

# Introduction to Heliophysics

Dana Longcope

Montana State University



Total Force.

Having since had occasion to examine the disturbances of the Declination at the same two stations in the three succeeding years 1846, 1847 and 1848, I have had the satisfaction of finding that the observations of these years confirm every deduction which

period: *Phil. Trans. Royal Soc.* **142**, 103 (1852)

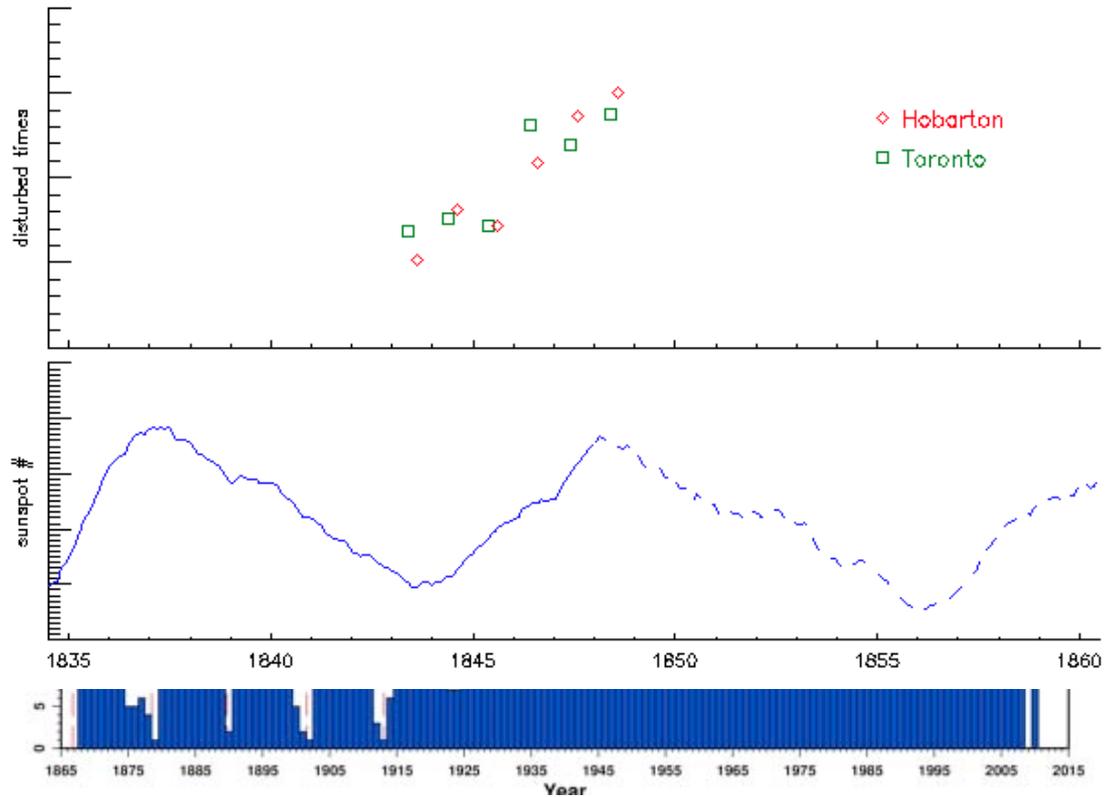
in comparison of the frequency and amount of the disturbances in different years, apparently indicating the existence of a *periodical variation*, which, either from a causal con-

some periodical affection of an outer envelope, (the photosphere,) of the sun; and it is certainly a most striking coincidence, that the period, and the epochs of minima and maxima, which M. SCHWABE has assigned to the variation of the solar spots, are absolutely identical with those which have been here assigned to the magnetic variations.

Edward Sabine  
(1788-1883)

- 1839+ : Helped establish global network of magnetic observatories
- 1852: Discovers correlation between disturbed times (@ Earth) and cycle of sunspots discovered previously by Schwabe (1843)
- 1861 – 1871: President of Royal Society (UK)

tion into the laws of the disturbances of the Inclination and of the Total Force, for which I have not yet been able to command the necessary leisure.





William Thomson  
aka. Lord Kelvin  
(1824 – 1907)

- 1890 – 1895: President of Royal Society (UK)
- 1892 – publishes demonstration that sunspots **cannot** [sic] influence Earth's magnetic field

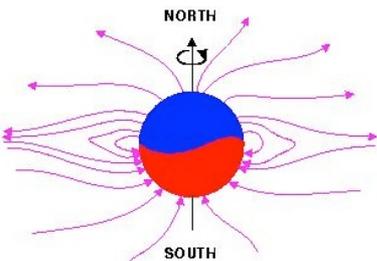
ether, we cannot say that the sun might not be 1000, or 10,000, or 100,000 times as intense a magnet as the earth. It is, therefore, a perfectly proper object for investigation to find whether there is, or is not, any disturbance of terrestrial magnetism, such as might be produced by a constant magnet in the sun's place with its magnetic axis coincident with the sun's axis of rotation. Neglecting for the present the seven degrees of obliquity of the sun's equator, and supposing the axis to be exactly perpendicular to the ecliptic, we have an exceedingly simple case of magnetic action to be considered: a magnetic force perpendicular to the ecliptic at every part of the earth's orbit and varying inversely as the cube of the earth's distance from the sun. The components of this force parallel and perpendicular to the earth's axis are respectively

i.e. a dipole

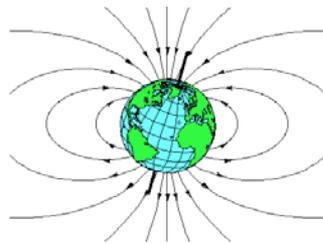
ergs per sec.) of the solar radiation. Thus, in this eight hours of a not very severe magnetic storm, as much work must have been done by the sun in sending magnetic waves out in all directions through space as he actually does in four months of his regular heat and light. This result, it seems to me, is absolutely conclusive against the supposition that terrestrial magnetic storms are due to magnetic action of the sun; or to any kind of dynamical action taking place within the sun, or in connection with hurricanes in his atmosphere, or anywhere near the sun outside.

It seems as if we may also be forced to conclude that the supposed connection between magnetic storms and sun-spots is unreal, and that the seeming agreement between the periods has been a mere coincidence.

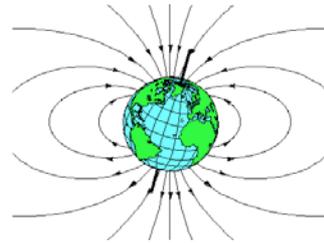
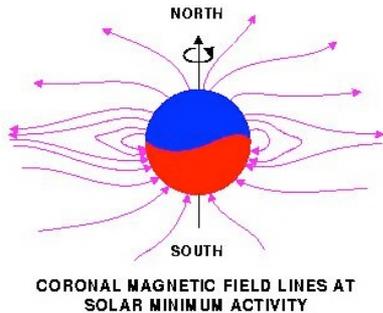
We are certainly far from having any reasonable explanation of any of the magnetic phenomena of the earth; whether the fact that the earth is a magnet; that its magnetism changes vastly, as it does from century to century; that it has somewhat



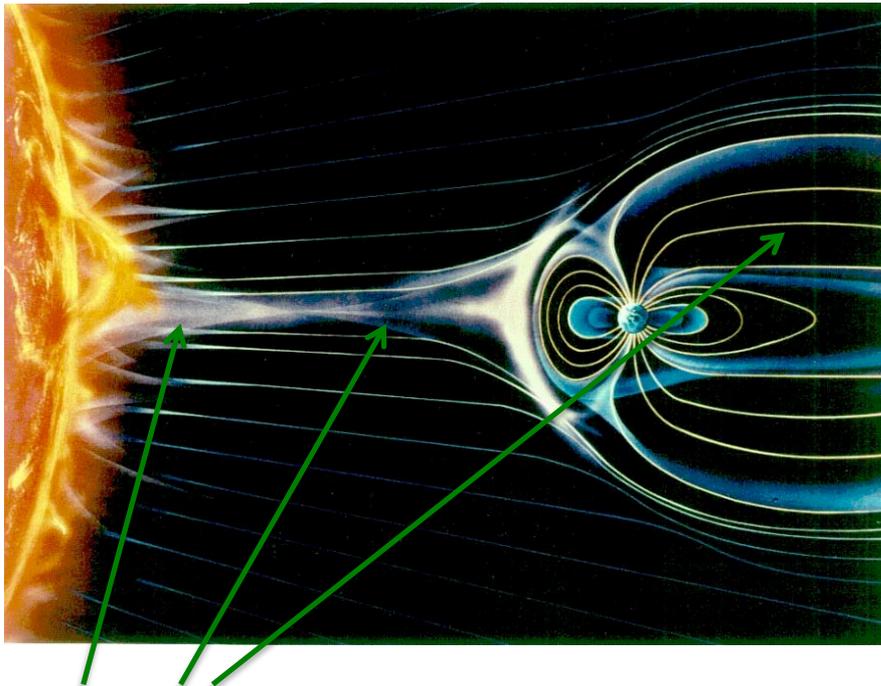
CORONAL MAGNETIC FIELD LINES AT SOLAR MINIMUM ACTIVITY



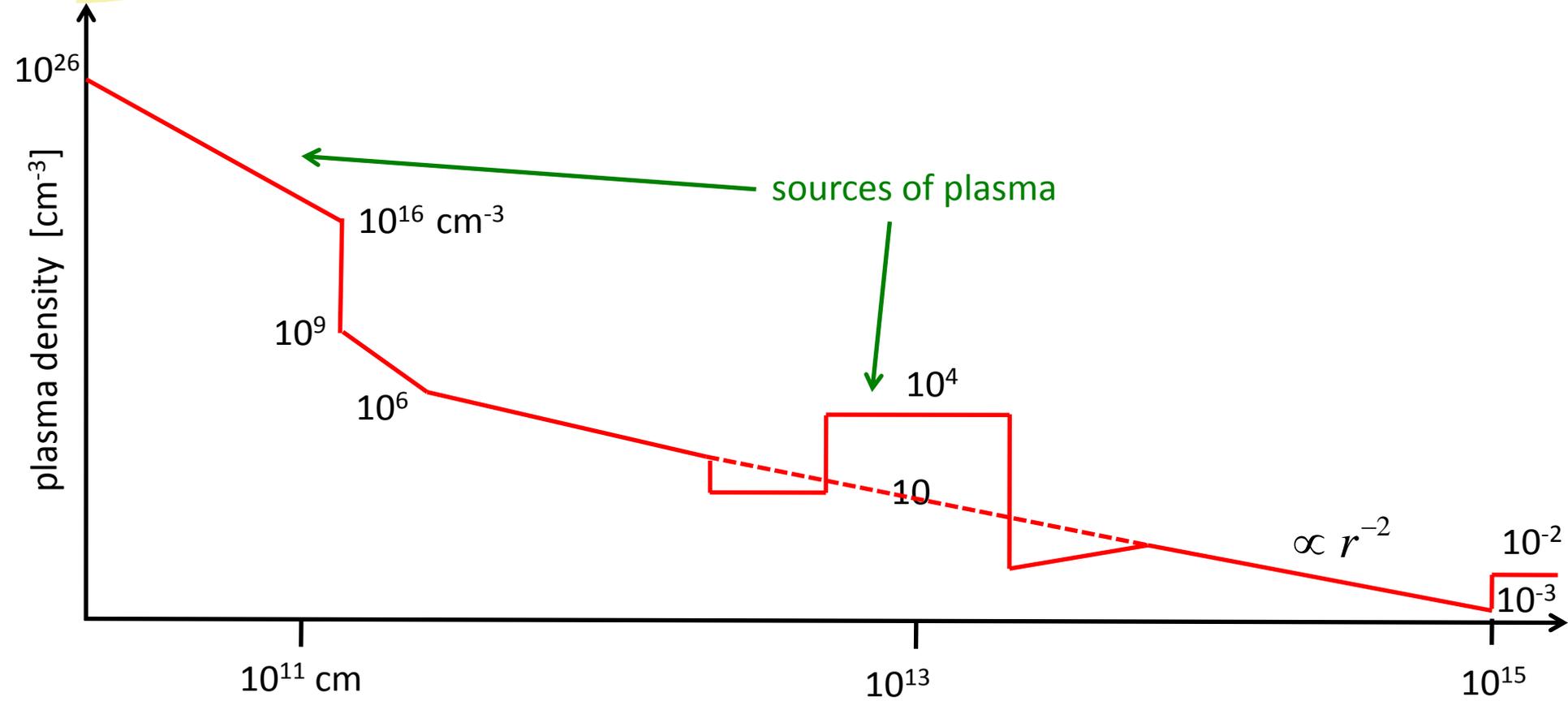
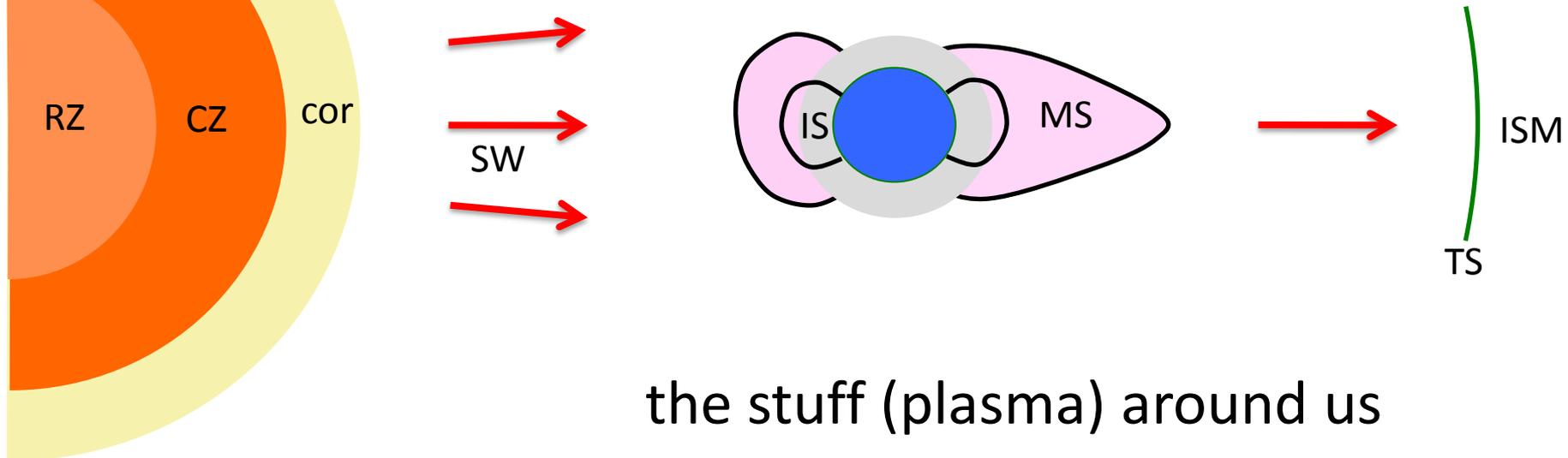
# Where did Kelvin go wrong?

