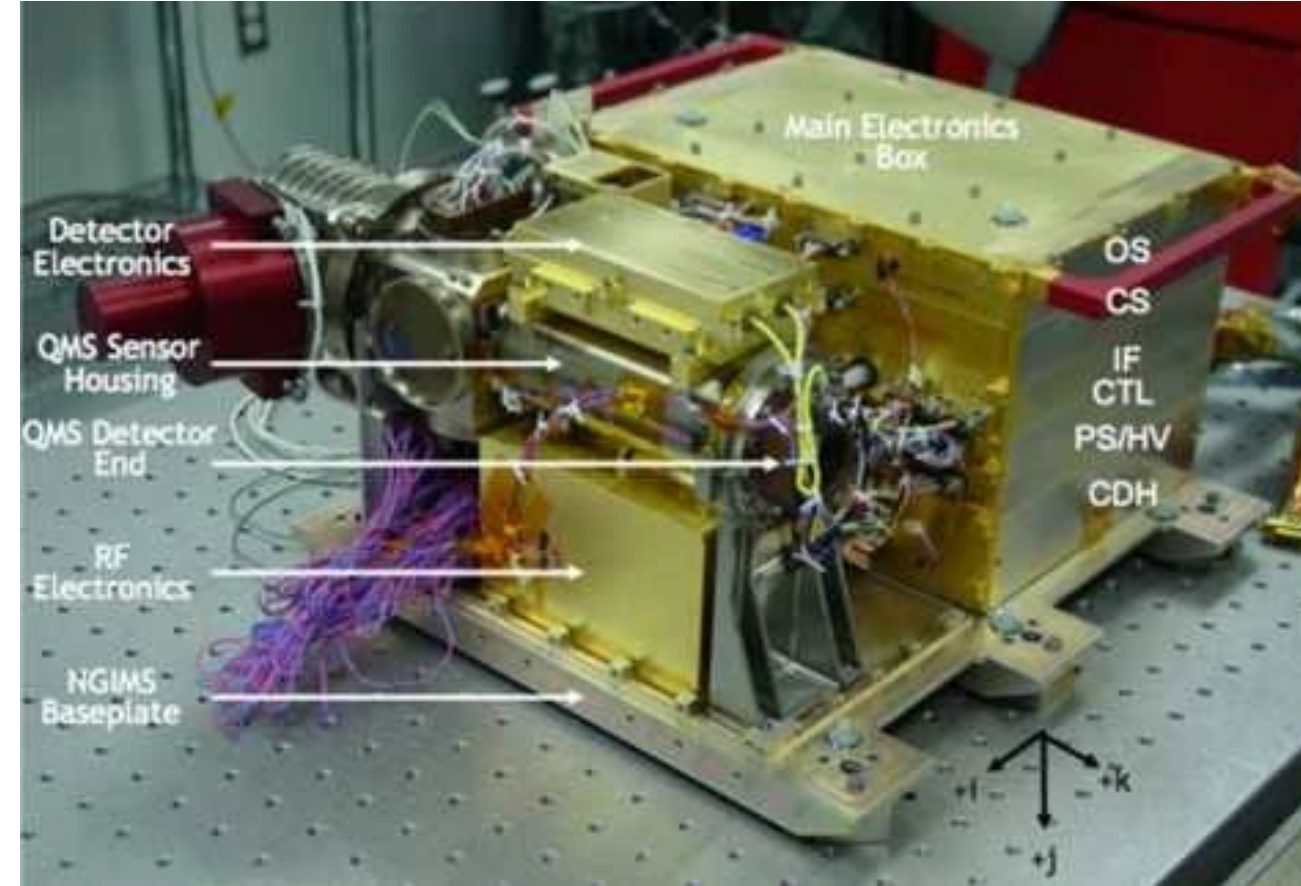
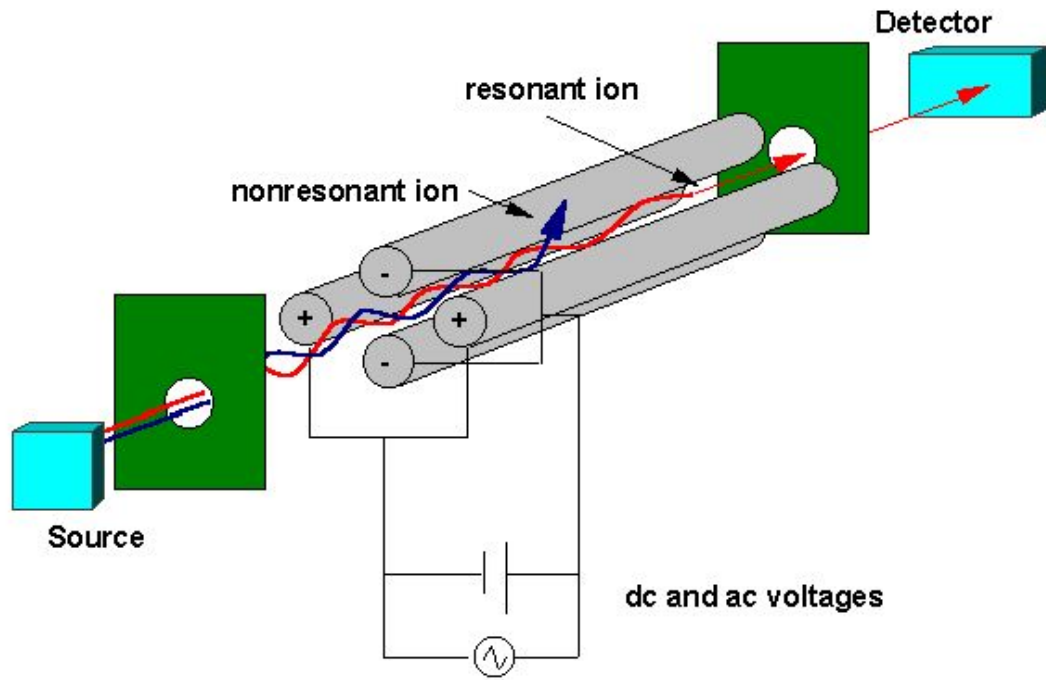
In Situ → Remote Techniques: ENA Imaging



