

Radiation Belts

Yuri Shprits^{1,2}

¹Institute of Geophysics and Planetary Physics, UCLA

²Department of Earth and Planetary Sciences, UCLA

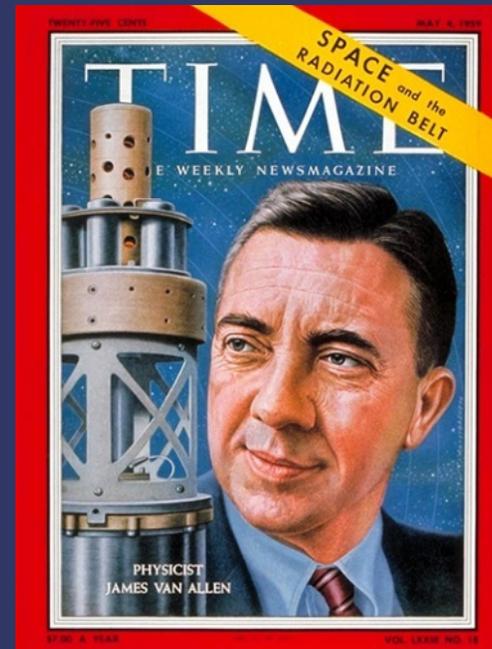
Acknowledgements: Alexander Ukhorskiy, Dmitri Subbotin,
Ksenia Orlova, Marianne Daae



Talk outline

1. Introduction: Discovery of the Radiation Belts.
Kinematics of radiation belt electrons and the quiet time structure of the radiation belts
2. Radiation acceleration mechanisms
3. Radiation belt particle losses
4. 3D simulations of the Radiation Belts using Versatile Electron Radiation Belt code

First discovery of space age



January 31, 1958: The US Explorer 1 satellite is launched into orbit with a Geiger counter on board. James Van Allen discovered that **space is radioactive**

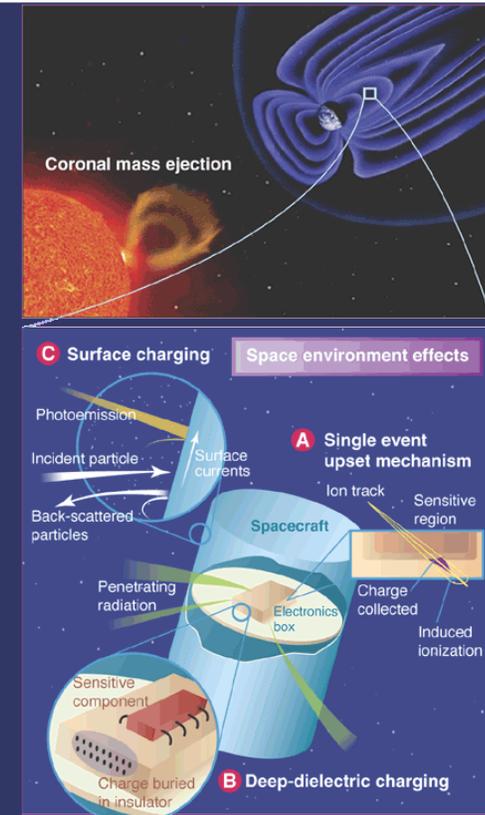
~250 satellites; Supporting \$25 billion/
year industry ;Replacement cost: \$75
billion;

Space weather–induced effects on an
Earth-orbiting spacecraft:

(A) Single-event upsets (SEUs) due to
energetic ions;