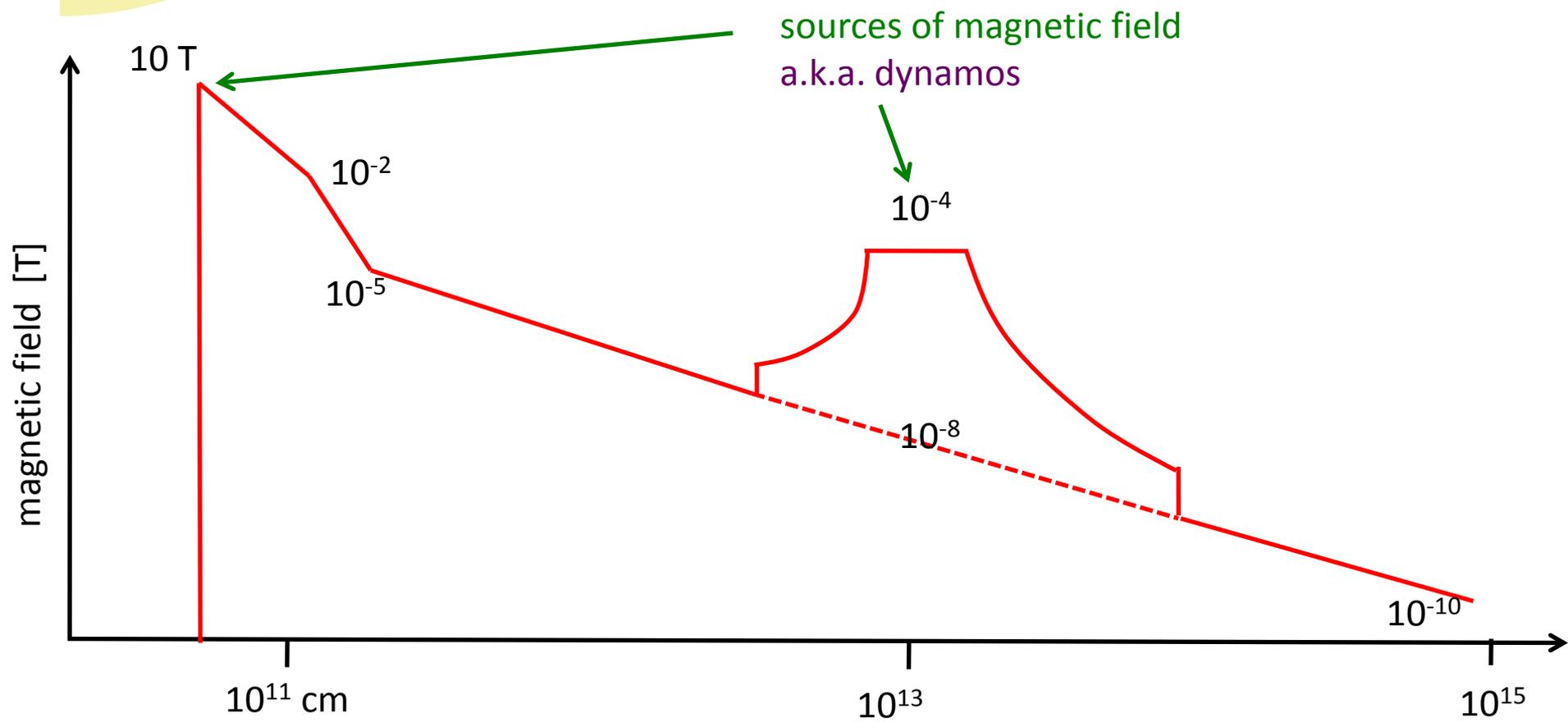
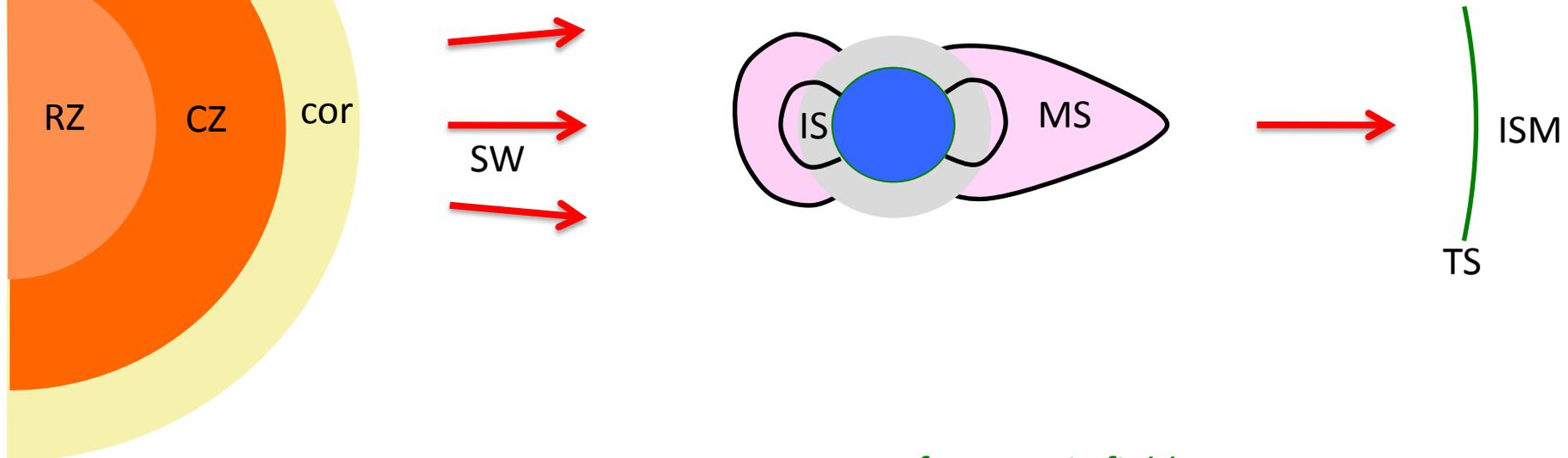


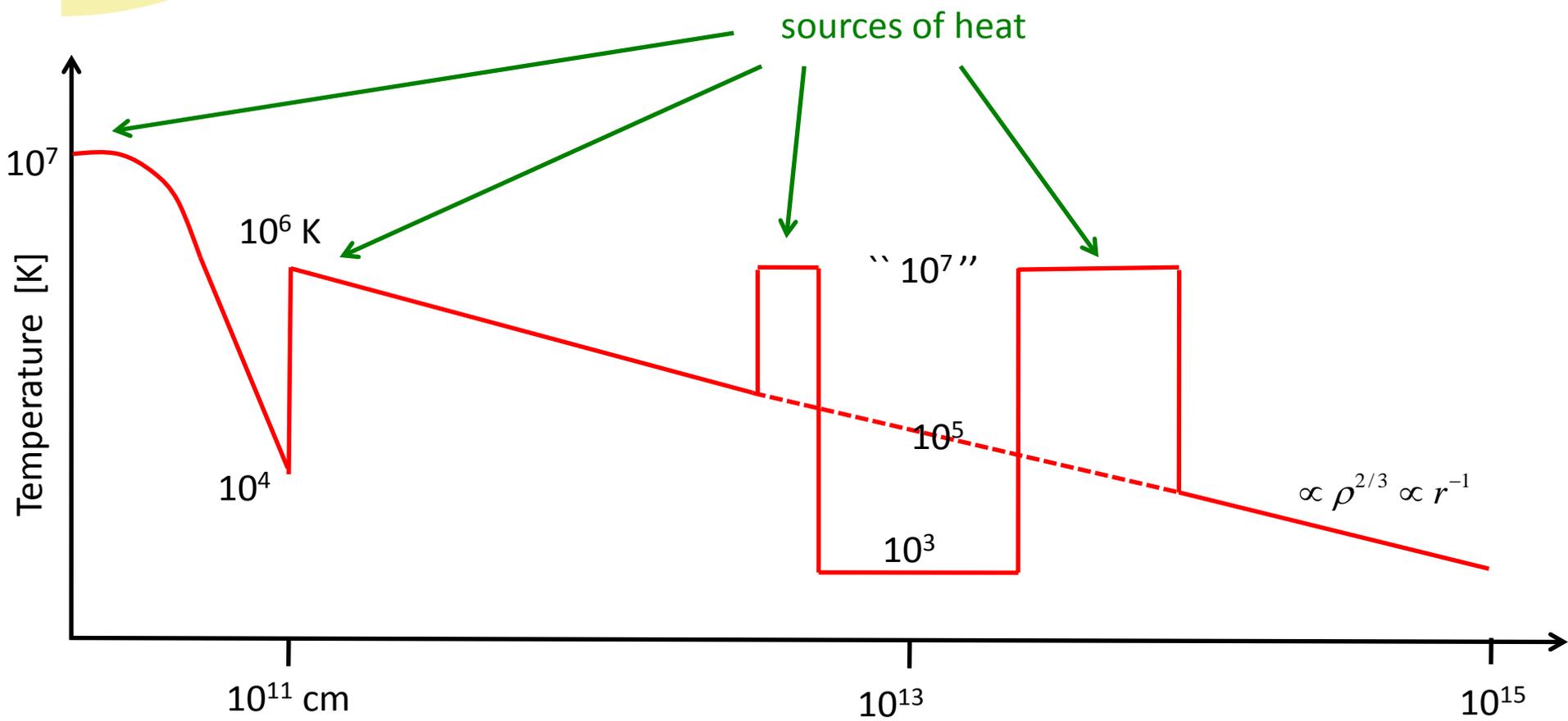
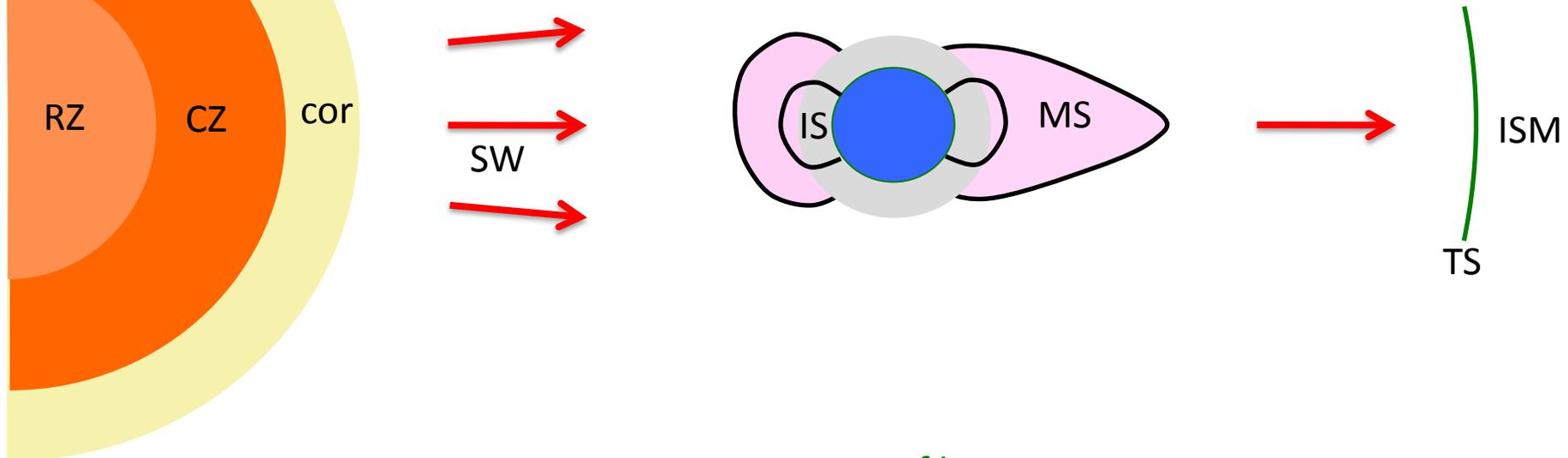
- wrong equations: “Maxwell’s”
- wrong magnetic fields
- too complex to model
- ...

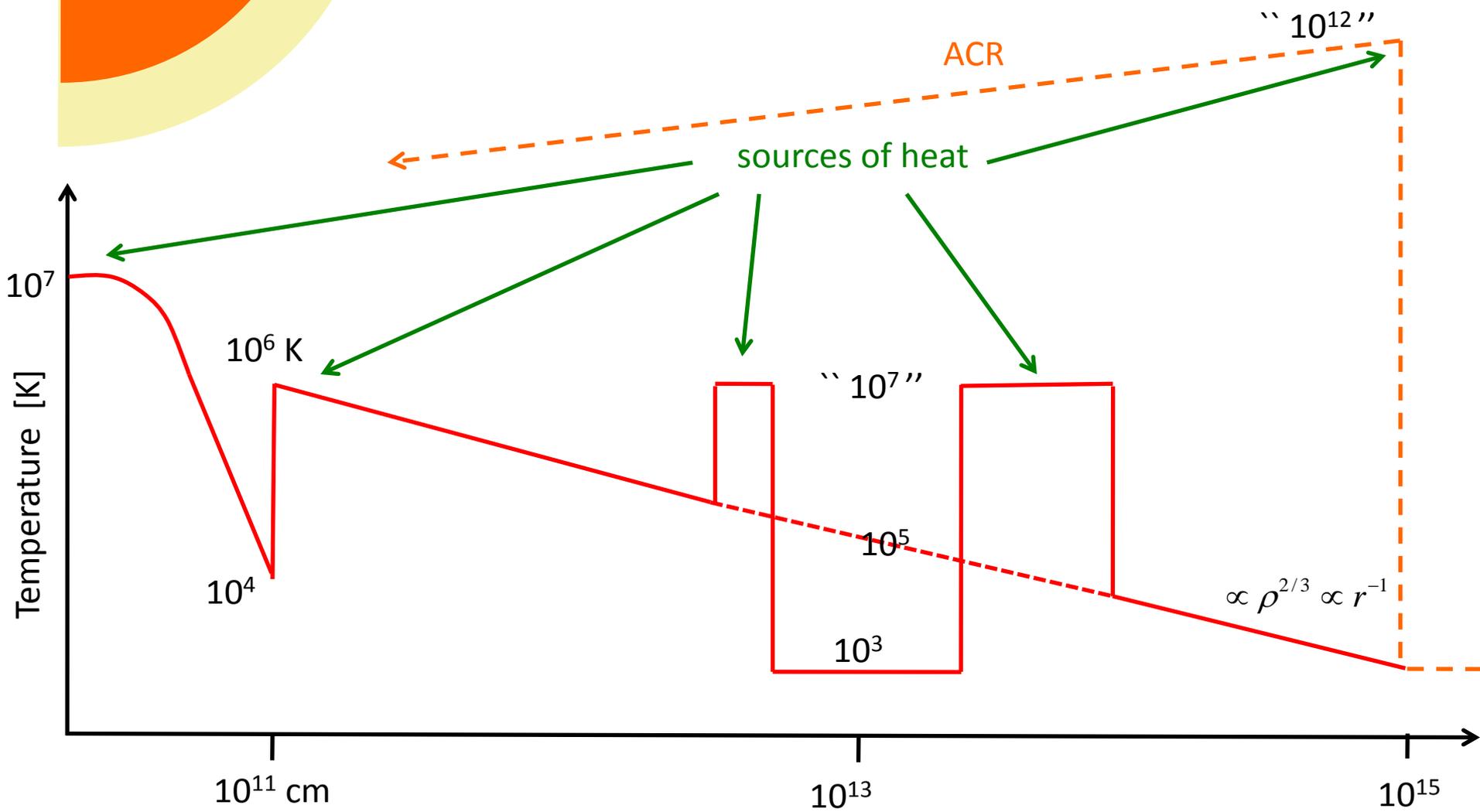
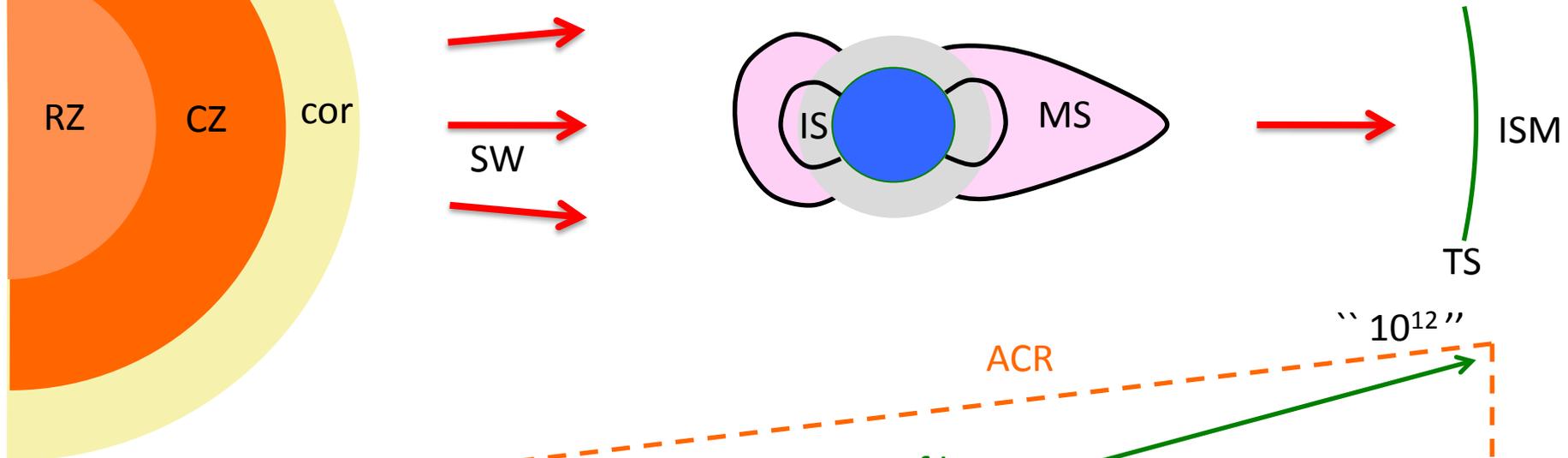