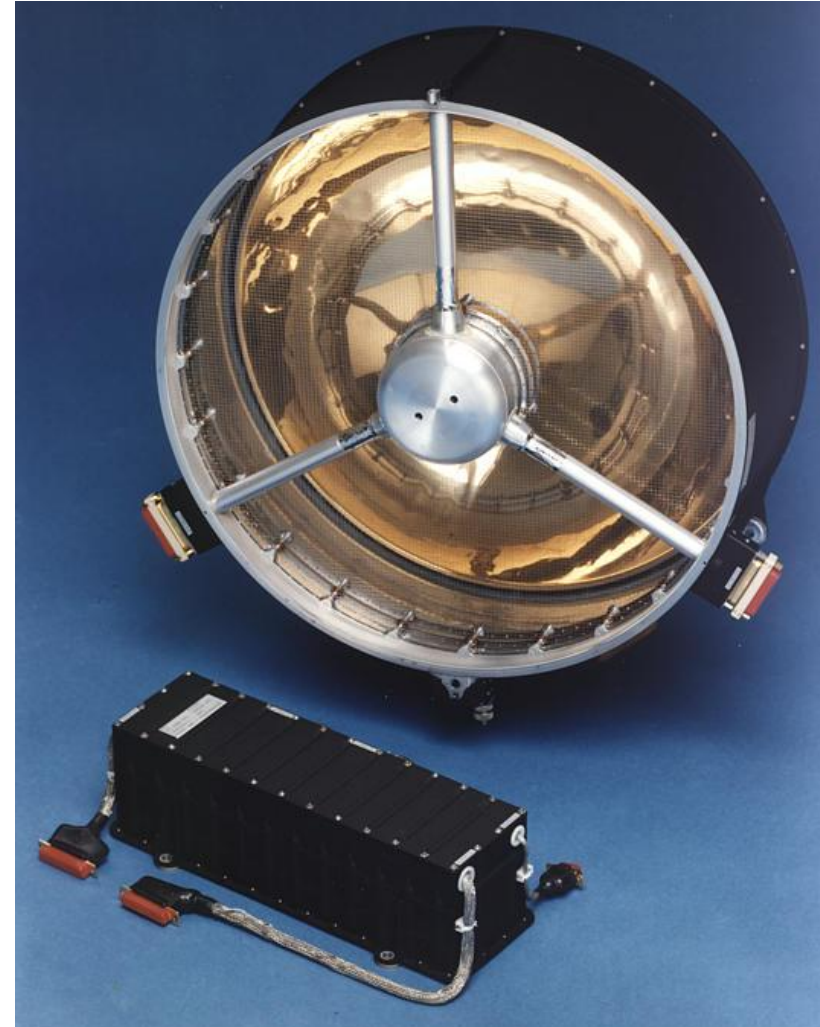
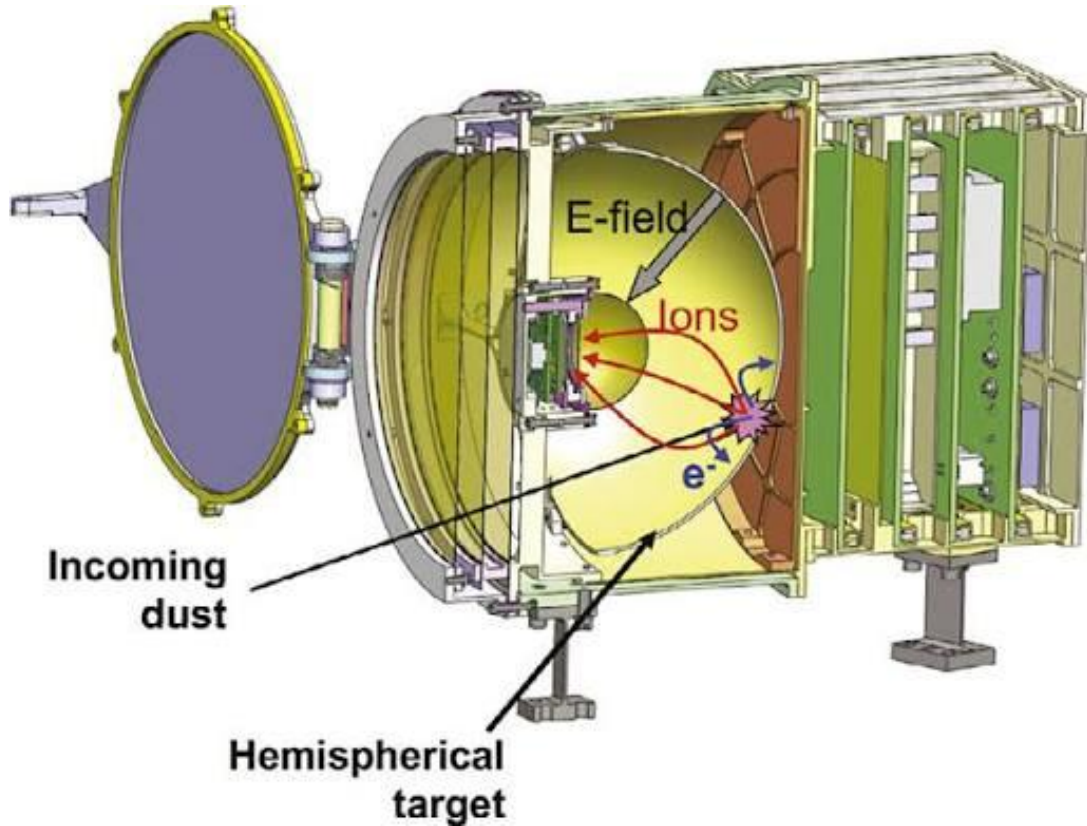
In Situ → Remote Techniques: ENA Imaging