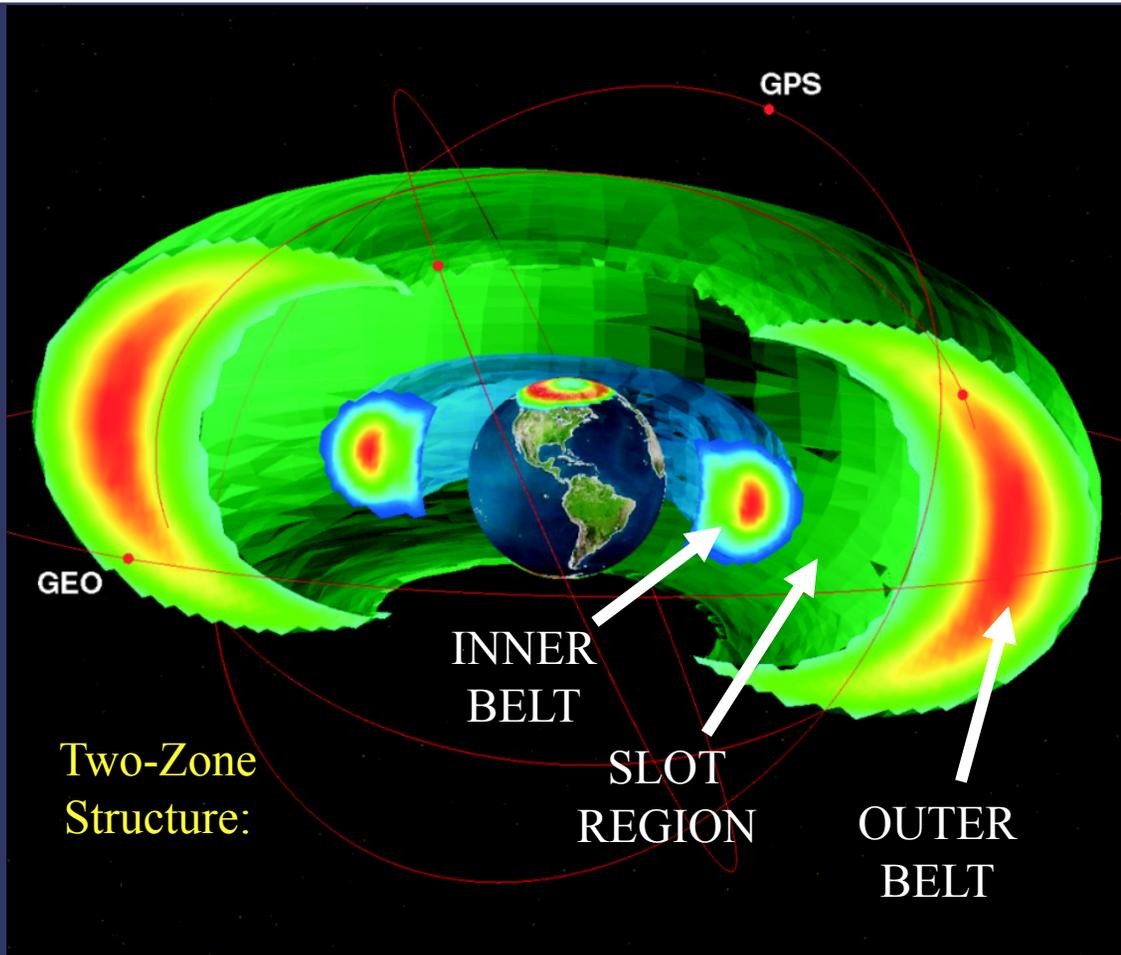
(B) Deep-dielectric charging due to
relativistic electrons;

(C) Surface charging due to moderate-
energy electrons.



Baker et al., 2002;
Odenwald., 2006]

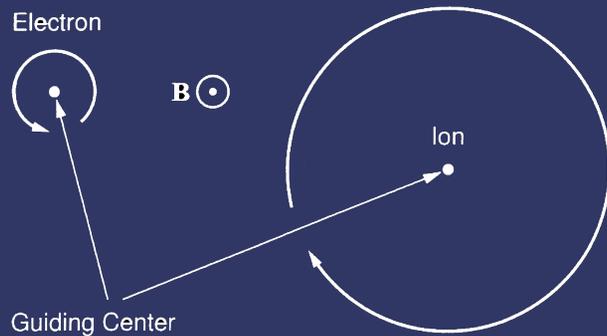
Two-Zone Structure



- Radiation belts
 - energies >100 keV
 - two-zone structure
- Inner belt: fairly stable
- Outer belt: can change on the time scale of an hour.