Stuff (plasma) – single system including Sun & Earth





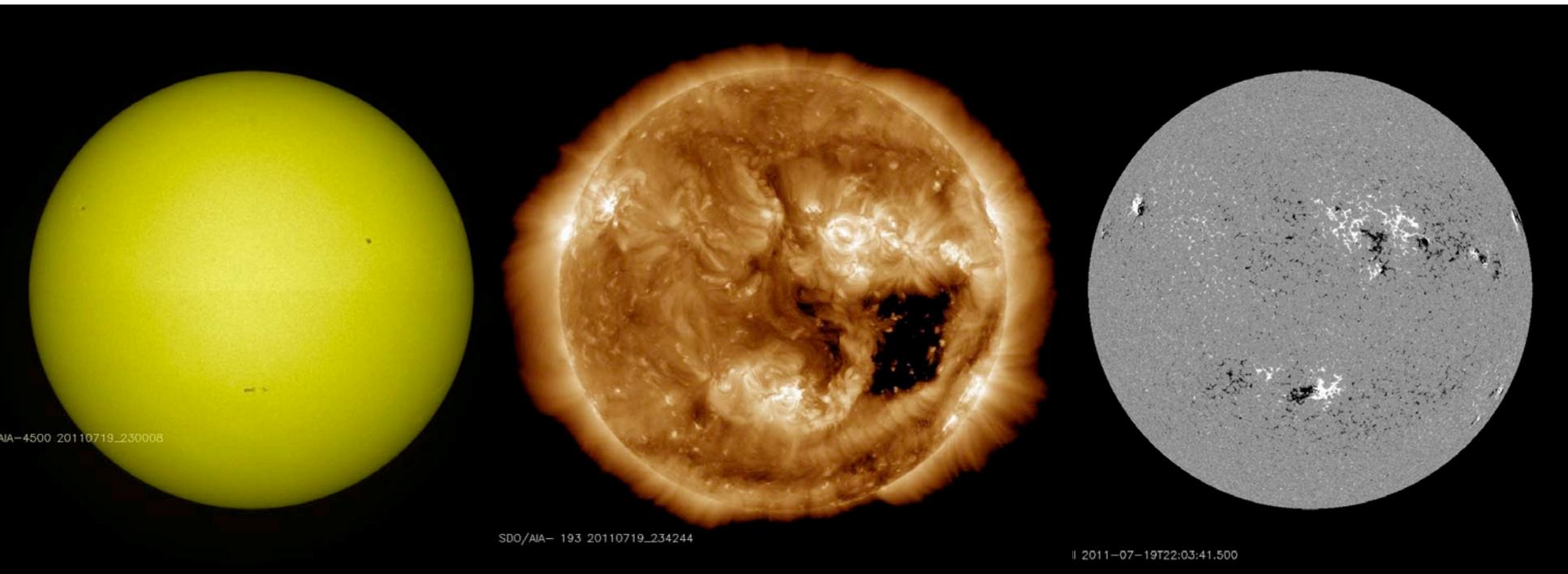




# The Sun's corona

- A heat source
- Source of the plasma flow = solar wind
- Are these unrelated features?

(vol. I, Ch. 9)

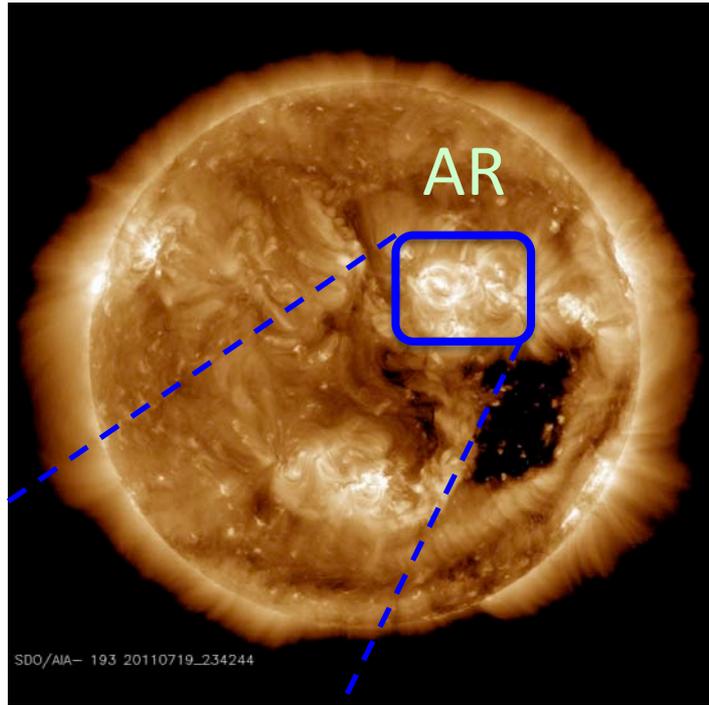


## Coronal (EUV) imaging – the basics:

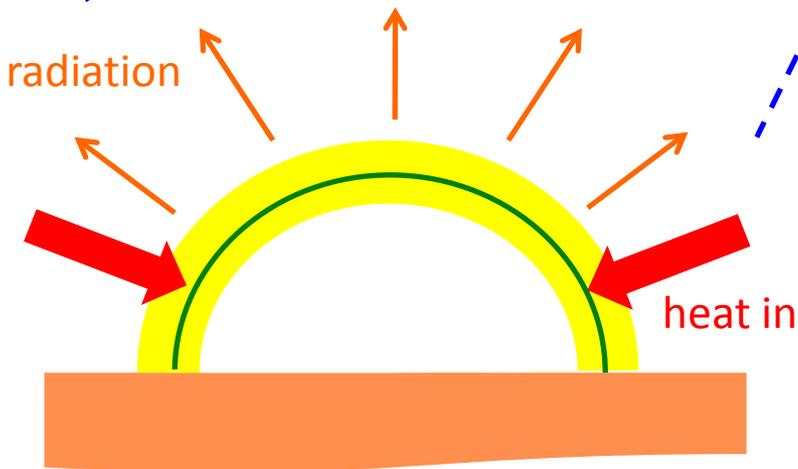
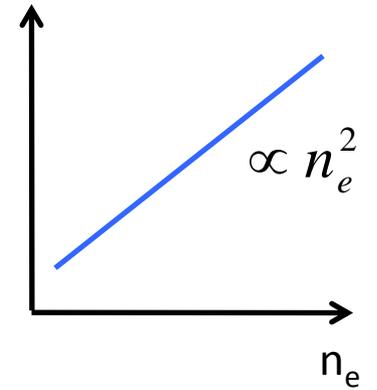
- what you see is all the same T ( $1.5 \times 10^6$  K)
- bright = dense plasma –  $n_e^2$
- heating **can**\* make plasma dense & thus bright
- heating is evidently magnetic

\* if magnetic field lines are closed – magnetic bottle

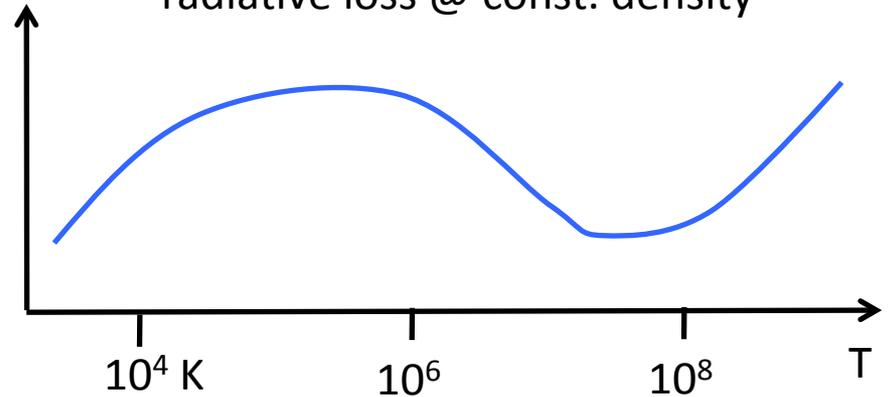
**B** large enough to restrict plasma motion: only along field lines



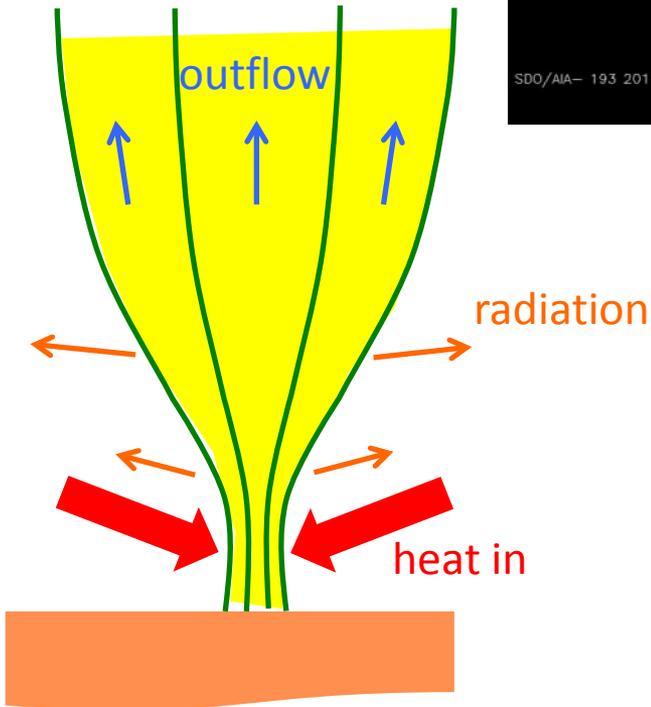
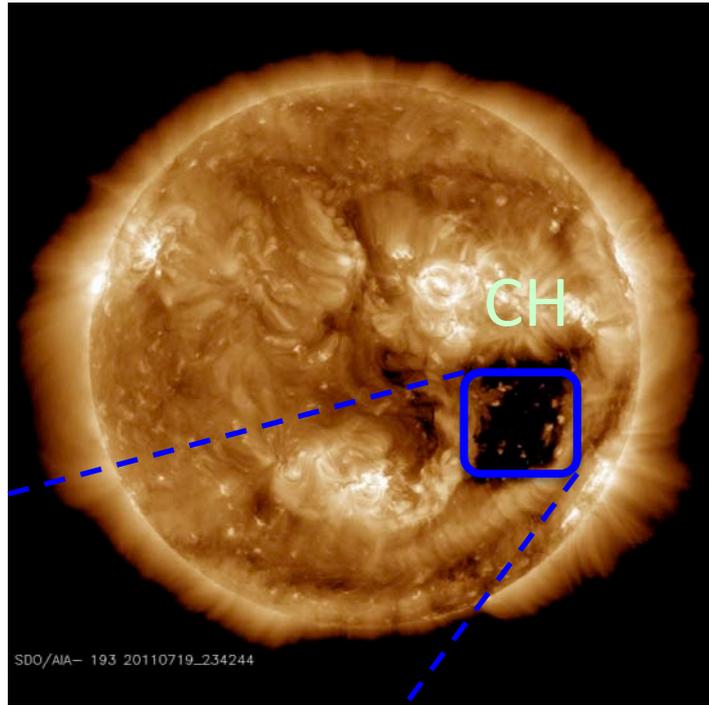
radiative loss @ const. temp.



radiative loss @ const. density



**B** large enough to restrict plasma motion: only along field lines



specific enthalpy

$$w(\rho) \propto \frac{\gamma}{\gamma - 1} \rho^{\gamma - 1}$$

Advective energy loss –

$$\frac{1}{2} \rho \mathbf{v} v^2 + \rho \mathbf{v} w(\rho)$$

>> radiative loss

Bernoulli's law:  $\frac{Q}{M} = \text{const.}$

Energy loss =  $A\rho v \left[ \frac{1}{2}v^2 + w(\rho) + \Psi(s) \right] = Q = \text{fixed \& given}$

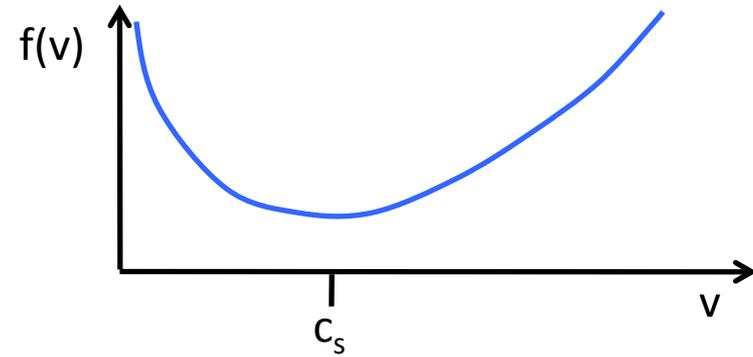
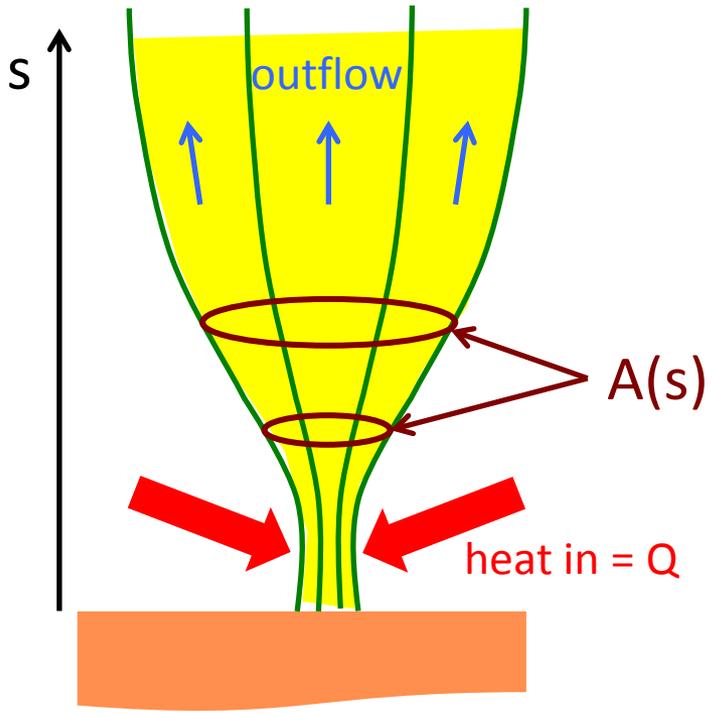
mass loss fixed & unknown

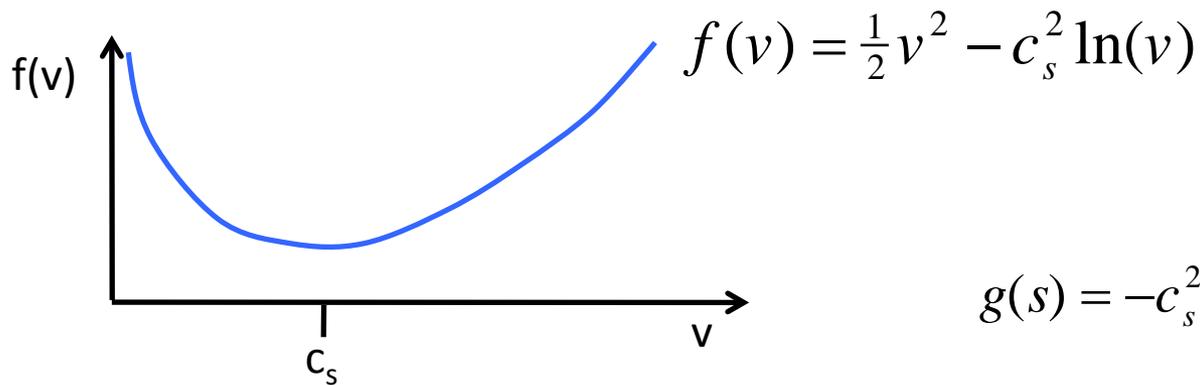
Simple case: Isothermal ...  $\gamma \rightarrow 1$