In Situ Techniques: RF Resonance

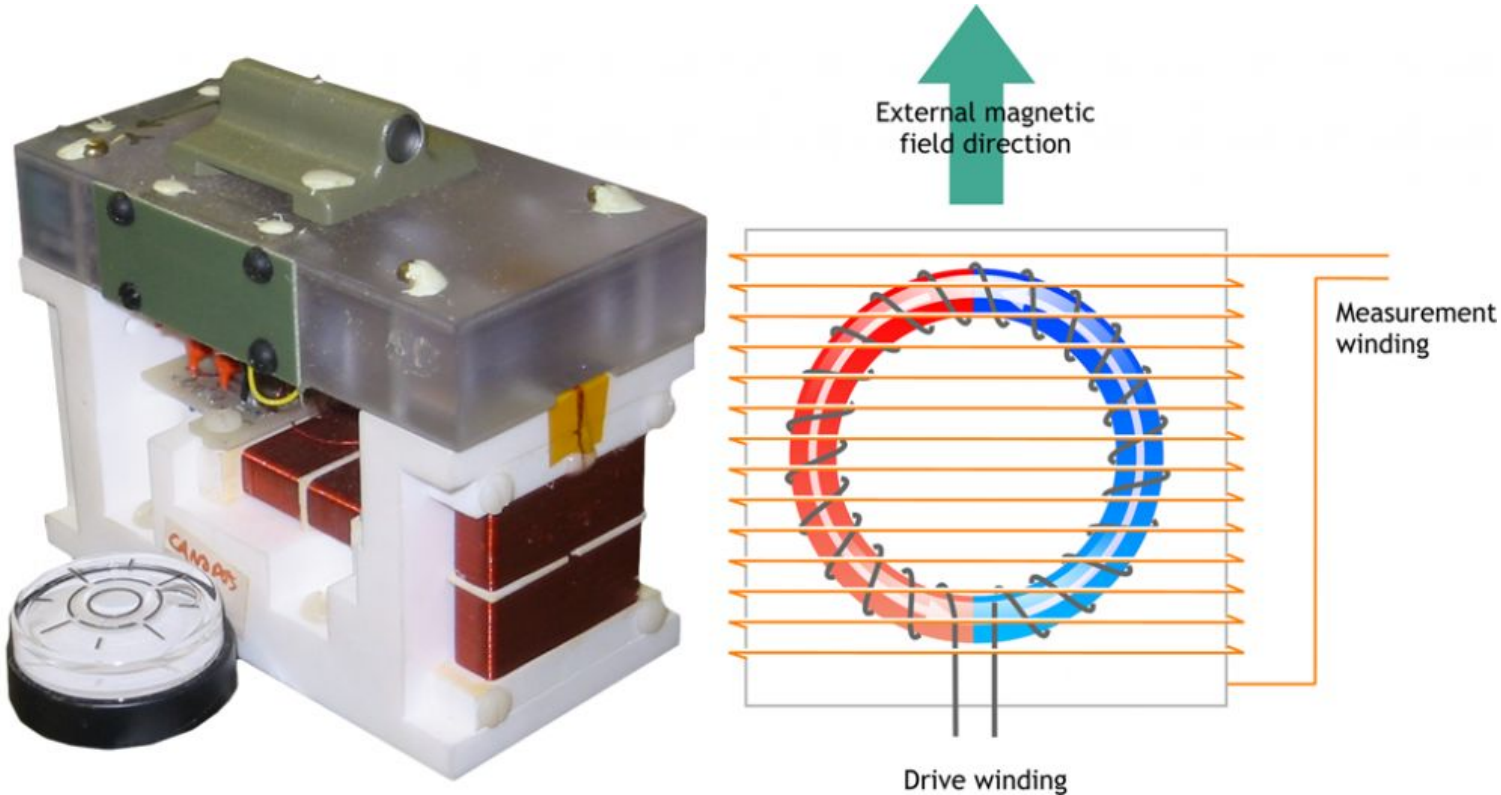


In Situ Techniques: Dust Impact Sensors



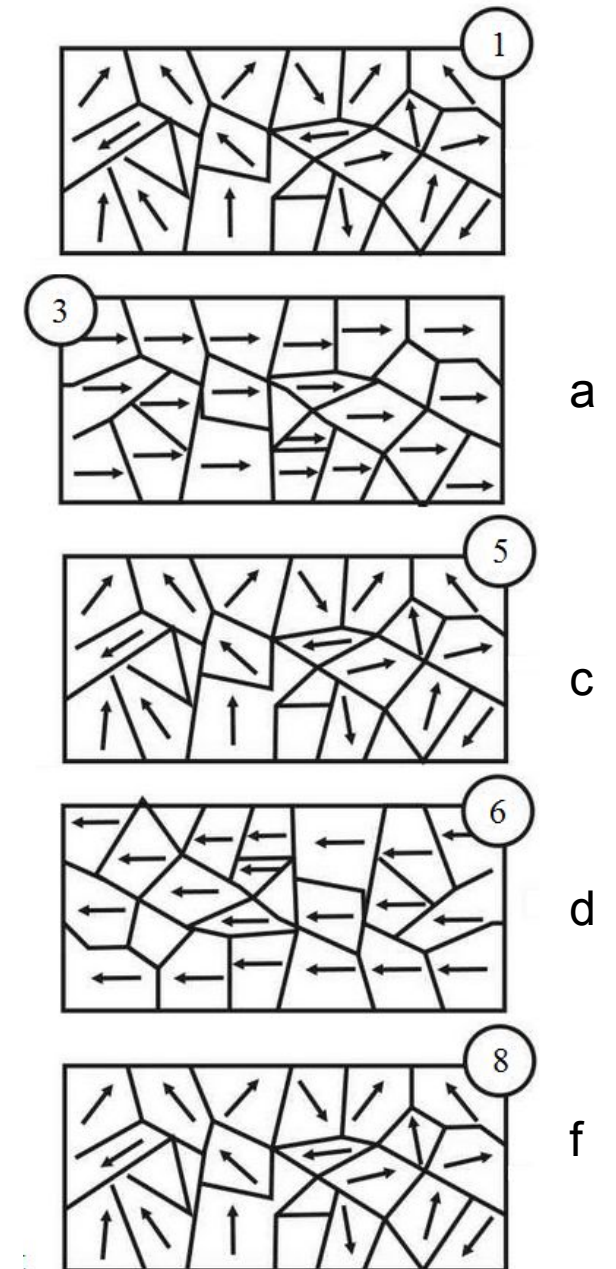
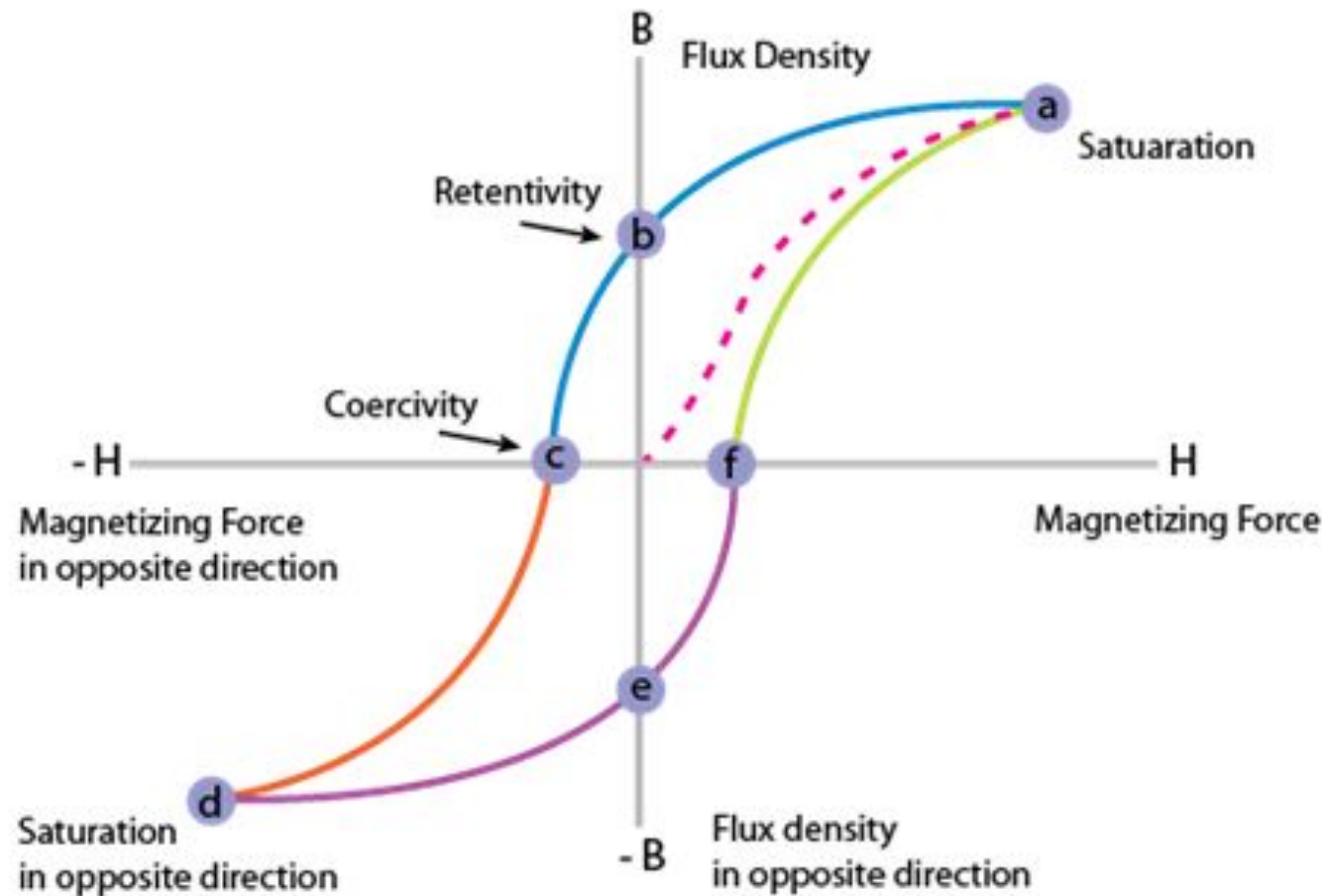
Fields