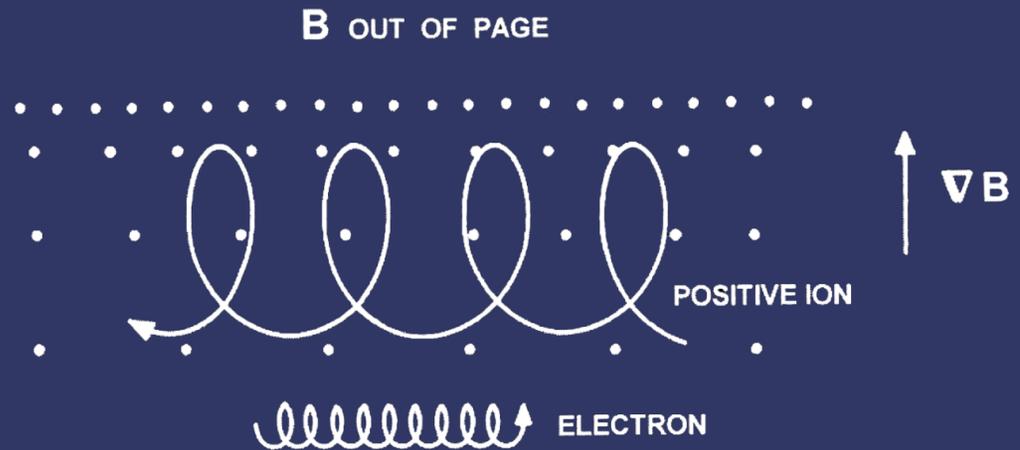
Charged Particle Motion

$E = 0, B = \text{const}$



$E = 0, B \neq \text{const}$

$\rho |\nabla B| \ll B$



Variation of the magnetic field intensity causes guiding center drift perpendicular to the field and its gradient:

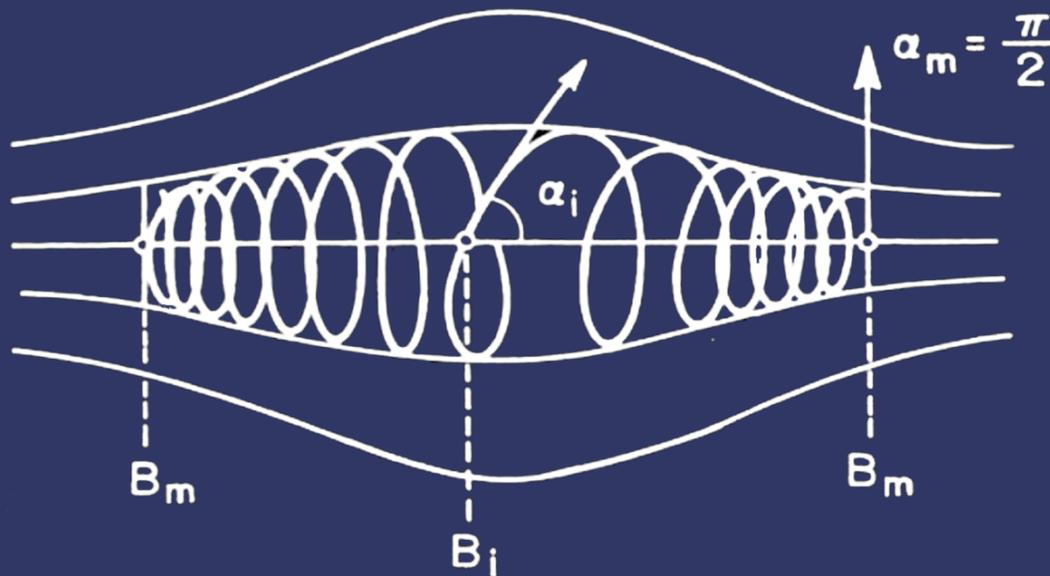
Uniform rotation at cyclotron frequency:

Bounce Motion in the Magnetic Bottle

$$\mu = \frac{p_{\perp}^2}{2mB} = \frac{p^2}{2mB_{\text{m}}} \quad F_{\parallel} = \frac{1}{\gamma} \mu \cdot \nabla B$$