$w(\rho) \propto \frac{\gamma}{\gamma-1} \rho^{\gamma-1} \rightarrow c_s^2 \ln(\rho) + \text{const.}$

$\rightarrow \frac{1}{2}v^2 - c_s^2 \ln(v) - c_s^2 \ln[A(s)] + \Psi(s) = \text{const.}$

$= f(v) + g(s) = \text{const.}$





$$g(s) = -c_s^2 \ln[A(s)] - \frac{R_o v_{\text{esc}}^2}{2r(s)}$$

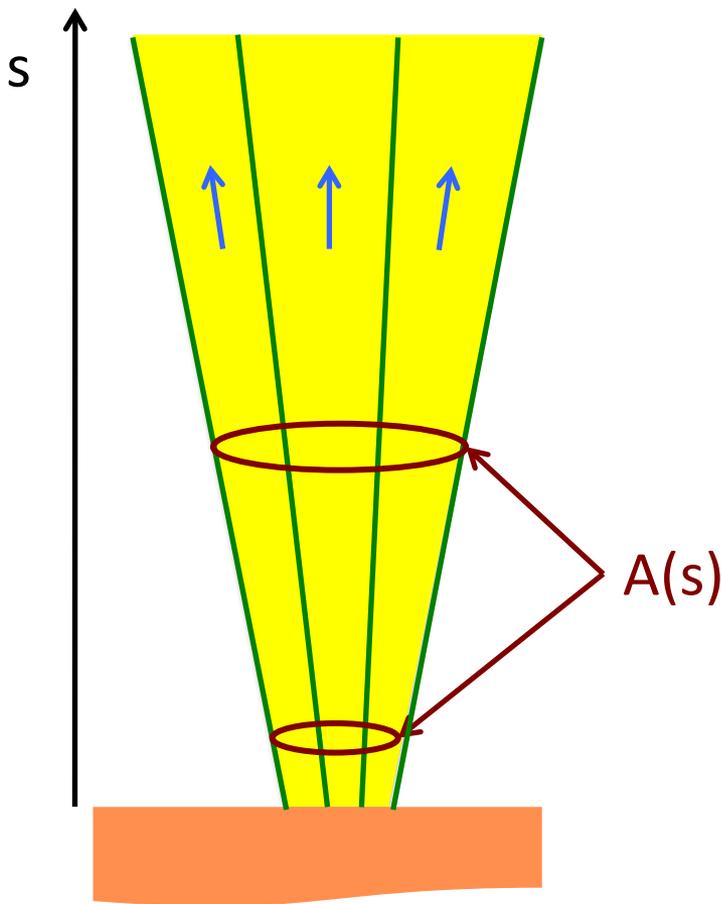
tube:

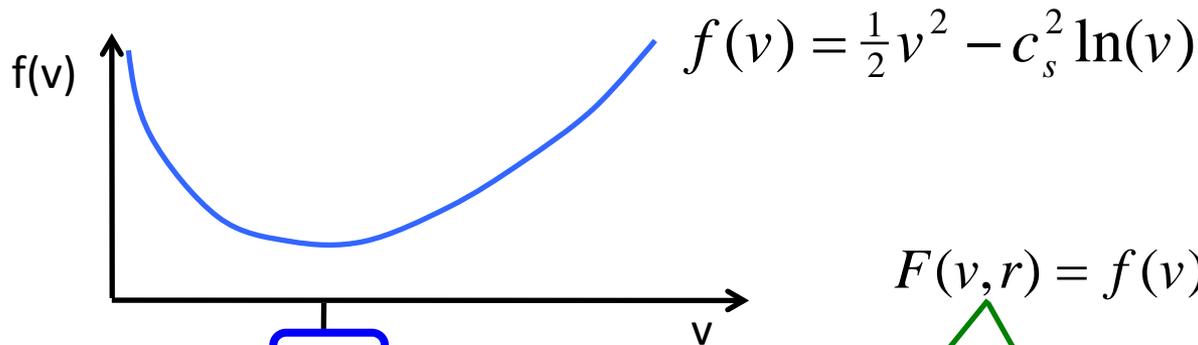
cone w/ vertical axis

$$A(s) \sim s^2$$

$$s = r$$

$$g(r) = -2c_s^2 \ln(r) - \frac{R_o v_{\text{esc}}^2}{2r}$$





$c_s$

tube:

cone w/ vertical axis

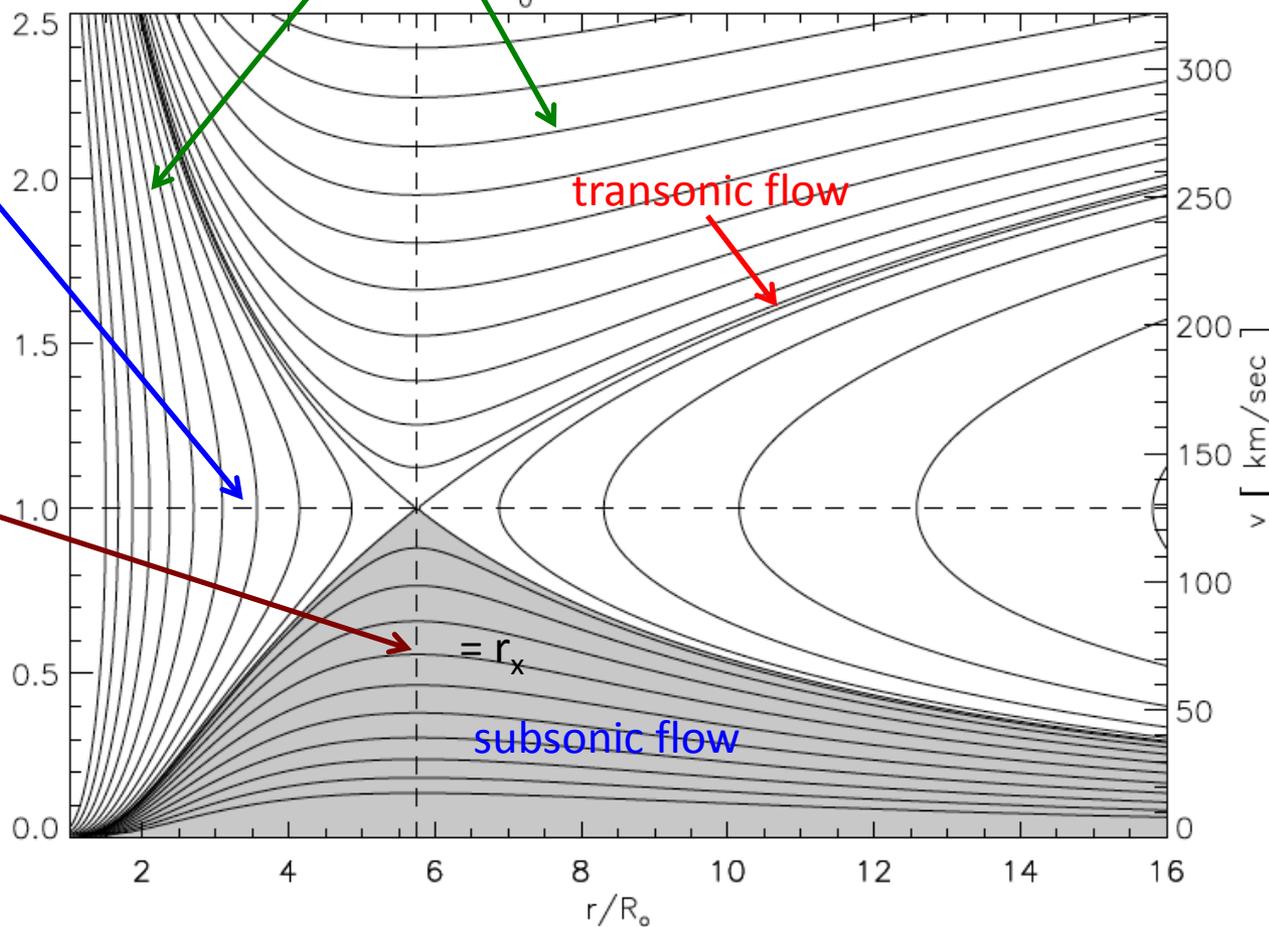
$A(s) \sim s^2$

$s = r$

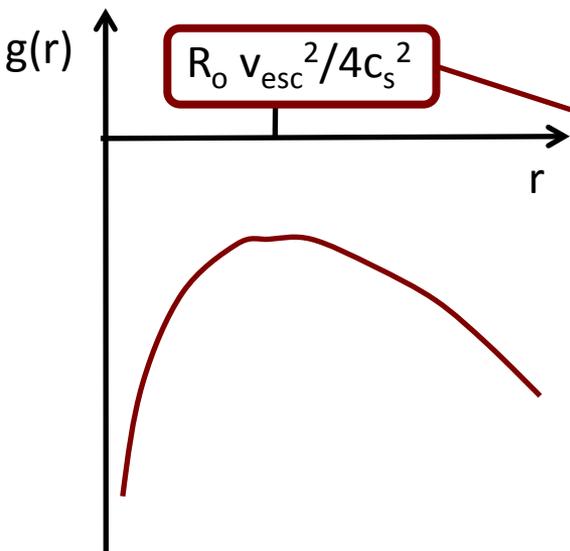
$$g(r) = -2c_s^2 \ln(r) - \frac{R_o v_{esc}^2}{2r}$$

$$F(v, r) = f(v) + g(r) = \frac{Q}{M} = \text{const.}$$

$T_0 = 1.0 \text{ MK}$



$R_o v_{esc}^2 / 4c_s^2$



tube:

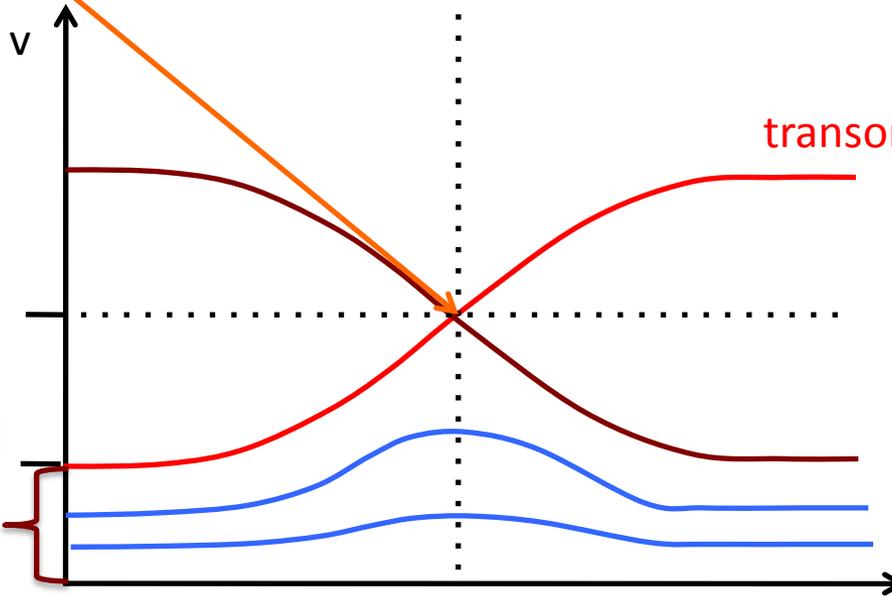
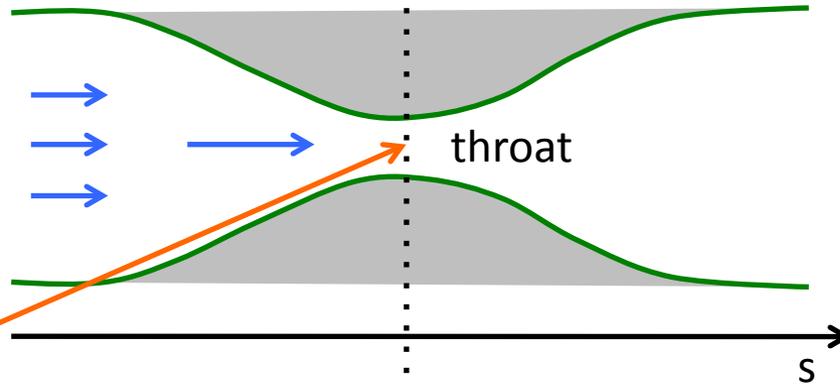
$$g(s) = -c_s^2 \ln[A(s)] + \Psi(s)$$

horizontal nozzle

$$\Psi(s) = \text{const.}$$

$$g(s) = -c_s^2 \ln[A(s)]$$

saddle @ max.  $g(s)$   
@ throat of nozzle



transonic flow

max. inflow speed

admissible inflow speeds

subsonic flow

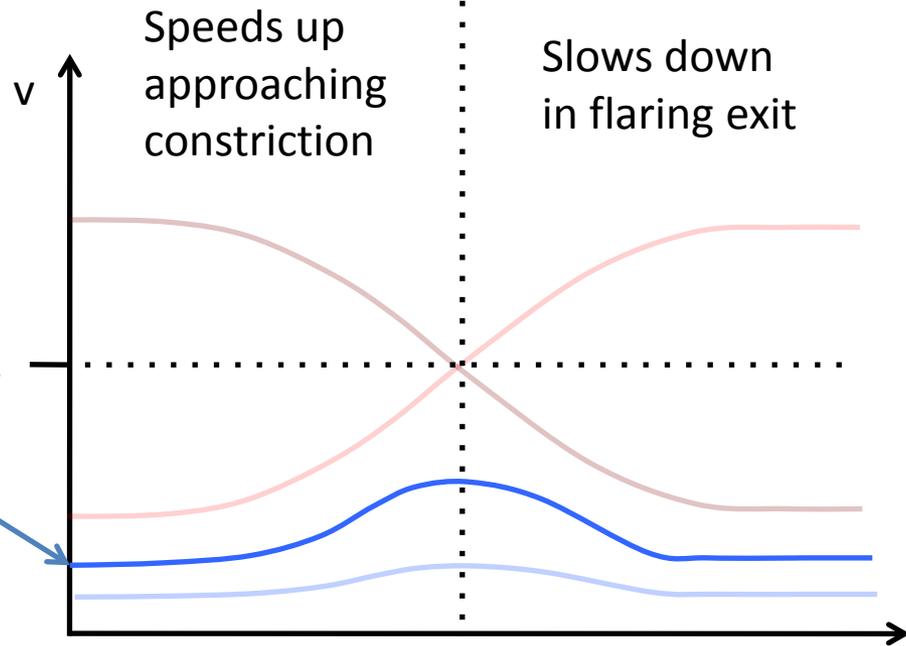
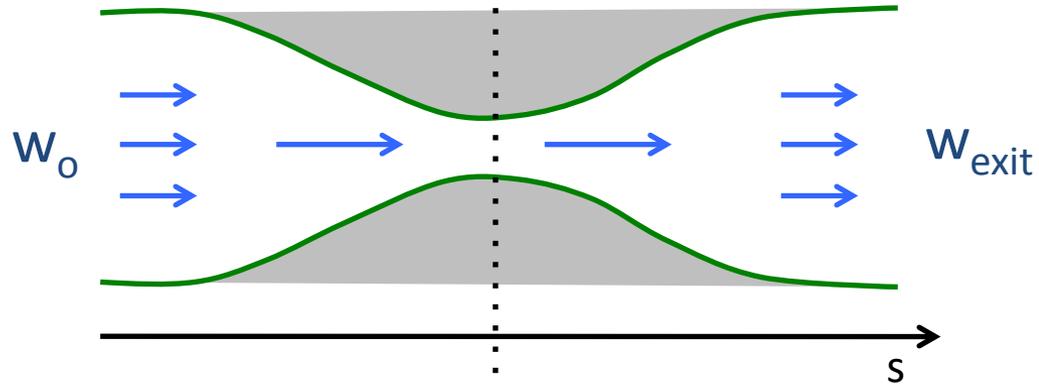
tube:

$$g(s) = -c_s^2 \ln[A(s)] + \Psi(s)$$

horizontal nozzle

$$\Psi(s) = \text{const.}$$

$$g(s) = -c_s^2 \ln[A(s)]$$



Inflow = mass loss rate

set by  
back-pressure

$W_{\text{exit}}$

subsonic flow

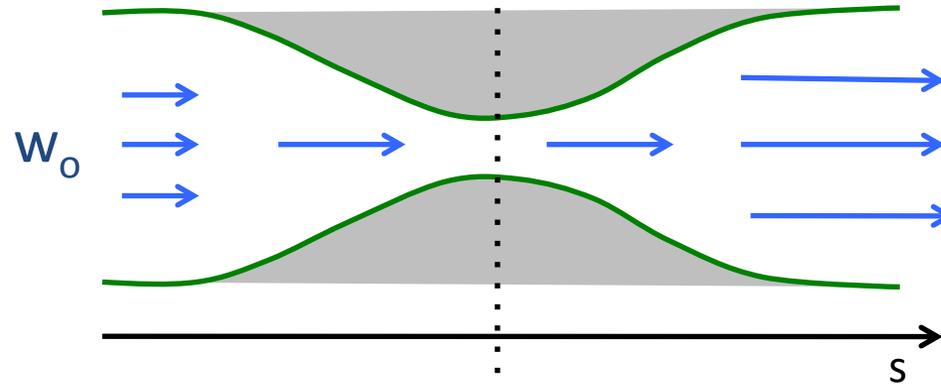
tube:

horizontal nozzle

$$\Psi(s) = \text{const.}$$

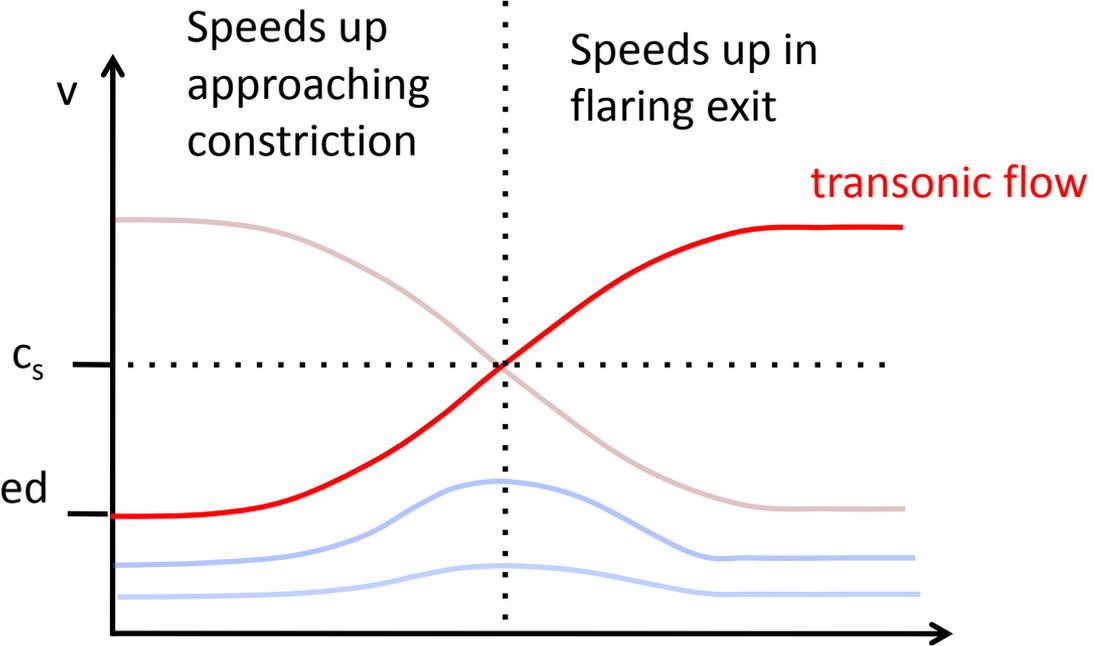
$$g(s) = -c_s^2 \ln[A(s)]$$

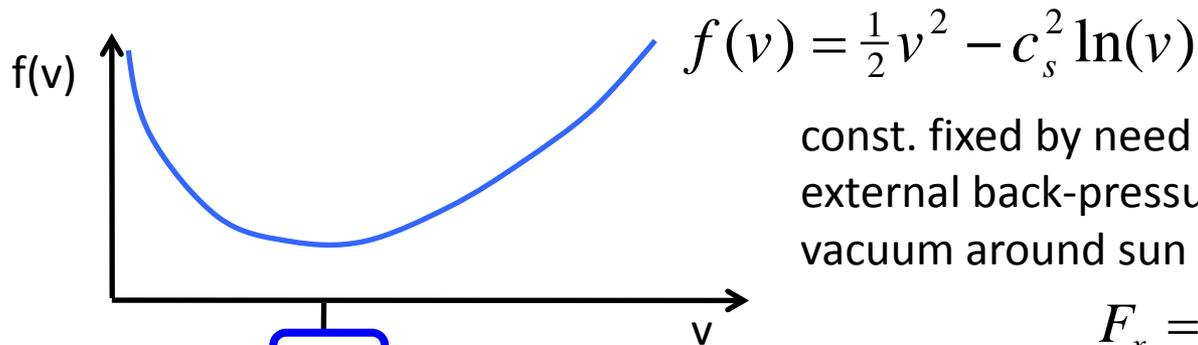
$$g(s) = -c_s^2 \ln[A(s)] + \Psi(s)$$



occurs for  
back-pressure  
insufficient to  
keep flow  
sub-sonic

max. inflow speed

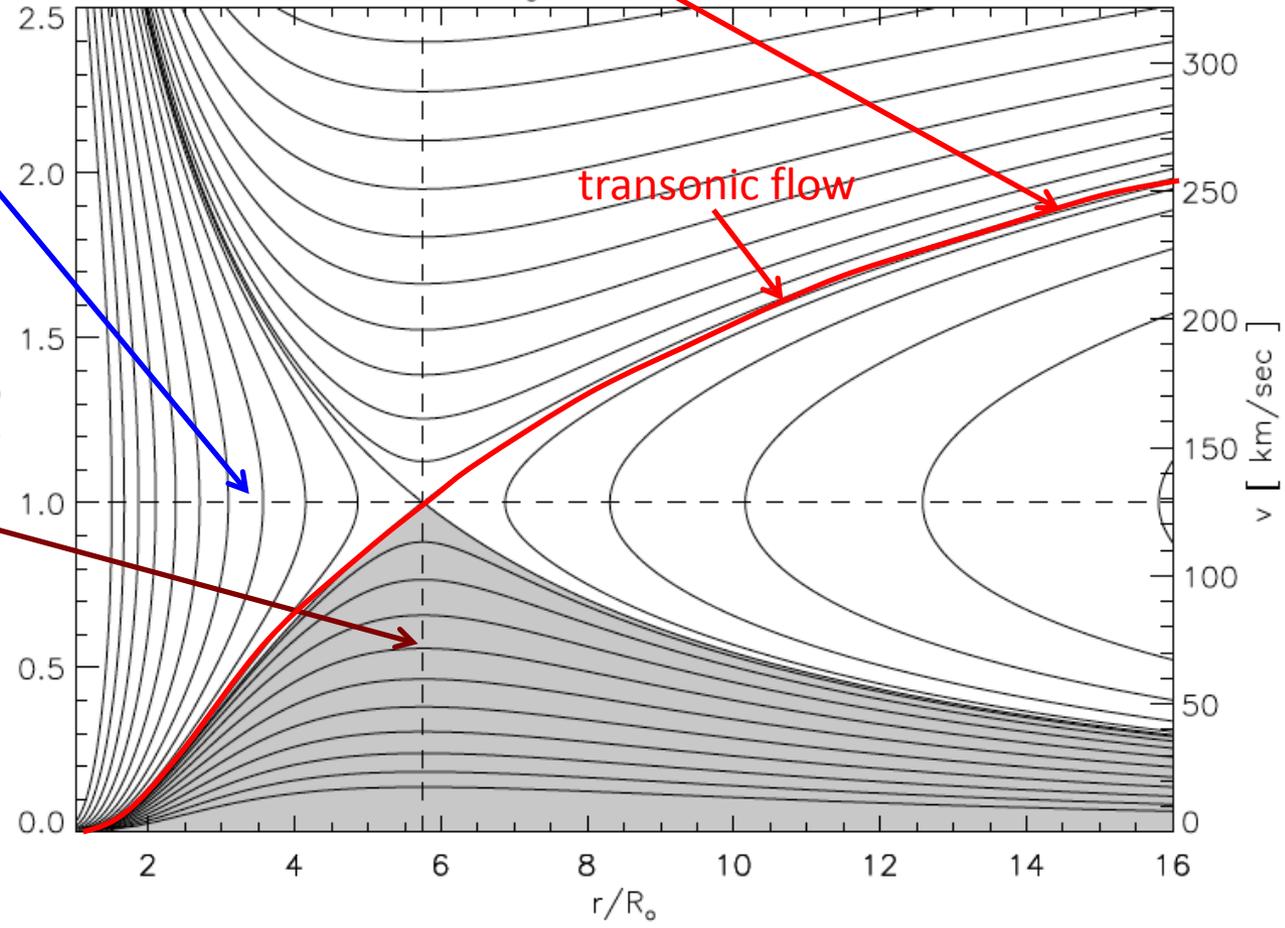
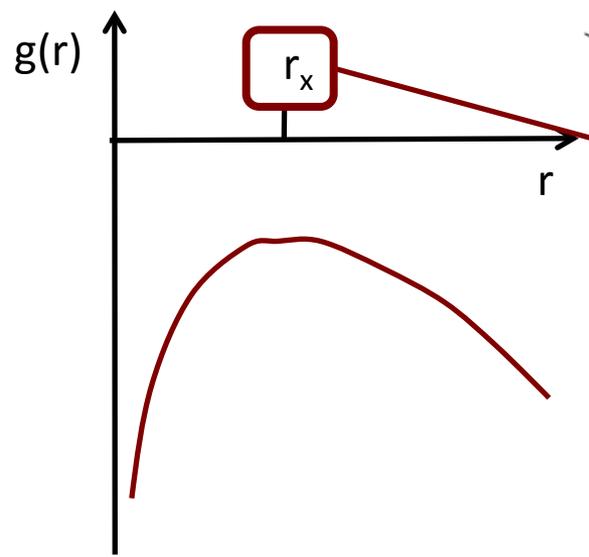




$$F_x = f(c_s) + g(r_x) = \frac{Q}{\dot{M}}$$

$T_0 = 1.0 \text{ MK}$

$$g(r) = -2c_s^2 \ln(r) - \frac{R_o v_{\text{esc}}^2}{2r}$$



→ Mass loss rate is set by heating rate\*

$$\dot{M} = \frac{Q}{F_x}$$

→ density everywhere is set by mass loss rate

$$\rho(r_x) = \frac{\dot{M}}{A(r_x)c_s}$$

→ density @ base is set by heating rate\*...

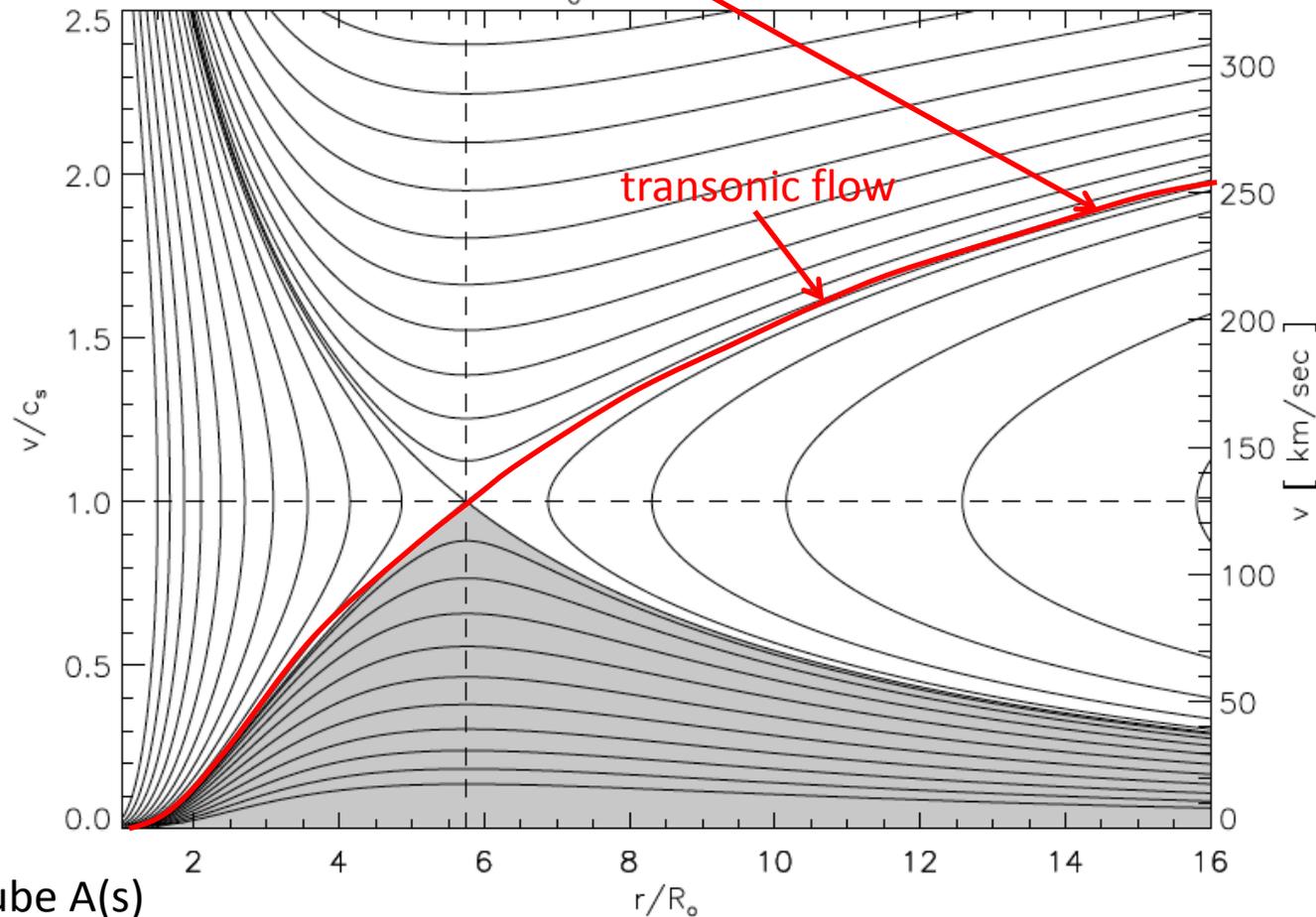
... and it will be lower than density on closed loops w/ same heating (Why?)