In Situ Techniques: Magnetic Field - Fluxgate

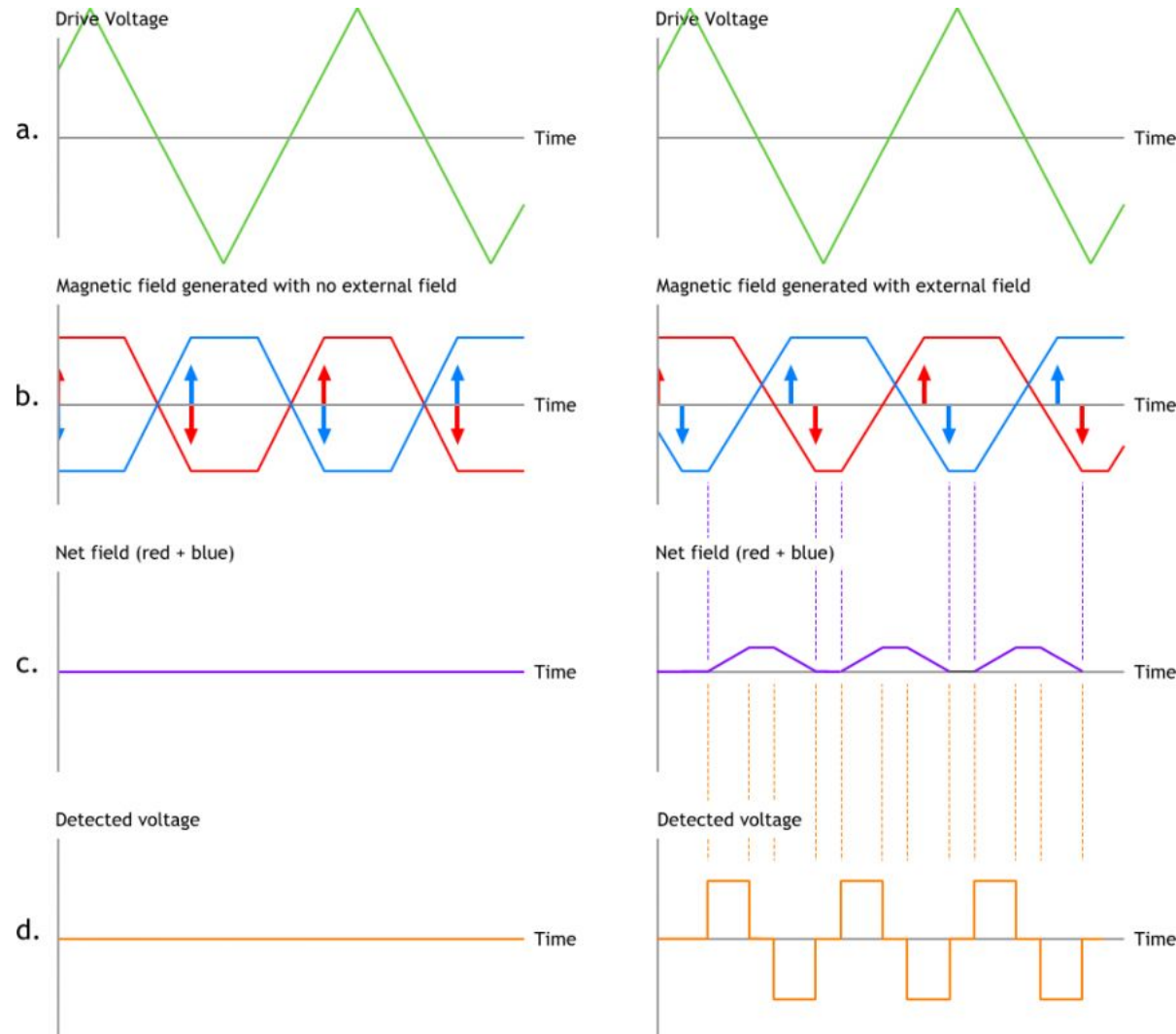


$$\begin{aligned}\mathcal{E} &= -\frac{d\Phi_B}{dt} \\ &= -\frac{d(\mu H N A)}{dt} \\ &= -(\mu N A) \frac{dH}{dt} - (N H A) \frac{d\mu}{dt}\end{aligned}$$

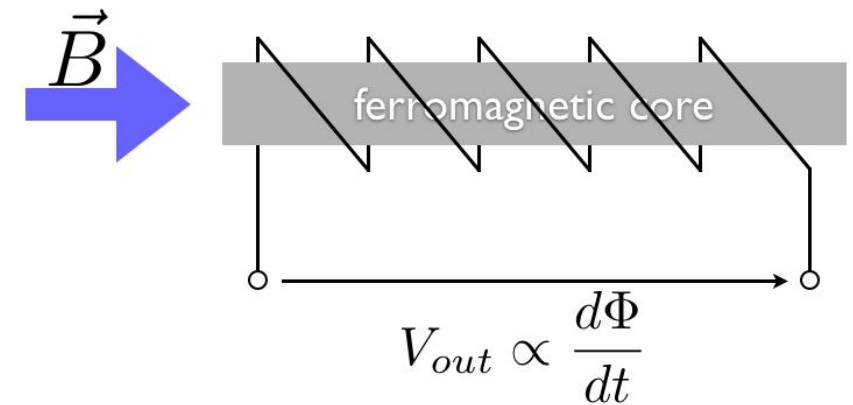
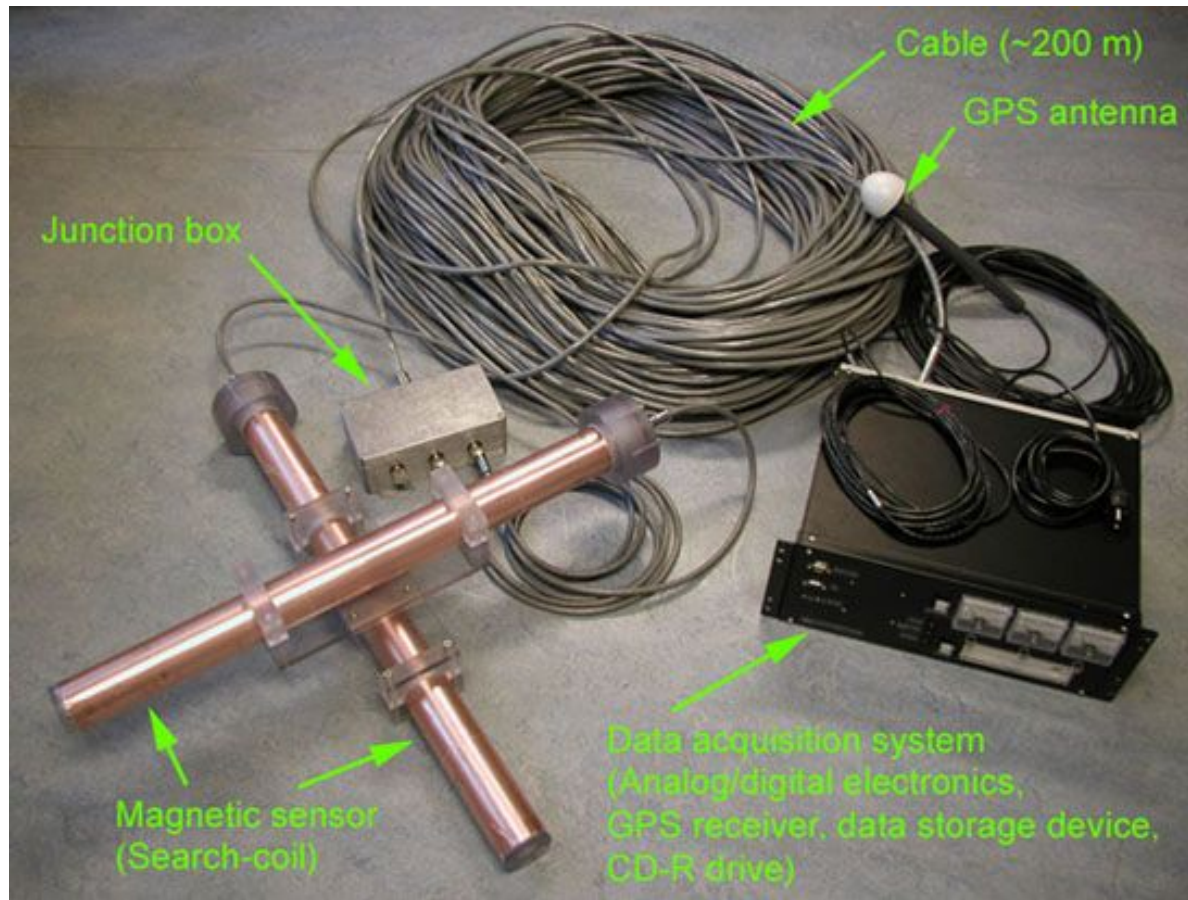
Magnetic Hysteresis



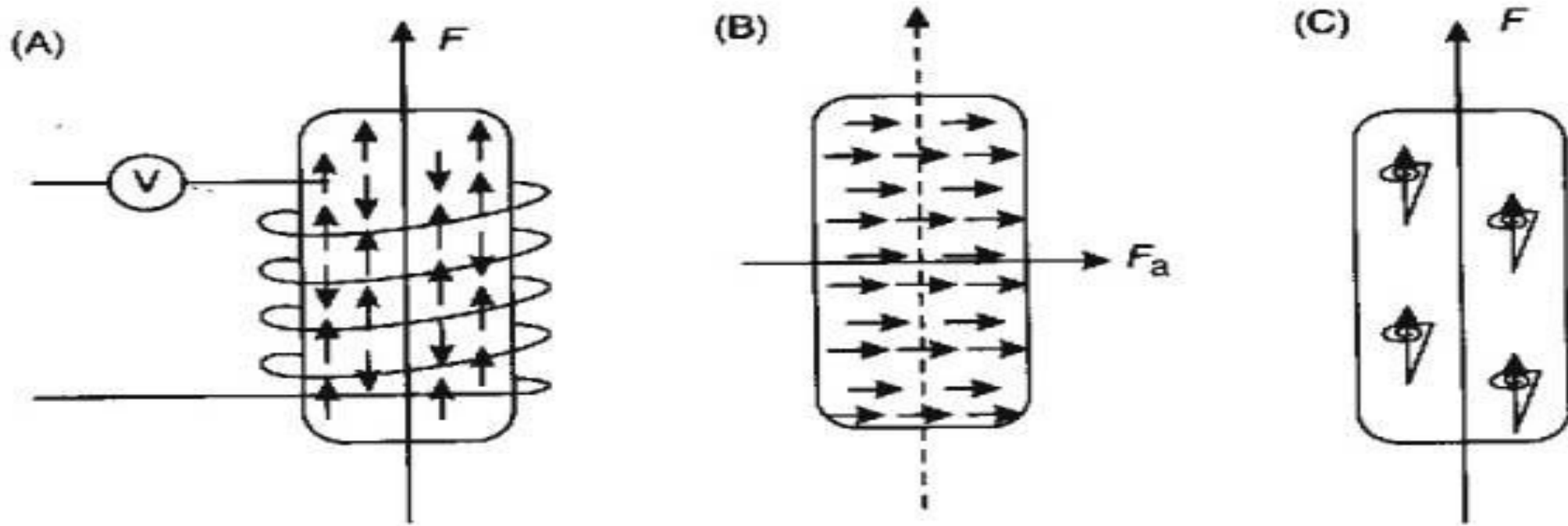
In Situ Techniques: Magnetic Field - Fluxgate