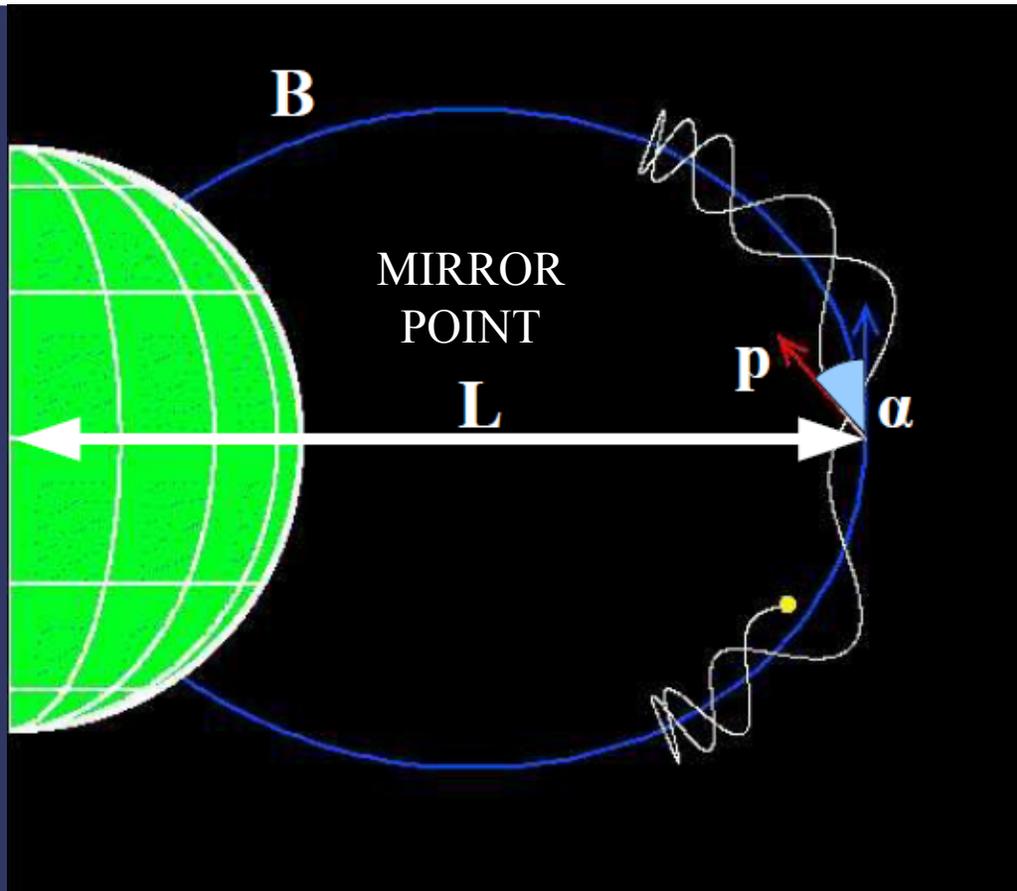
Charged particles undergo bounce motion in the magnetic bottle.

Pitch-angle of the particle changes from the maximum value on the equator to 90° in the mirror point.