\* ... and geometry of flux tube  $A(s)$

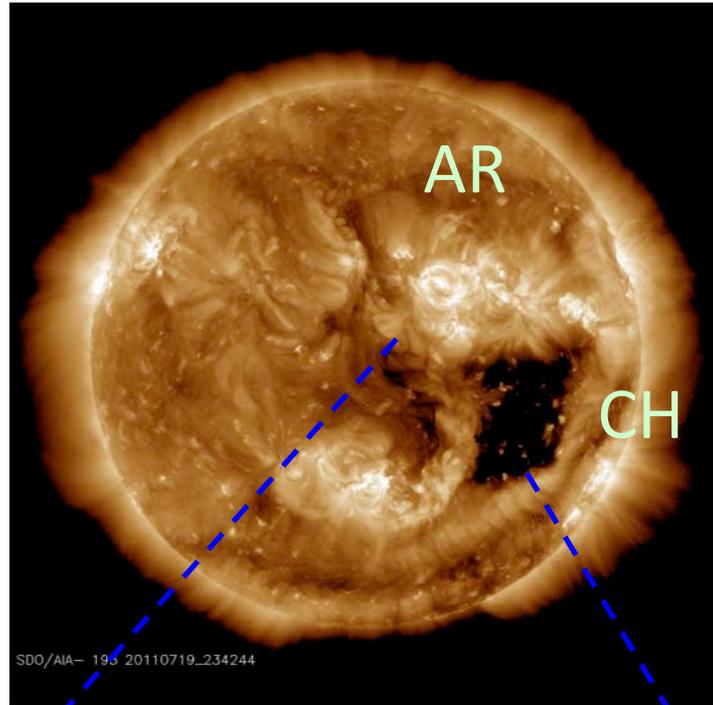
const. fixed by need to become transonic when external back-pressure is insufficient – i.e. vacuum around sun

$$F_x = f(c_s) + g(r_x) = \frac{Q}{\dot{M}}$$

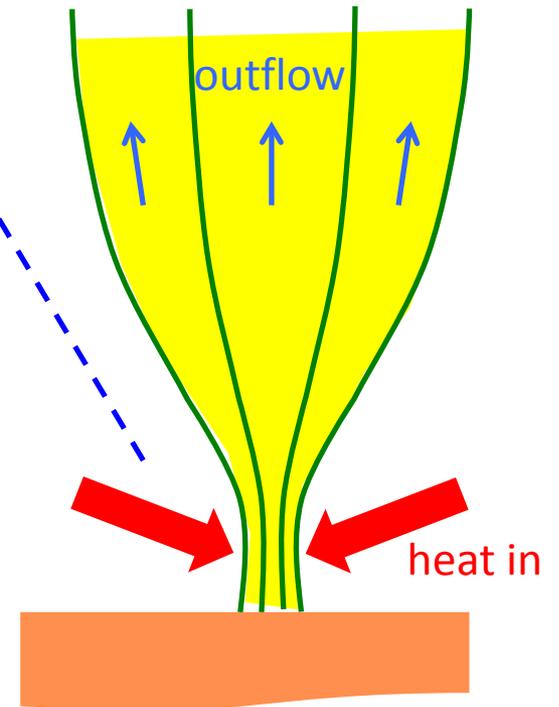
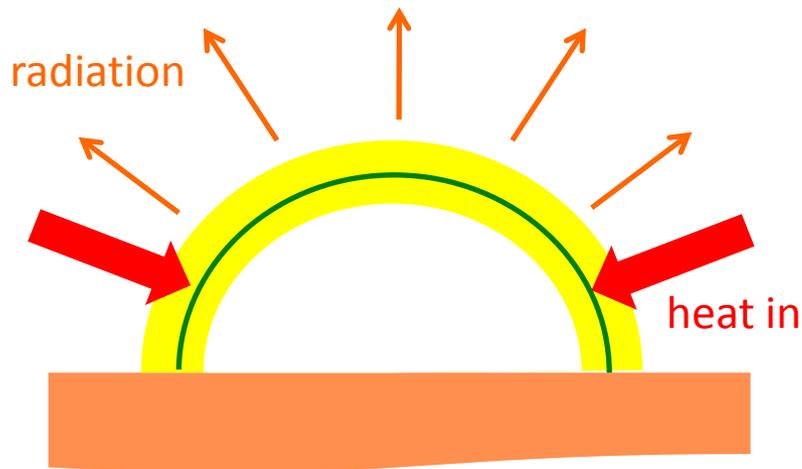
$T_0 = 1.0 \text{ MK}$



**B** large enough to restrict plasma motion: only along field lines



Different coronae from different magnetic topology: open vs. closed



# Why are some field lines open & others closed?

Magnetic field dominates:  
nothing capable of countering its force so...

$$(\nabla \times \mathbf{B}) \times \mathbf{B} = 0$$

$$\Rightarrow \nabla \times \mathbf{B} = \alpha \mathbf{B} \quad (\text{i.e. } \parallel \mathbf{B})$$

simplest version:  $\alpha = 0$  (by fiat)

$$\Rightarrow \nabla \times \mathbf{B} = 0 \quad \Rightarrow \mathbf{B} = -\nabla \chi \quad \text{potential field}$$

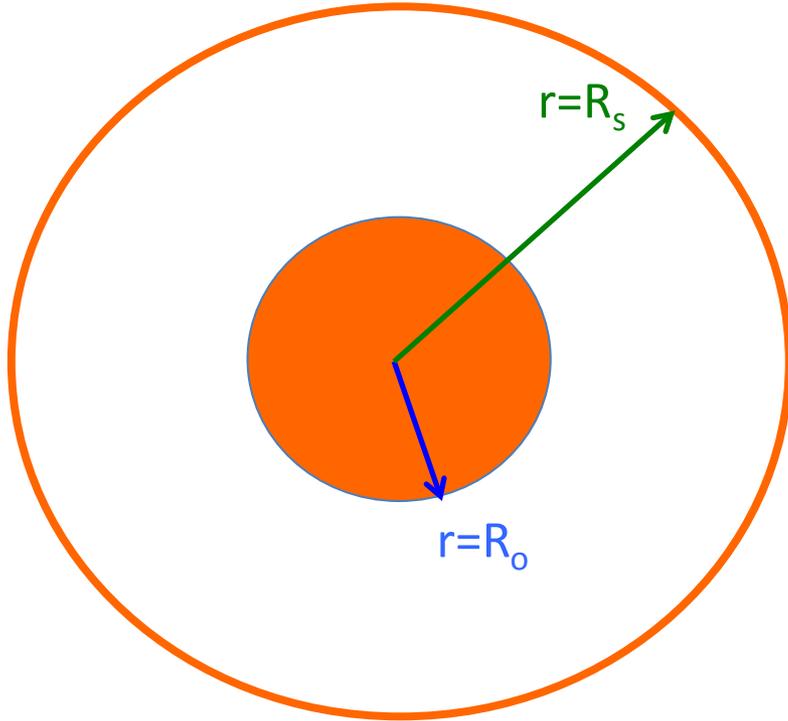
(cf. electrostatics)

$$\nabla \cdot \mathbf{B} = 0 \quad \Rightarrow \quad \nabla^2 \chi = 0 \quad \text{harmonic potential}$$

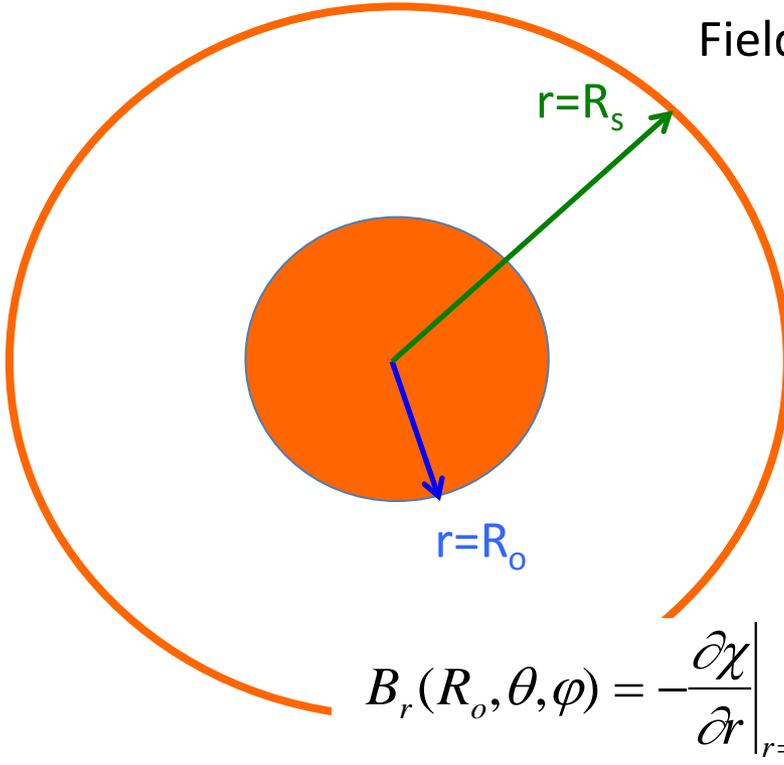
(cf. electrostatics in vacuum)

$$\mathbf{B} = -\nabla\chi \quad \& \quad \nabla^2\chi = 0$$

potential field outside  
sphere  $r=R_0$



$$\mathbf{B} = -\nabla\chi \quad \& \quad \nabla^2\chi = 0 \quad \text{potential field outside sphere } r=R_o$$



Field: purely radial @  $r=R_s$  (by fiat)

$$(B_\theta, B_\phi) = 0 \Rightarrow \left( \frac{\partial\chi}{\partial\theta}, \frac{\partial\chi}{\partial\phi} \right) = 0$$

$$\Rightarrow \chi(R_s, \theta, \phi) = 0 \quad \text{Dirichlet}$$

$$\chi(r, \theta, \phi) = \sum_{l,m} A_{l,m} \left[ \left( \frac{R_s}{r} \right)^{l+1} - \left( \frac{r}{R_s} \right)^l \right] Y_{l,m}(\theta, \phi)$$

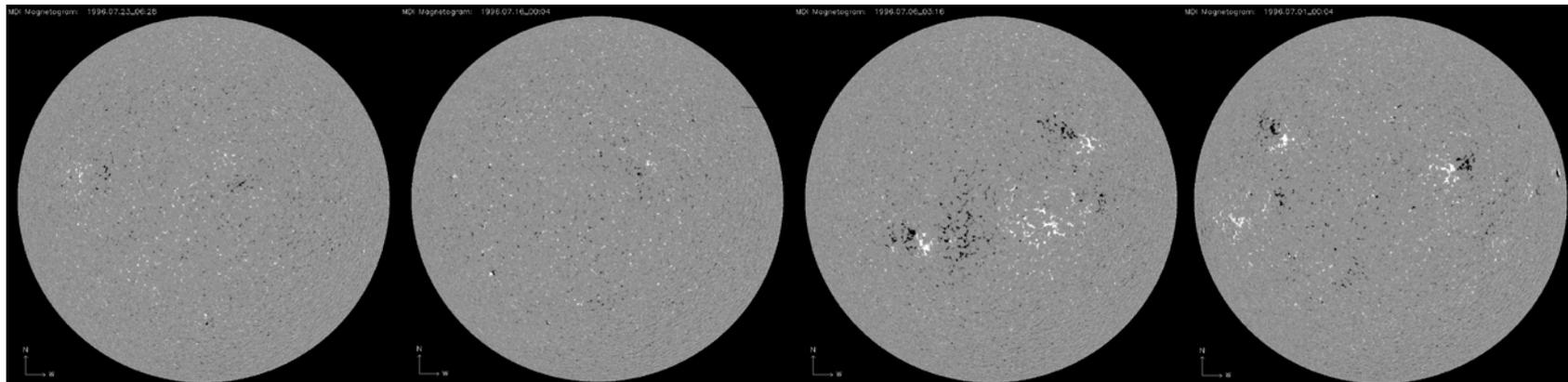
$$B_r(R_o, \theta, \phi) = - \left. \frac{\partial\chi}{\partial r} \right|_{r=R_o}$$

Observed (Neumann)

$$B_r(R_o, \theta, \phi) = \sum_{l,m} \frac{A_{l,m}}{R_s} \left[ (l+1) \left( \frac{R_s}{R_o} \right)^{l+2} + l \left( \frac{R_o}{R_s} \right)^{l-1} \right] Y_{l,m}(\theta, \phi)$$

- Observe  $B_r(\theta, \phi)$   
@ photosphere
- decompose w/ spherical harmonics
- coeffs.  $\rightarrow A_{l,m}$

← time

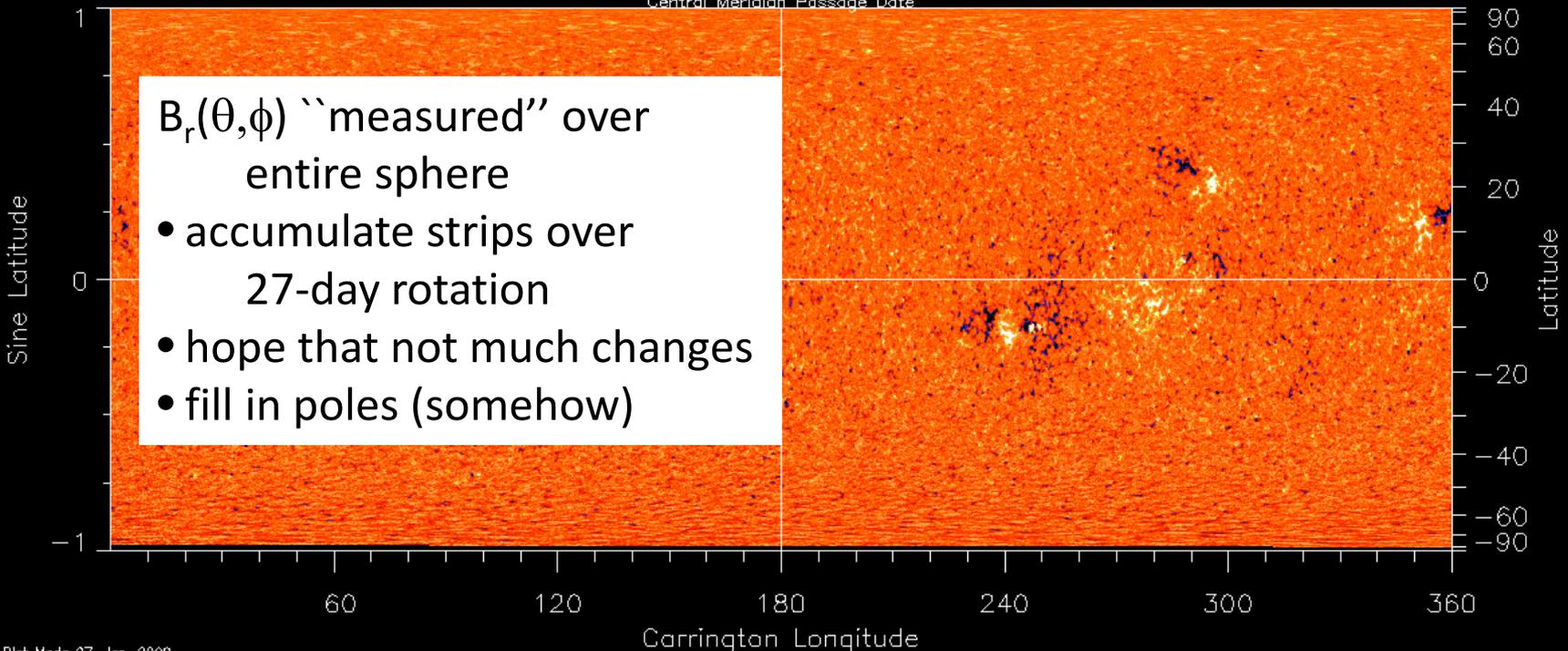


MDI Synoptic Chart for Carrington Rotation 1911



$B_r(\theta, \phi)$  "measured" over entire sphere

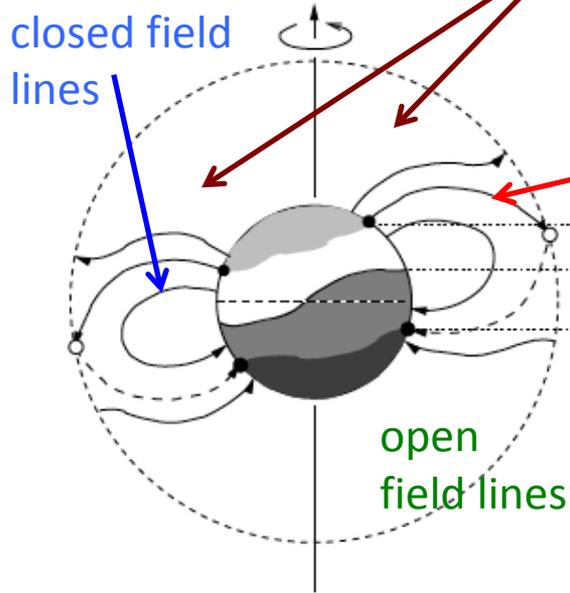
- accumulate strips over 27-day rotation
- hope that not much changes
- fill in poles (somehow)