In Situ Techniques: Magnetic Field – Search Coil



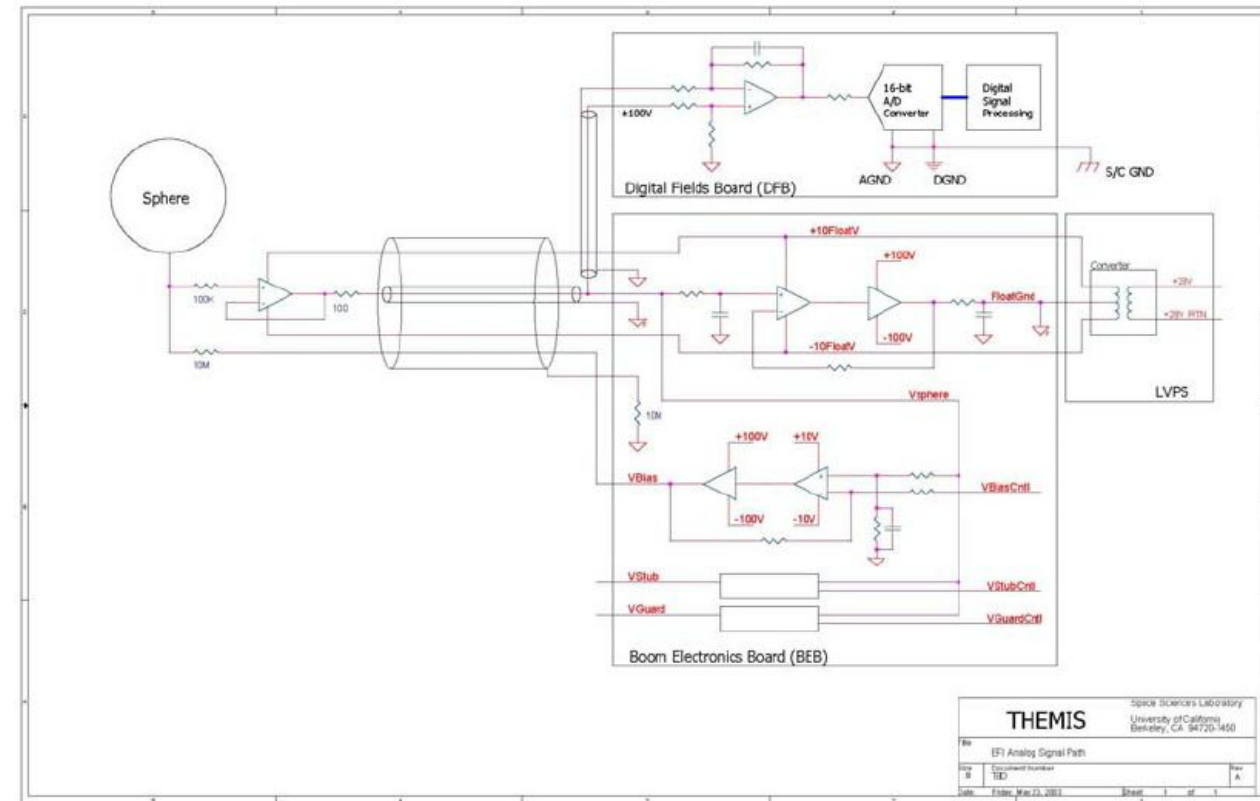
In Situ Techniques: Magnetic Field – Proton Precession



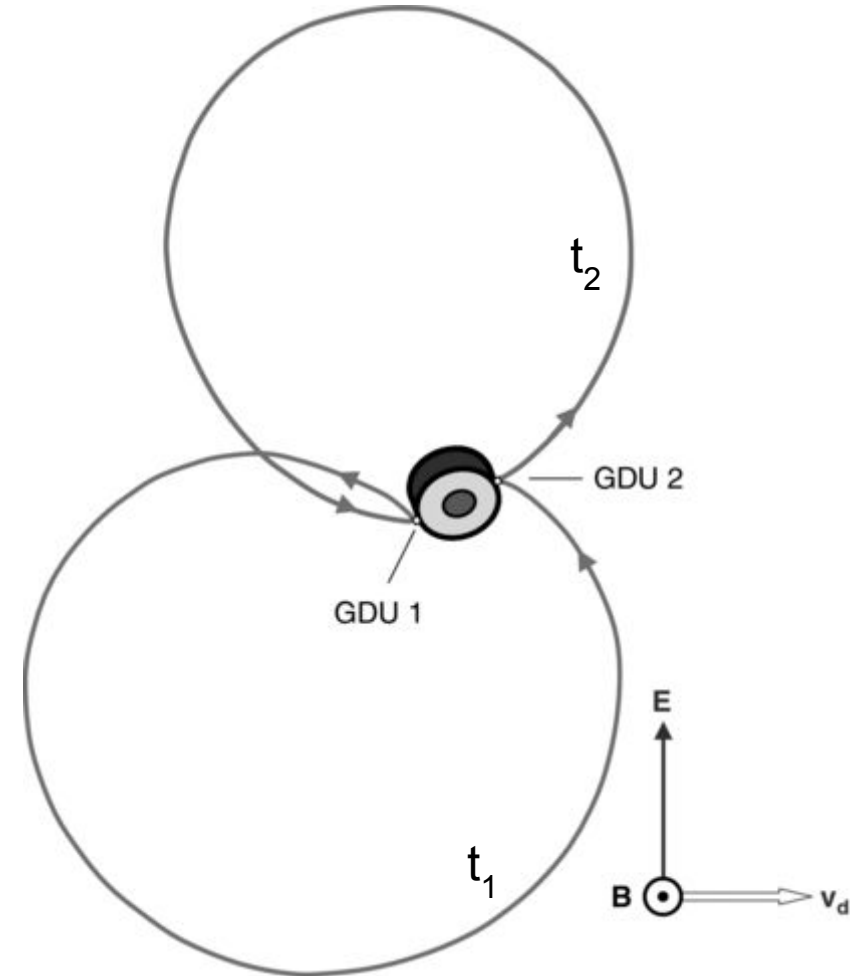
IOWA

Diagram illustrating the boom configuration of the MMS spacecraft:

- 2x $\sim 5\text{m}$ long axial EFI booms
- 1m long SCM boom
- 2m long FGM boom
- 4x $\sim 20\text{m}$ long radial EFI booms