Courtesy of A. Ukhorskiy

Bounce Motion

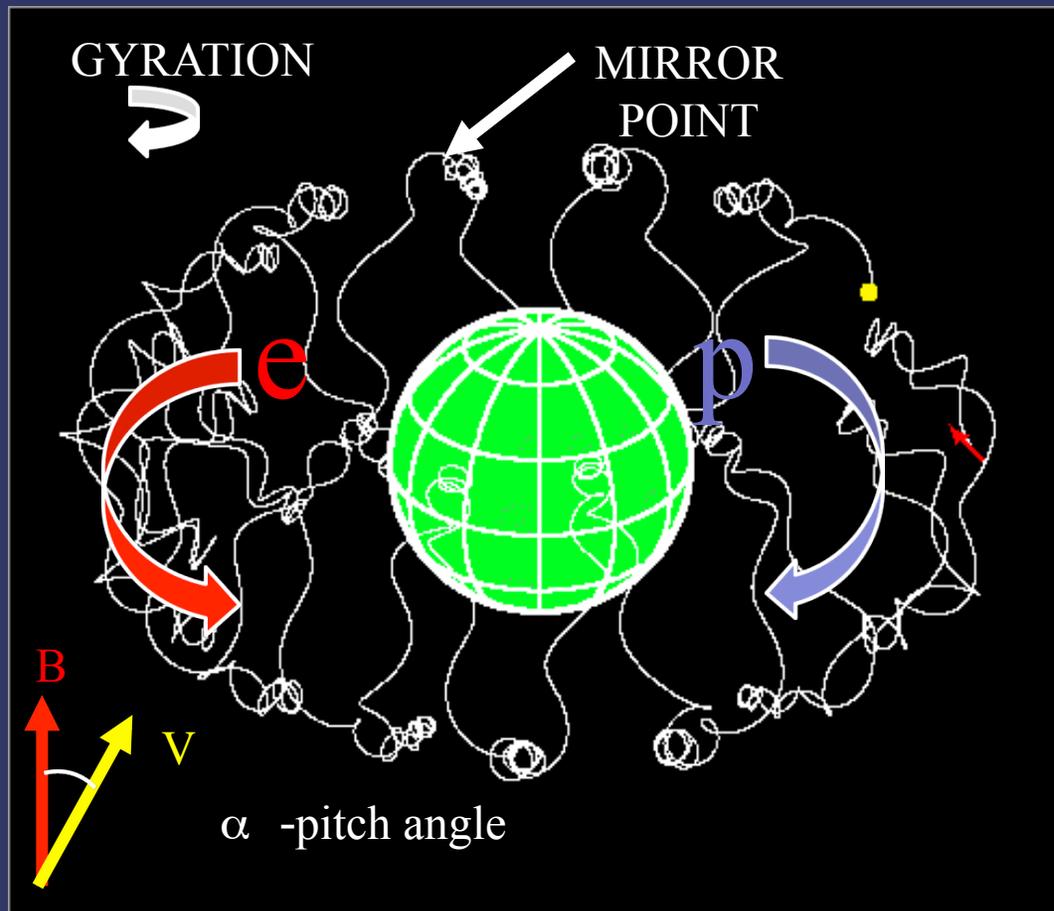


Pitch-angle – the angle between the magnetic field and particle's velocity.

90° pitch-angle particles will stay in the equatorial plane

Small pitch-angle particles **will be lost to the atmosphere.**

Trajectories of the Radiation Belt Particles



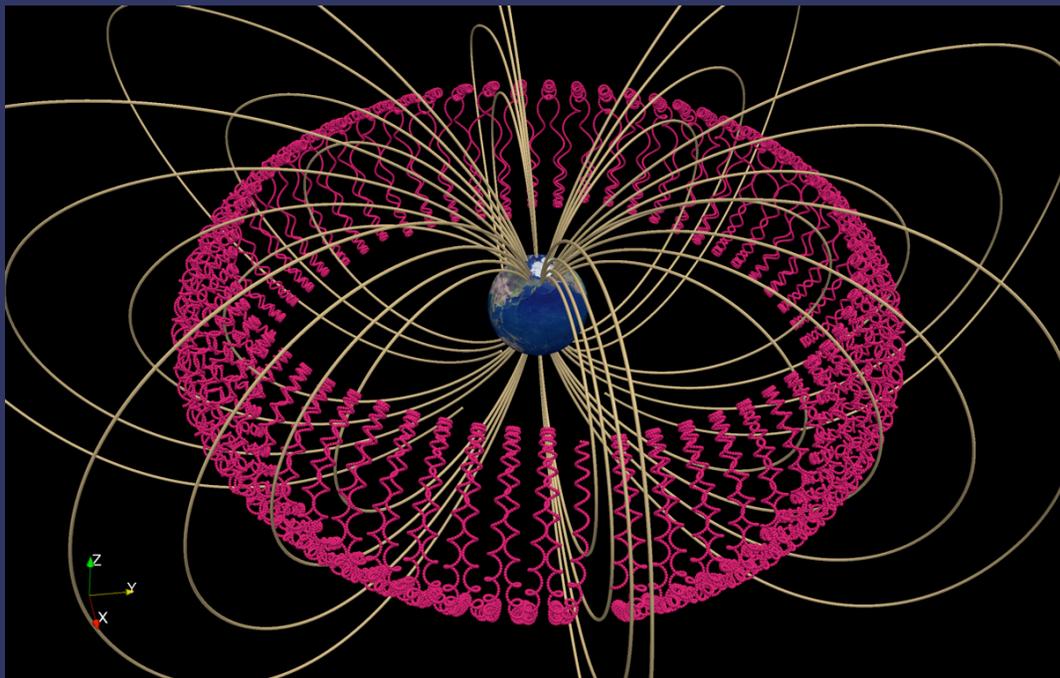
Gyro motion around the field line,

Bounce motion between the mirror points

Slow drift around the Earth.