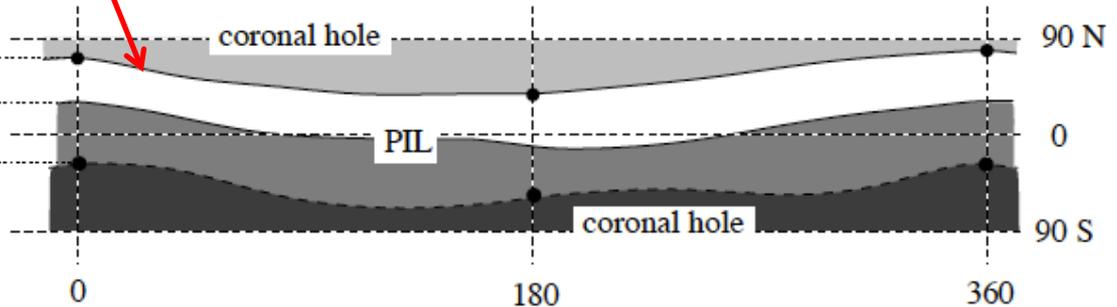
$$\chi(r, \theta, \varphi) = \sum_{l,m} A_{l,m} \left[ \left( \frac{R_s}{r} \right)^{l+1} - \left( \frac{r}{R_s} \right)^l \right] Y_{l,m}(\theta, \varphi)$$

# PFSS model

(potential field source surface)

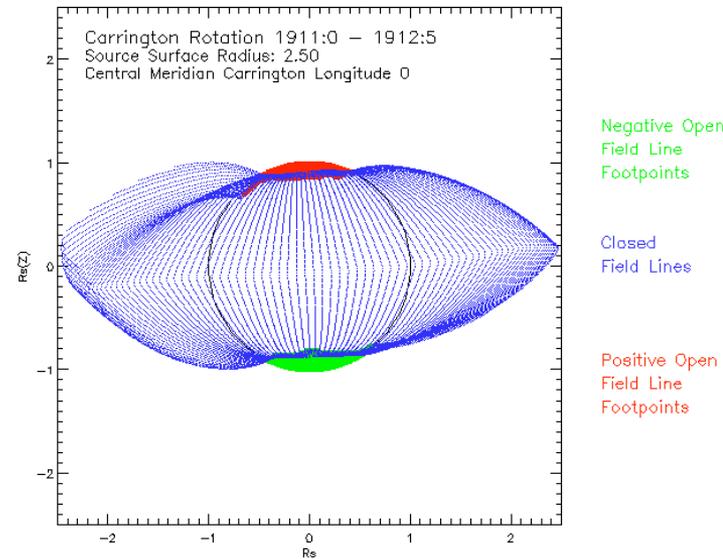
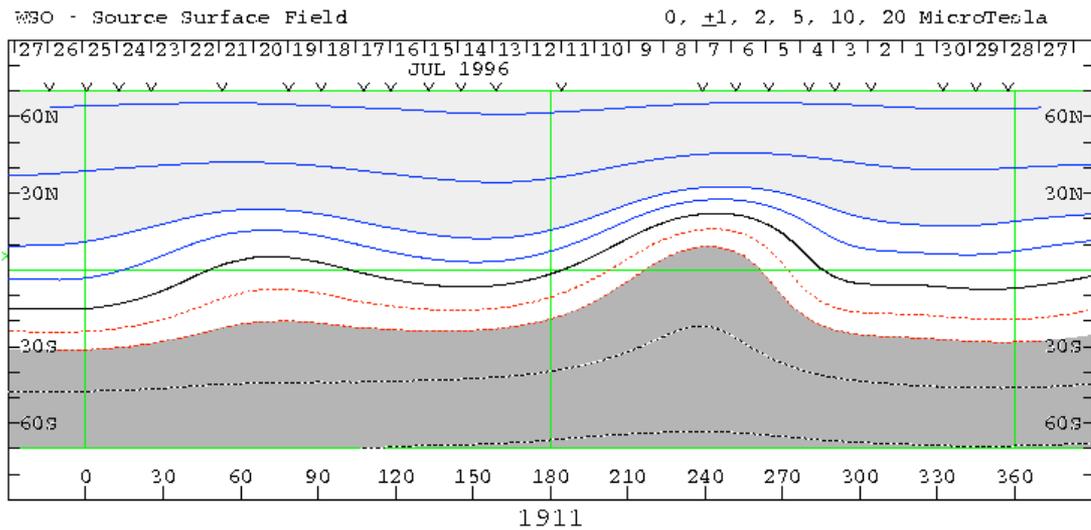


Separatrix dividing open from closed

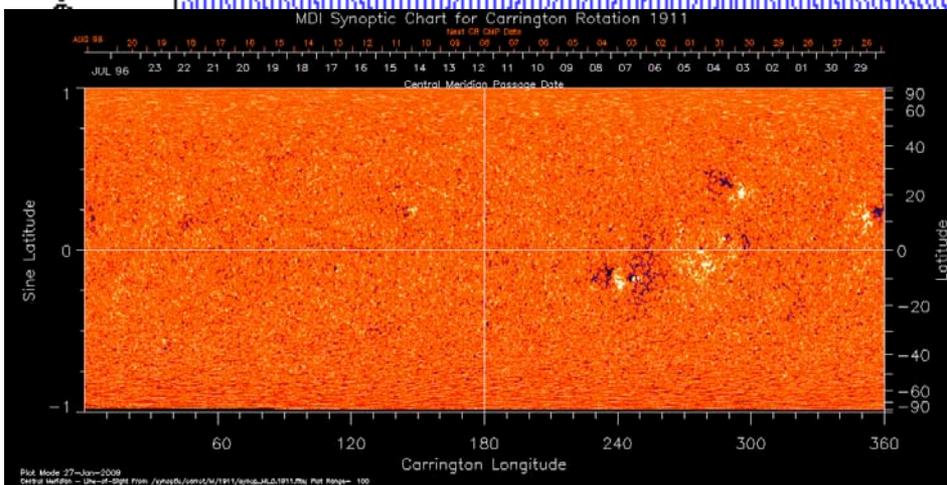
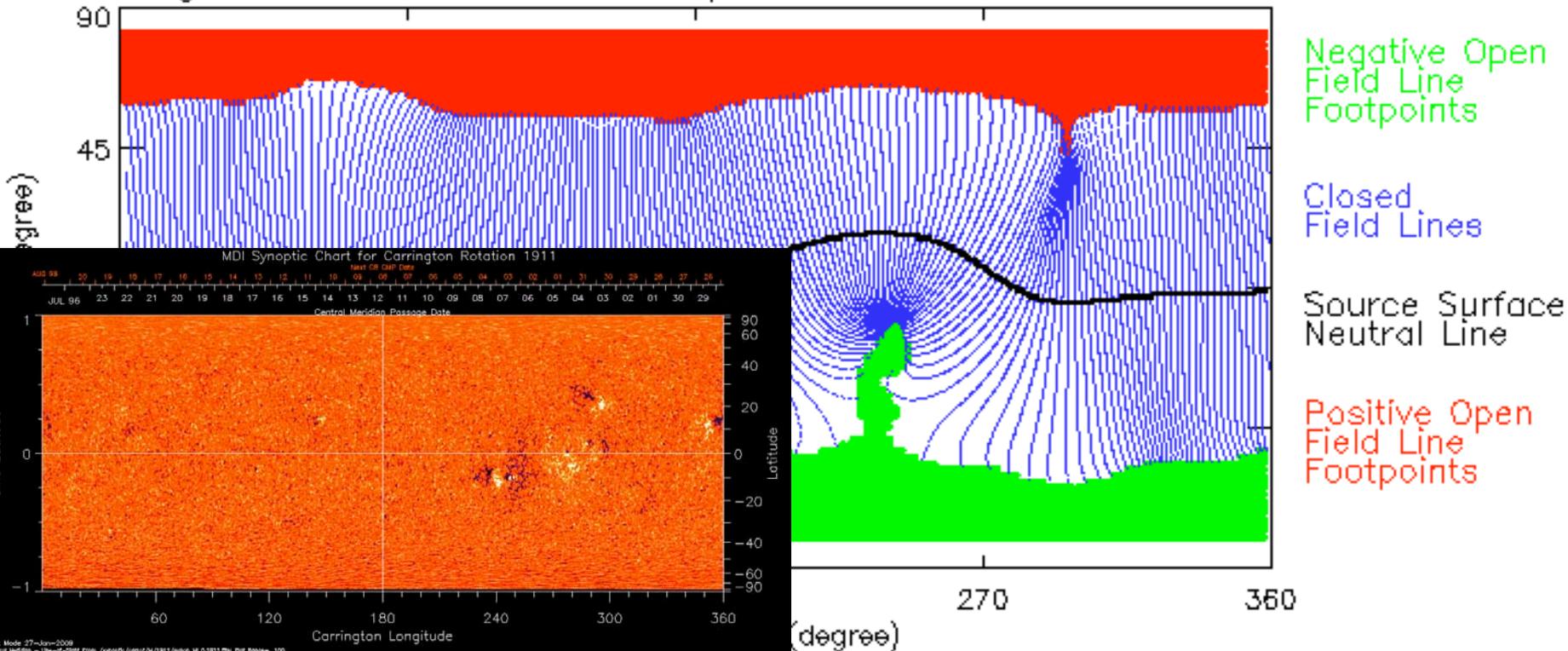


Solar wind flows from open field crossing  $r=R_s$  ... the 'source' of the wind → the 'source surface'

- $B_r(\theta, \phi)$  "measured" over entire sphere
- accumulate strips over 27-day rotation
  - hope that not much changes
  - fill in poles (somehow)
  - decompose w/ spherical harmonics
  - coeffs. →  $A_{l,m}$

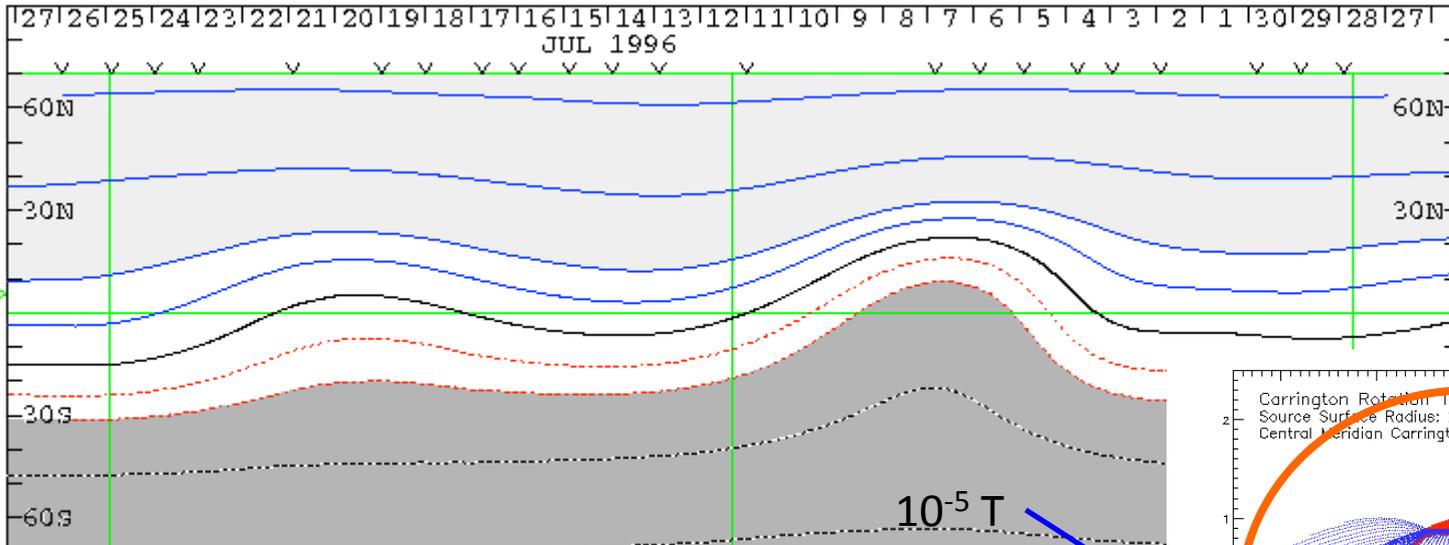


Carrington Rotation 1911:0-1912:5 / Source Surface Radius: 2.50



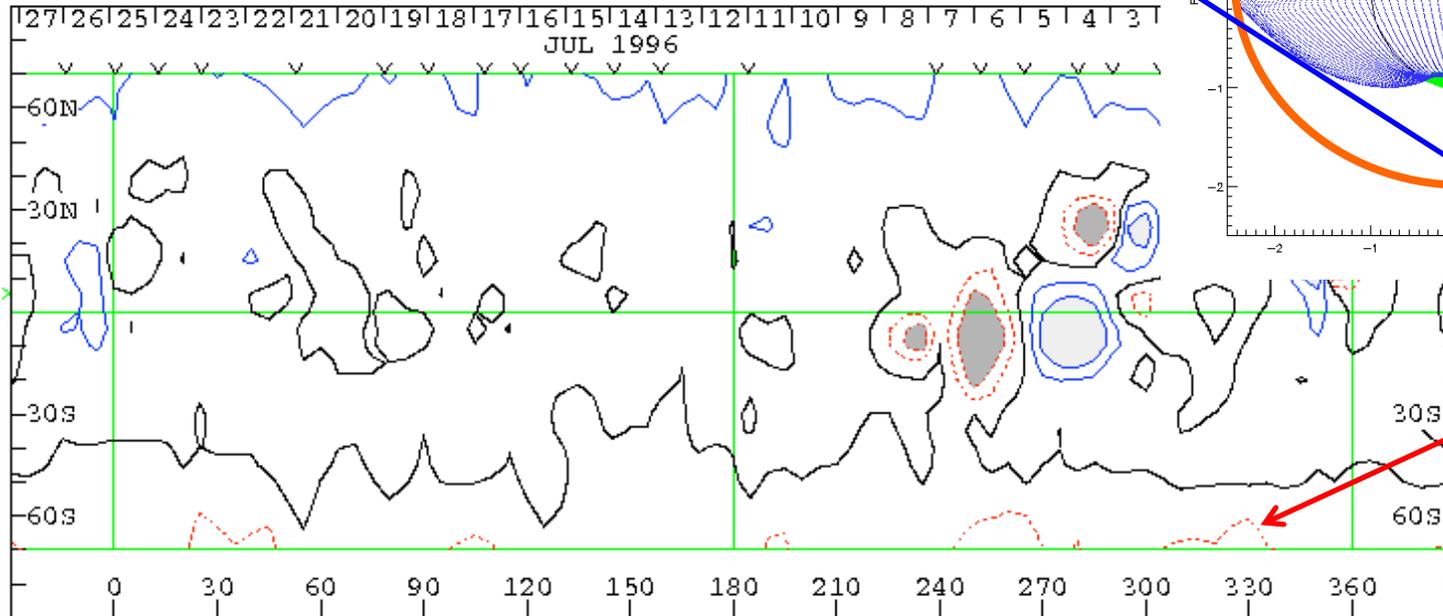
WSO - Source Surface Field

0, +1, 2, 5, 10, 20 MicroTesla

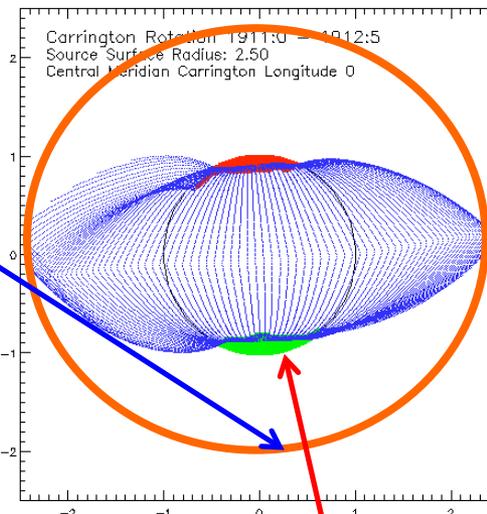


WSO - Photospheric Magnetic Field

0, +100, 200, 500, 1000, 2000



1911



Negative Open Field Line Footpoints

Closed Field Lines

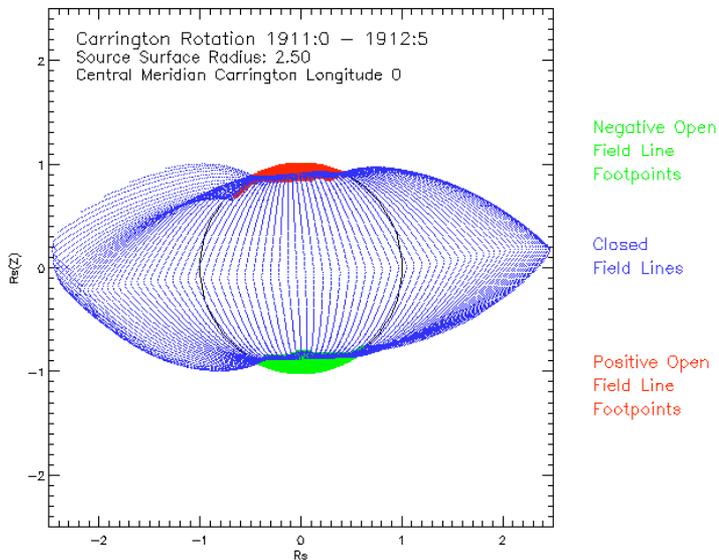
Positive Open Field Line Footpoints

# Assumptions of the PFSS

- No currents in coronal field (simplest equilibrium)