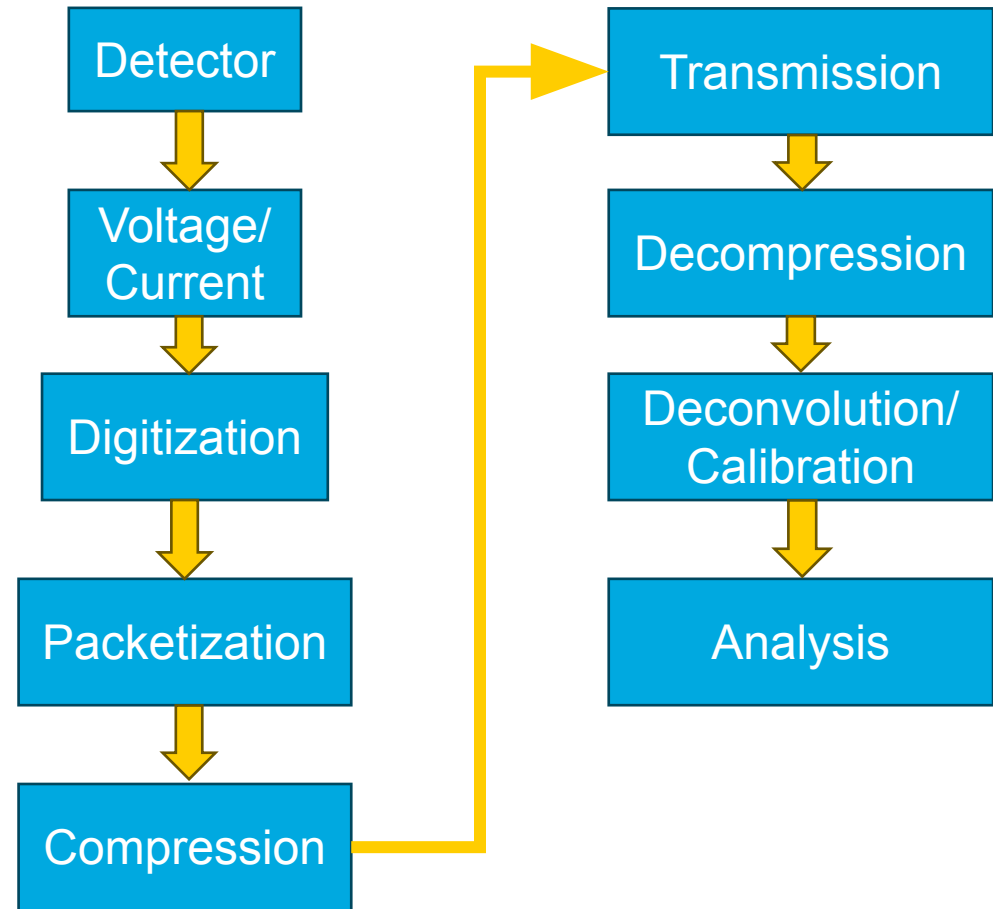
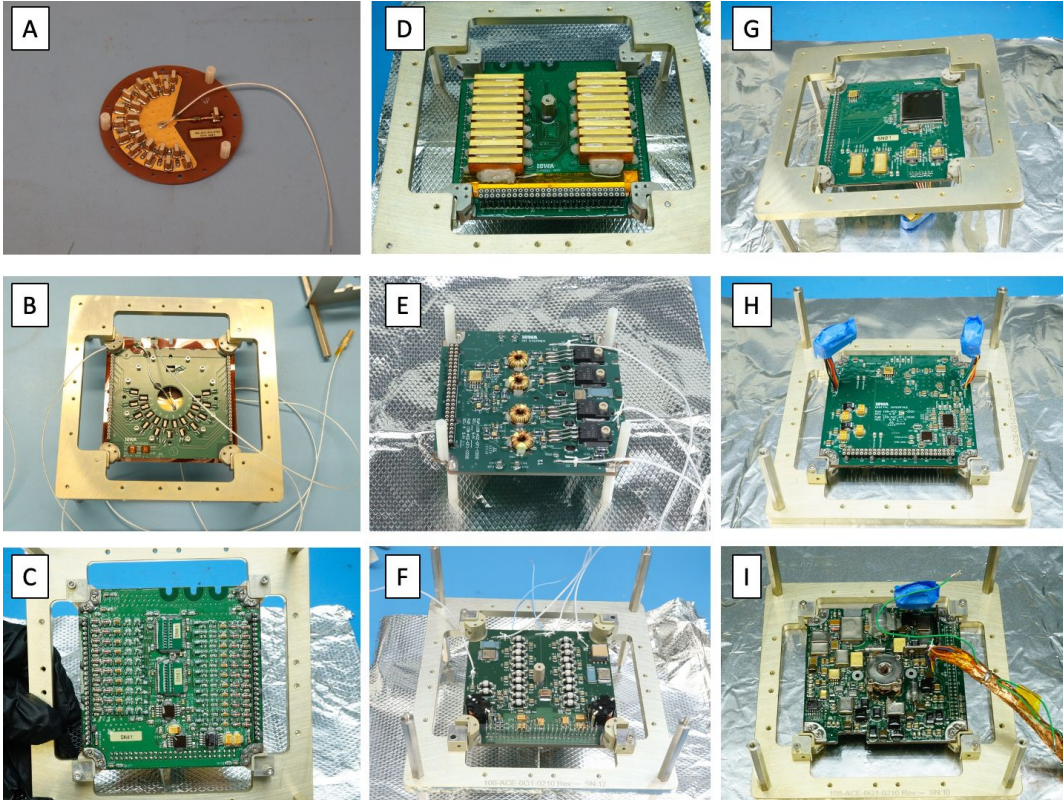


In Situ Techniques: Electric Field Measurements w/ Particles



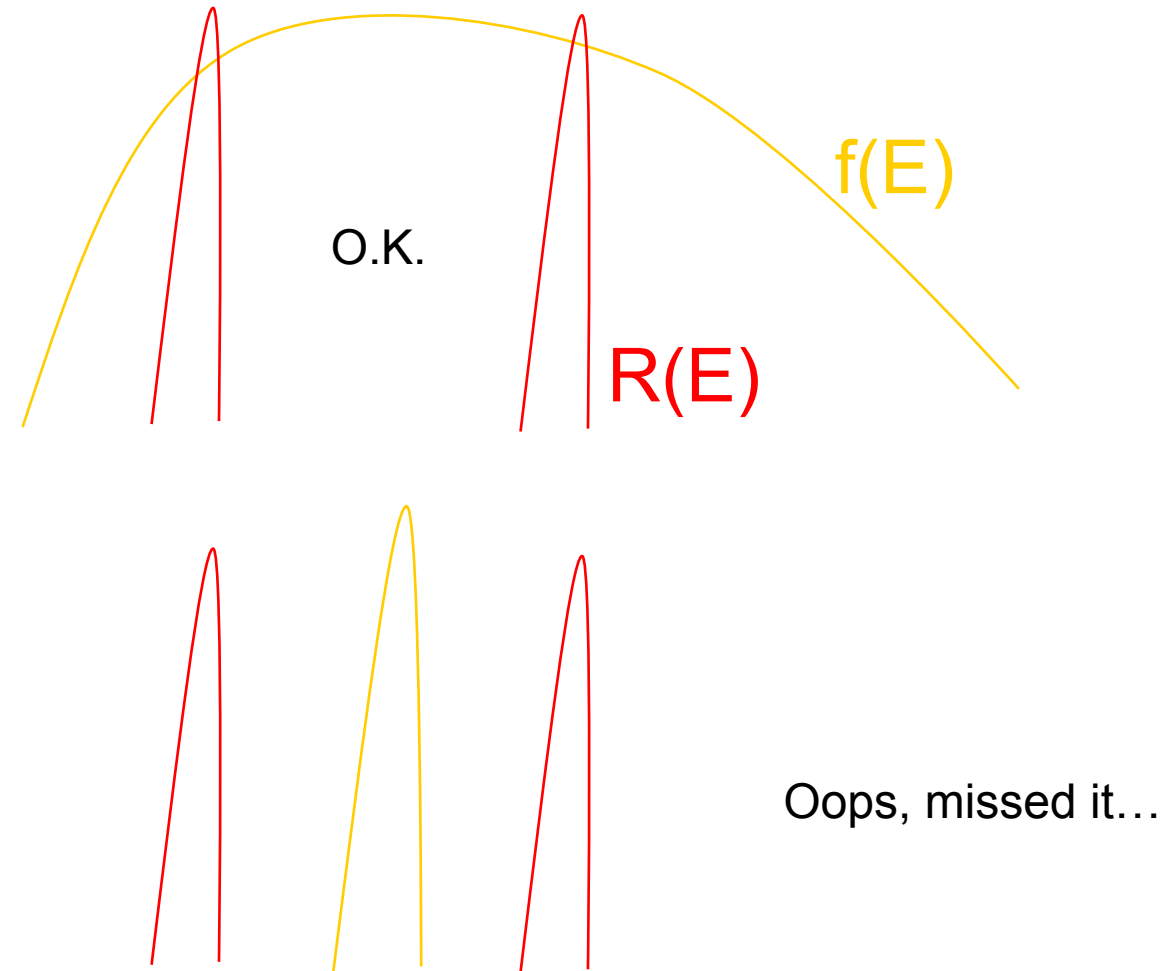
Some Issues

How to Make a Measurement

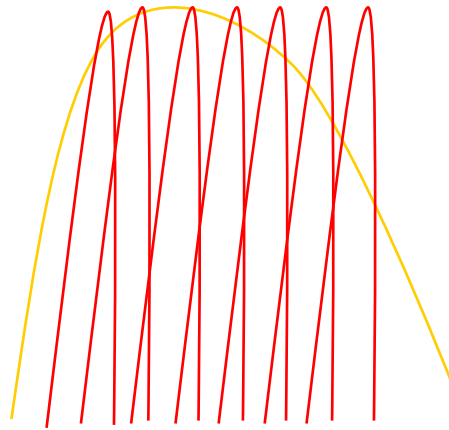


Every step may introduce errors...