Each type of periodic motion can be associated with an adiabatic invariant.

Simulations of the drift of 1 MeV electron, $m=20m_e$



Gyro motion around the field line,

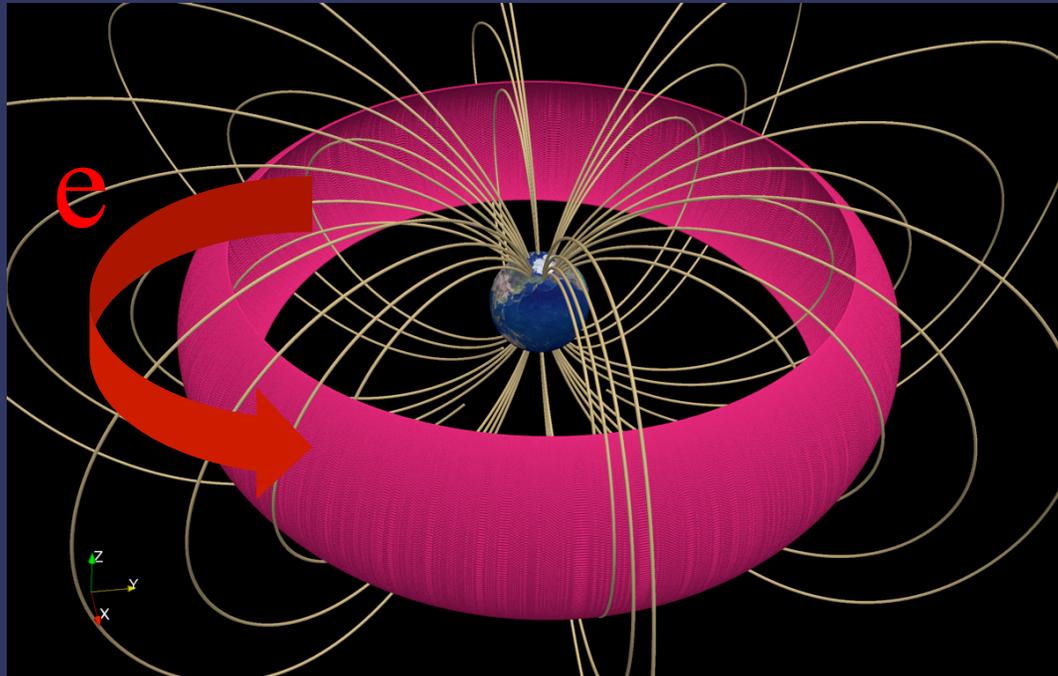
Bounce motion between the mirror points

Slow drift around the Earth.

Each type of periodic motion can be associated with an adiabatic invariant.

Courtesy of Sasha Ukhorskiy

Simulations of the drift of 1 MeV electron



Gyro motion around the field line,

Bounce motion between the mirror points

Slow drift around the Earth.

Each type of periodic motion can be associated with an adiabatic invariant.

Courtesy of Sasha Ukhorskiy

Adiabatic invariants

$$\mu = \frac{p_{\perp}^2}{2m_0B}$$

$$J = \oint p_{\parallel} ds$$

$$\Phi = \int B dS$$

$$L^* = \frac{2\pi M}{\Phi R_E}$$

$$J(E, \alpha) = f(\mu, J, L^*) / p^2$$

$$\frac{\partial f}{\partial t} = \sum_{m, n=1}^3 \frac{\partial}{\partial J_m} D_{J_m J_n} \frac{\partial}{\partial J_n} f ;$$

Fokker-Planck equation

Radial diffusion

Pitch-angle diffusion

Energy diffusion

$$\frac{\partial f}{\partial t} = L^2 \frac{\partial}{\partial L} L^{-2} D_{LL} \frac{\partial f}{\partial L} + \frac{1}{T(y)y} \frac{\partial}{\partial y} T(y)y D_{yy} \frac{\partial f}{\partial y} + \frac{1}{p^2} \frac{\partial}{\partial p} p^2 D_{pp} \frac{\partial f}{\partial p}$$