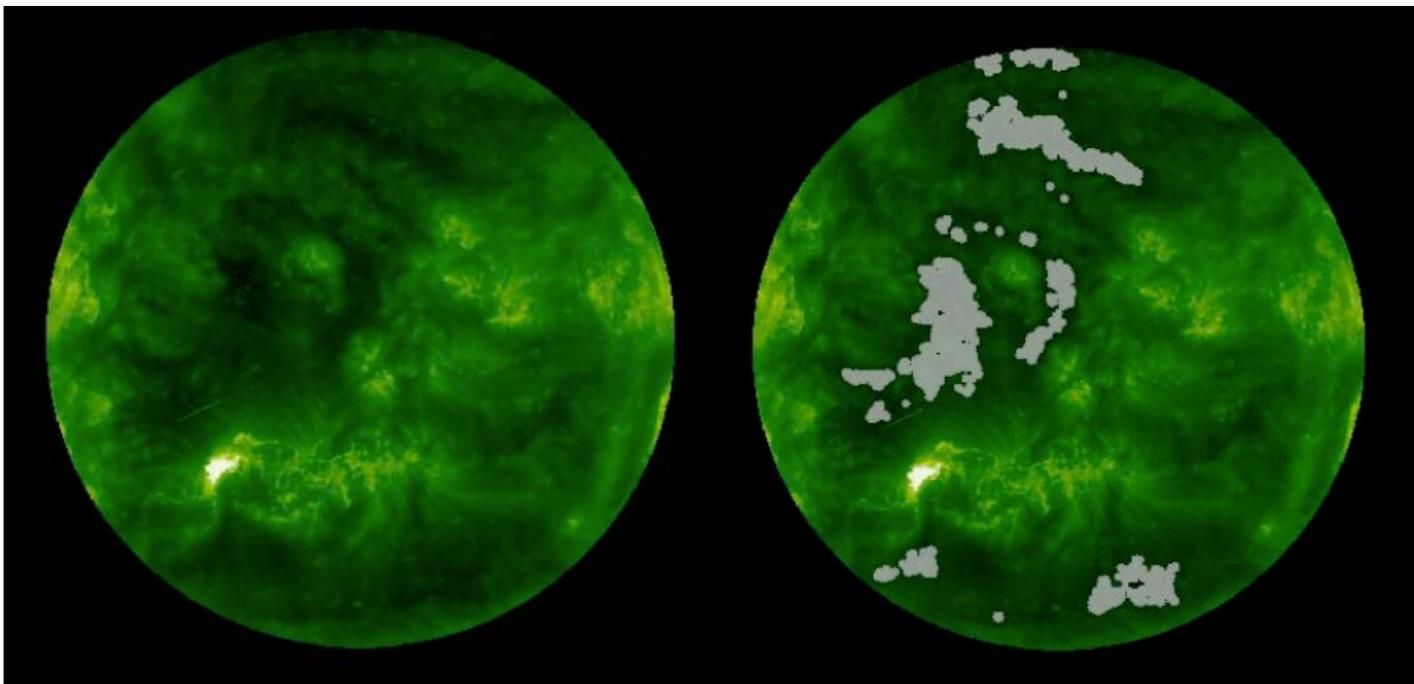
$$\nabla \times \mathbf{B} = 0 \quad R_o < r < R_s$$

- Field becomes open (radial) @ fixed radius  $r=R_s$
- Not much change during 27-day accumulation



➔ **Model** distinguishing open/closed coronal field

➔ Field **actually** open will be source of solar wind, less dense & dark in EUV & SXR

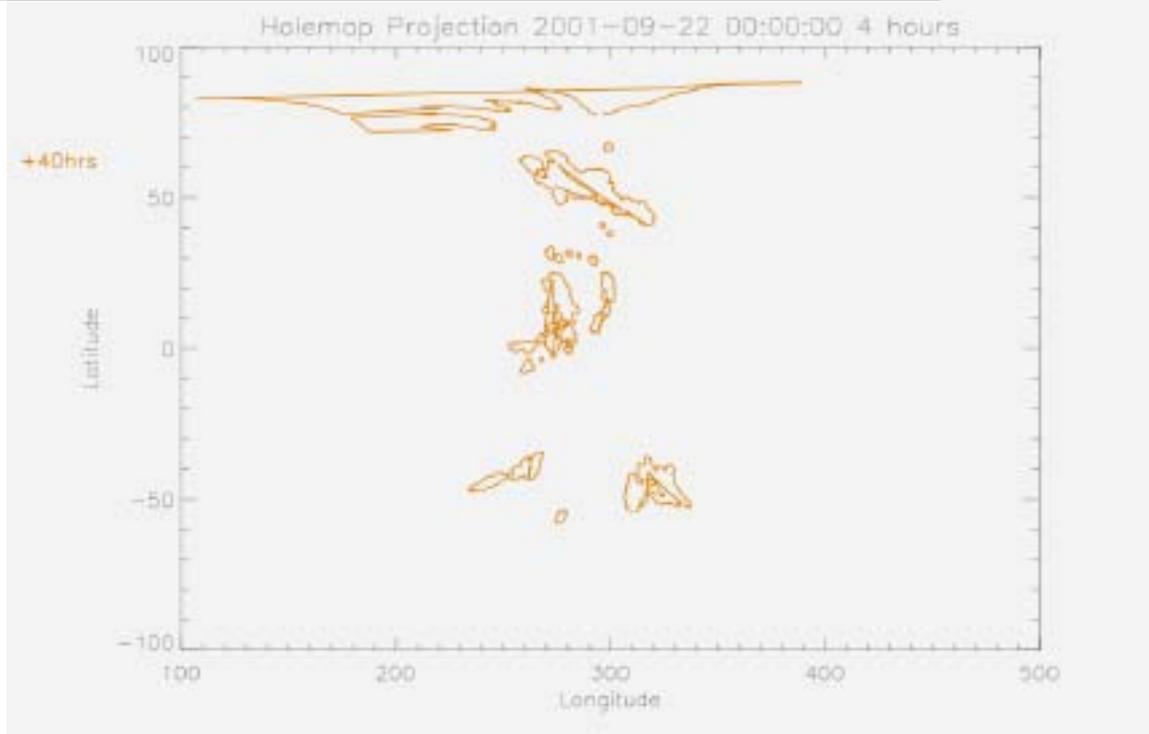


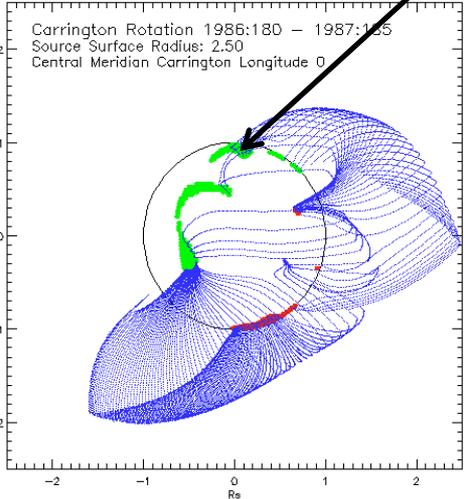
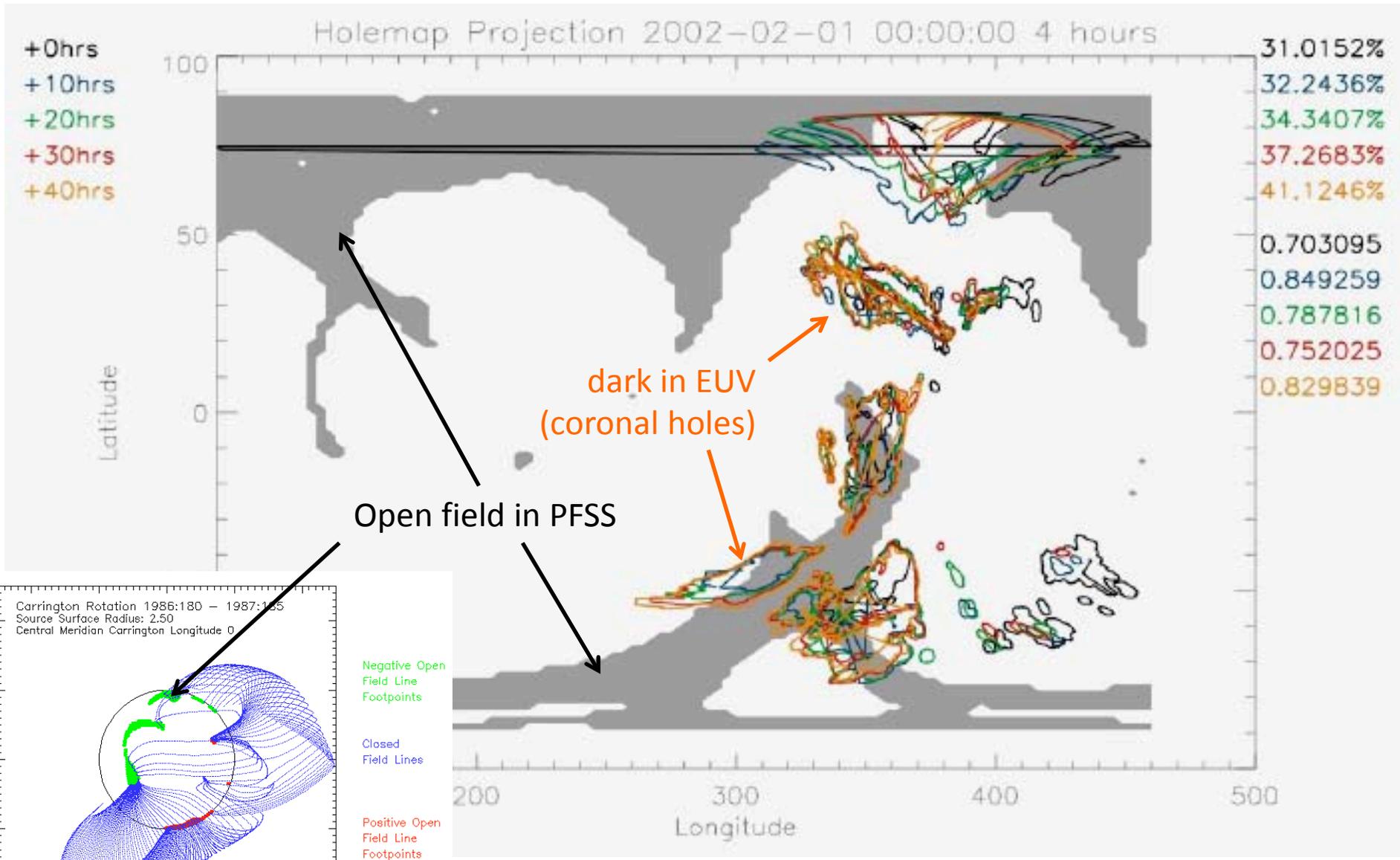
## finding coronal holes

Dustin Hickey

Chris Lowder

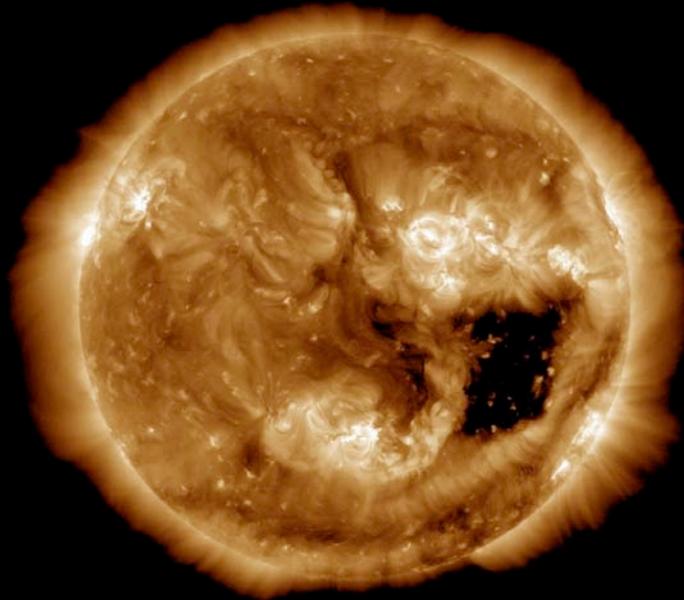
Jiong Qiu & DWL



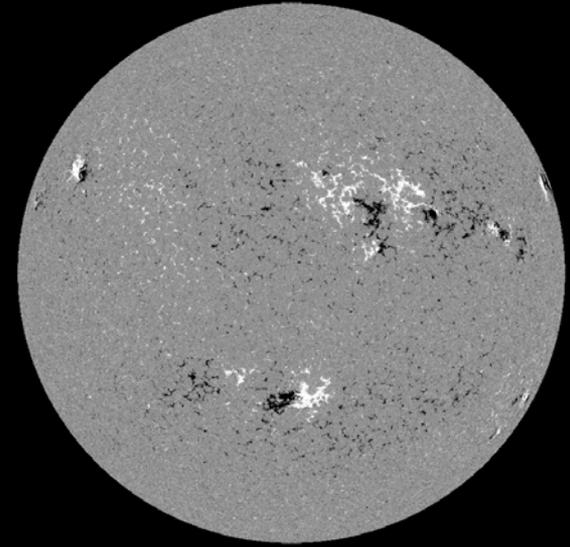




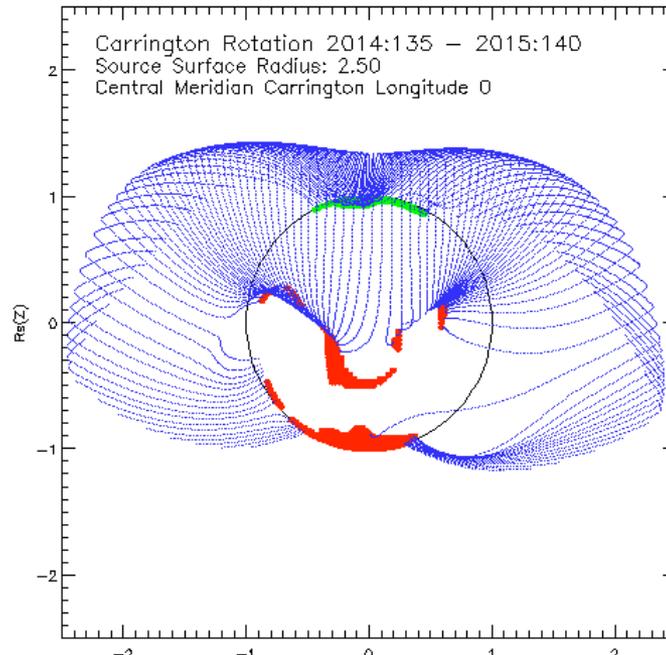
AIA-4500 20110719\_230008



SDO/AIA- 193 20110719\_234244



I 2011-07-19T22:03:41.500



Negative Open  
Field Line  
Footpoints

Closed  
Field Lines

Positive Open  
Field Line  
Footpoints

# Summary

- Heliosphere is a system of (mostly) plasma coupling Sun, & planets
- Includes sources of plasma, magnetic field & heat
- Corona is a source of heat & solar wind
- Energy dissipation drives flow along open field lines: the solar wind
- Coronal field composed of closed & open field according to conditions of magnetic equilibrium