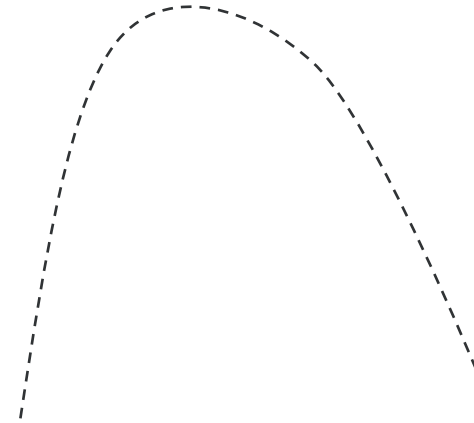
Sampling Issues



Deconvolution Issues



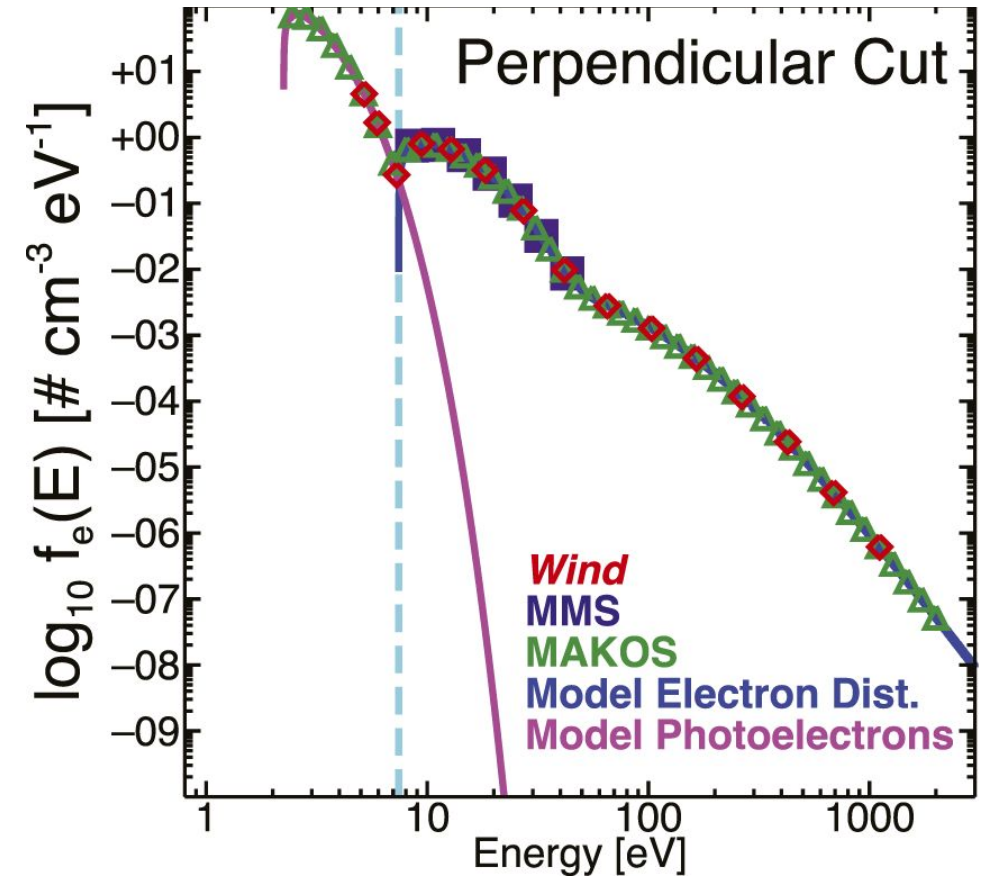
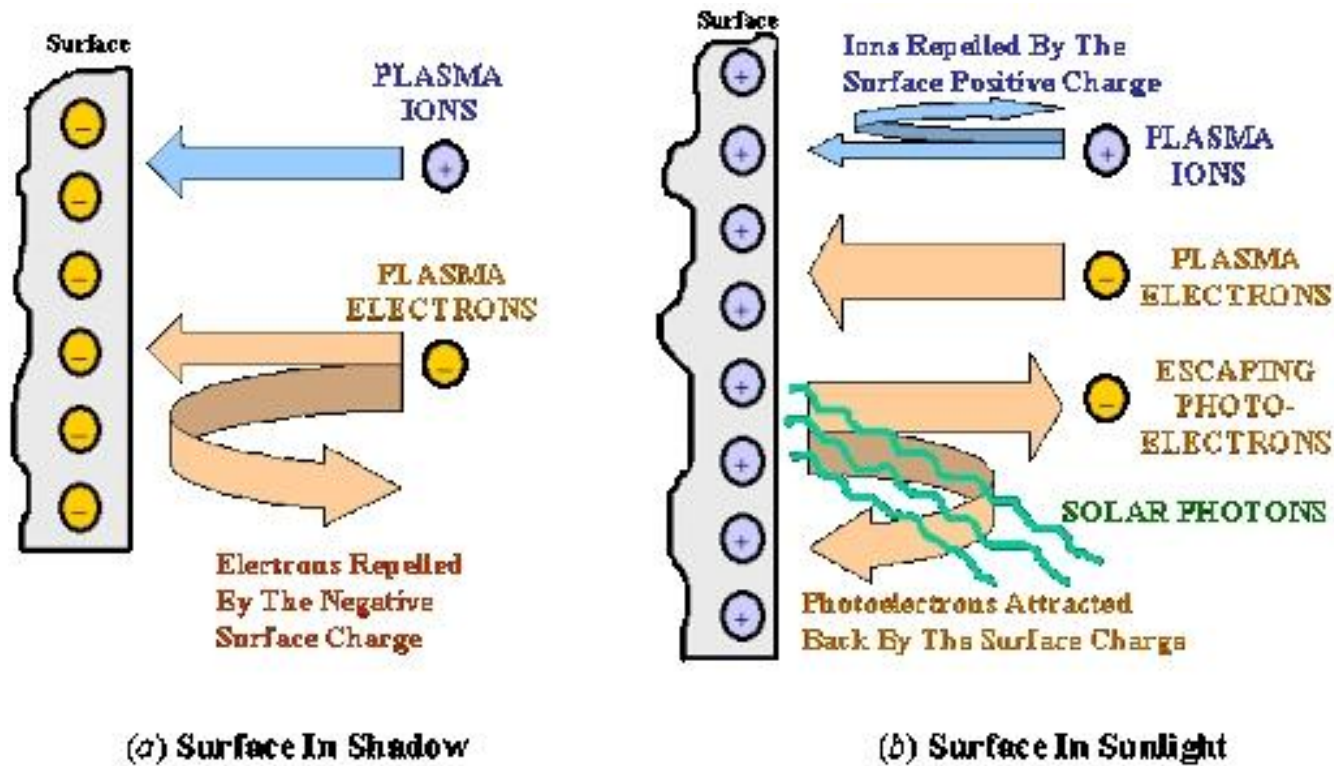
O.K.