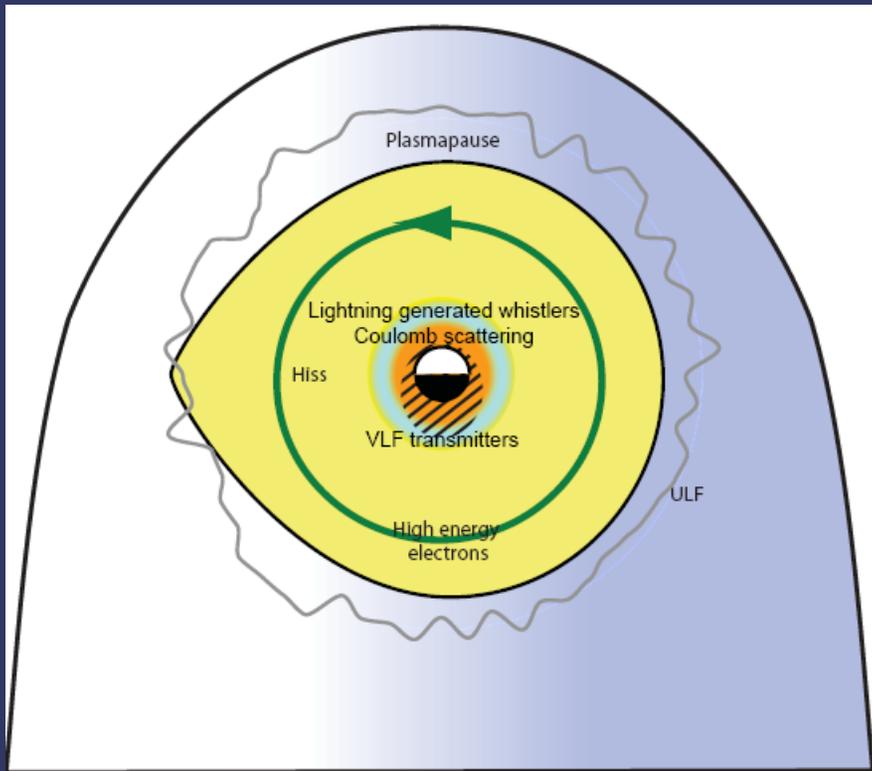
$y = \sin(\alpha)$

Inward **radial diffusion** can redistribute fluxes and provide energization or loss

Pitch angle scattering produces loss of electrons to the atmosphere

Energy diffusion provides local source of particles

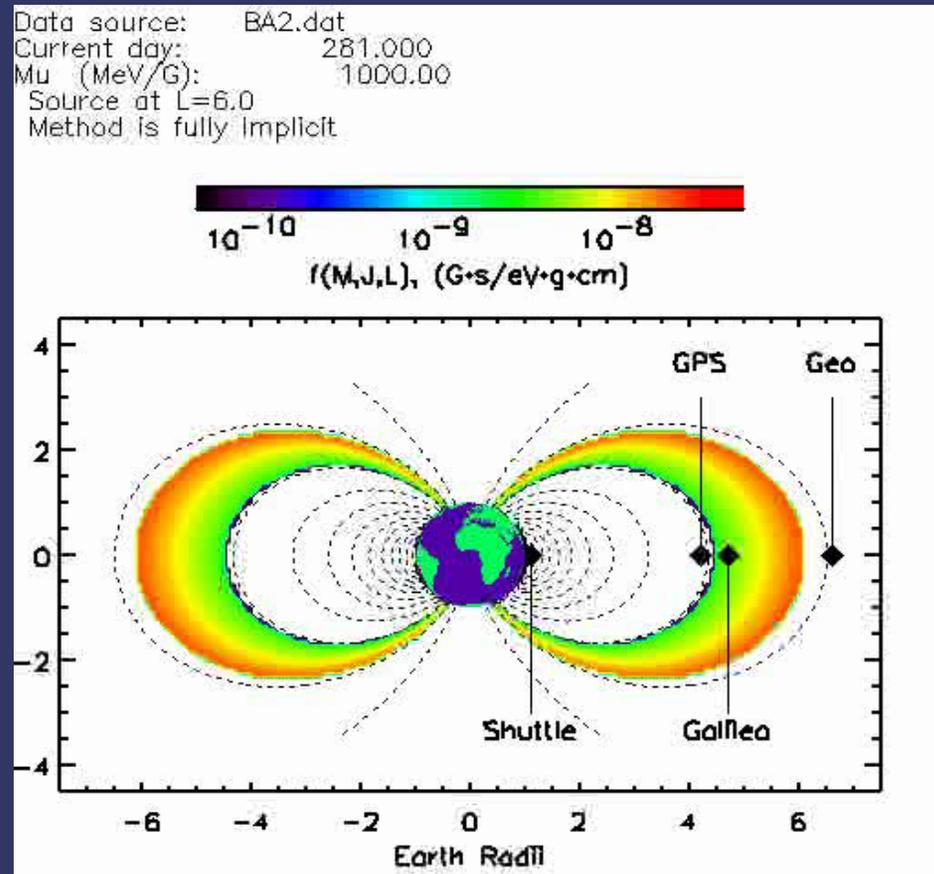
Dominant Acceleration and Loss Mechanisms during Quiet Times