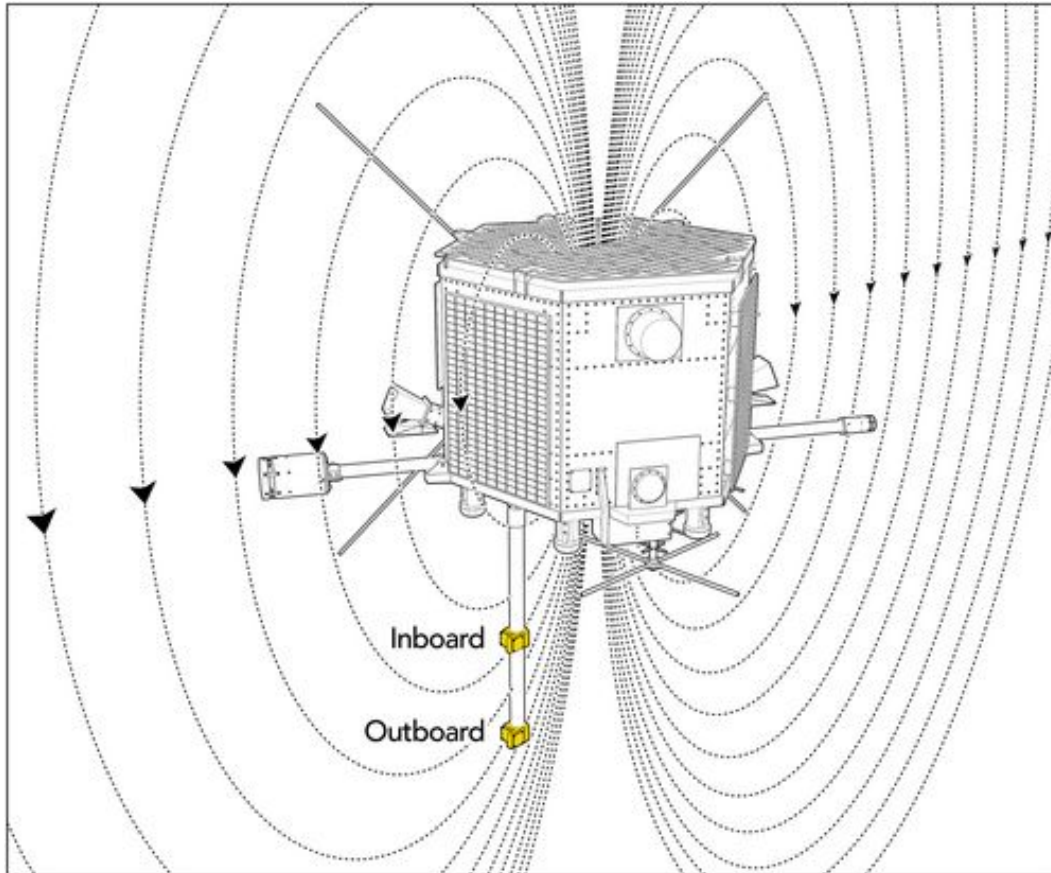
Under-resolved



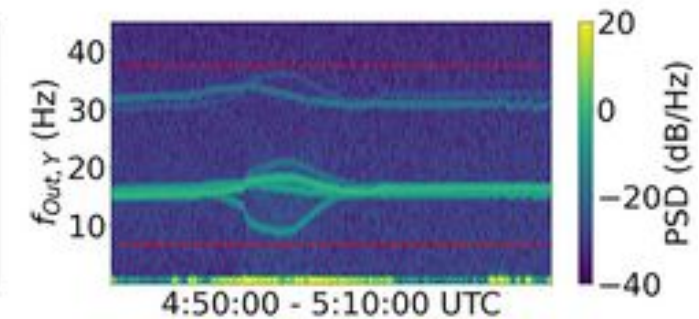
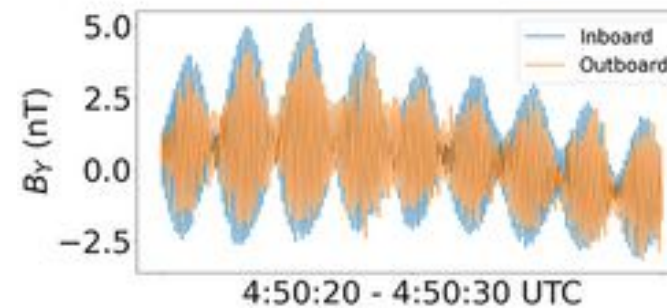
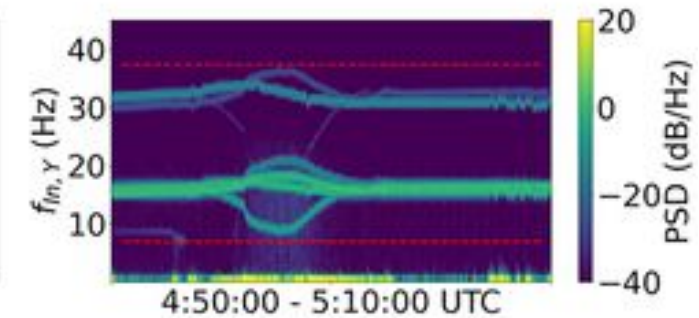
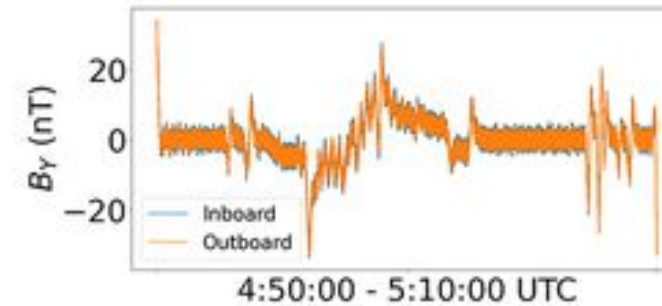
There's a Spacecraft in the Way



There's a Spacecraft in the Way



e-POP MGF (16 March 2016)



Key Takeaways

- The physics of measurement devices can be as complex as the physics of the systems they measure
 - To understand the measurement, you have to understand the device and its limitations
- No measurement device is perfect
 - Limited by sensitivity/dynamic range
 - Limited by finite resolution – in all dimensions
 - Limited by counting statistics
 - Limited by systematic errors, offsets, and/or background
 - Limited by calibration uncertainties
- The measurement platform affects the measurement
 - Every measurement platform produces its own electromagnetic fields
 - Every measurement platform is both a source and a sink of particles
 - The measurement platform affects the surrounding plasma over at least a Debye length, and sometimes more

Backup

The Problem

- Plasma instruments do not measure the distribution function
- Plasma instruments usually measure something related to flux
 - Typical observable:
 - Number of particles entering aperture with area A per measurement period T
 - Normalize by A , T to get flux $J = \#/[cm^2 s]$

Differential Measurements

- Charged particle measurements are differential
 - We only measure part of the distribution at a time
- Measurement are often stated in terms of differential energy flux $[dJ_E/dE]$:
 - $\text{eV}/[\text{cm}^2 \text{ s sr eV}]$
 - Count Rate = $\#/s = dJ_E/dE * A * \Delta E/E * \Delta\Omega = G * dJ_E/dE$
 - G = Geometric Factor [units of area]

Converting Units

- Measured flux $J = n \langle \mathbf{v} \cdot d\mathbf{A} \rangle = \int [f(v) \mathbf{v} \cdot d\mathbf{A} d^3\mathbf{v}] = \# / (\text{cm}^2 \text{ s})$
- $\mathbf{v} \cdot d\mathbf{A} d^3\mathbf{v} = \mathbf{v} \cdot d\mathbf{A} v^2 dv \sin\theta d\theta d\phi = v^3 dA dv d\Omega$
- $E = \frac{1}{2} mv^2 \Rightarrow mv dv = dE$
- $v^3 dA dv d\Omega = 2E/m^* dE/m dA d\Omega = 2E/m^2 dE dA d\Omega$
- $J = \int [f * 2E/m^2 dE dA d\Omega]$
- $dJ/dE = f * 2E/m^2$ [differential in E, Ω , A]
- $dJ_E/dE = f * 2E^2/m^2$