Inward radial diffusion driven by ULF waves. Plasmaspheric Hiss is a dominant scattering mechanism at $L > 3$. Coulomb collisions, lightning generated whistlers, anthropogenic VLF zone.

Shprits et al, 2009

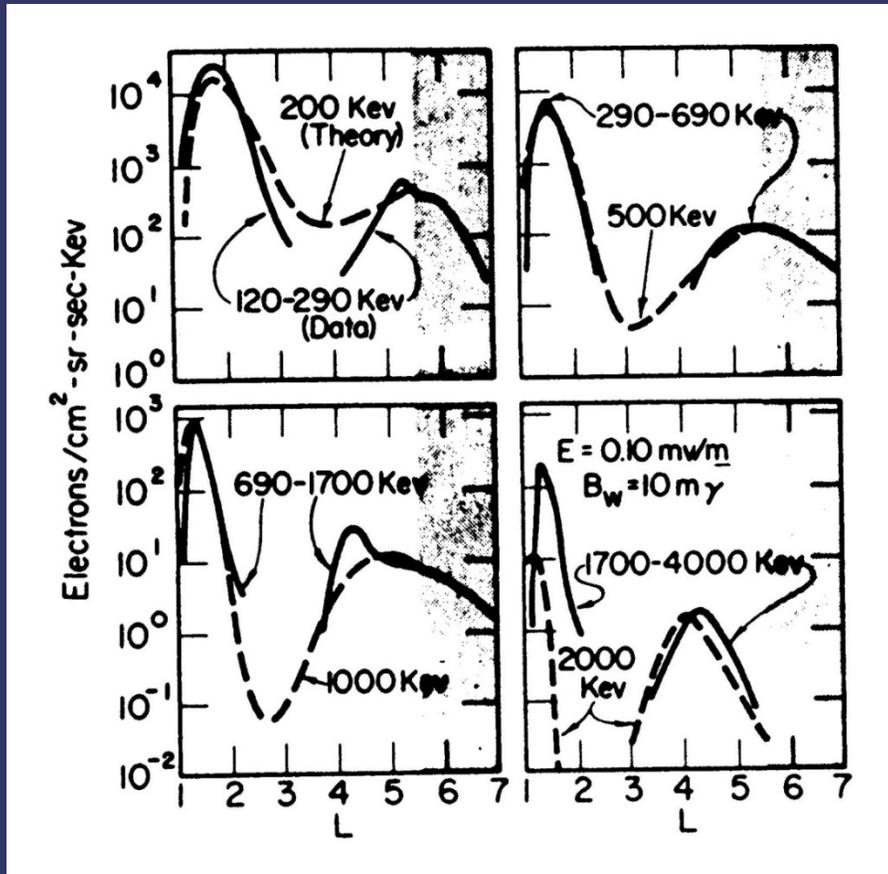
Radial Diffusion



Interactions with ULF waves causes random radial displacements that can result in the radial diffusion. Particles that are diffused into the inner region of strong magnetic field will gain energy.

Courtesy of Richard Horne

Theoretical flux profiles at fixed energy

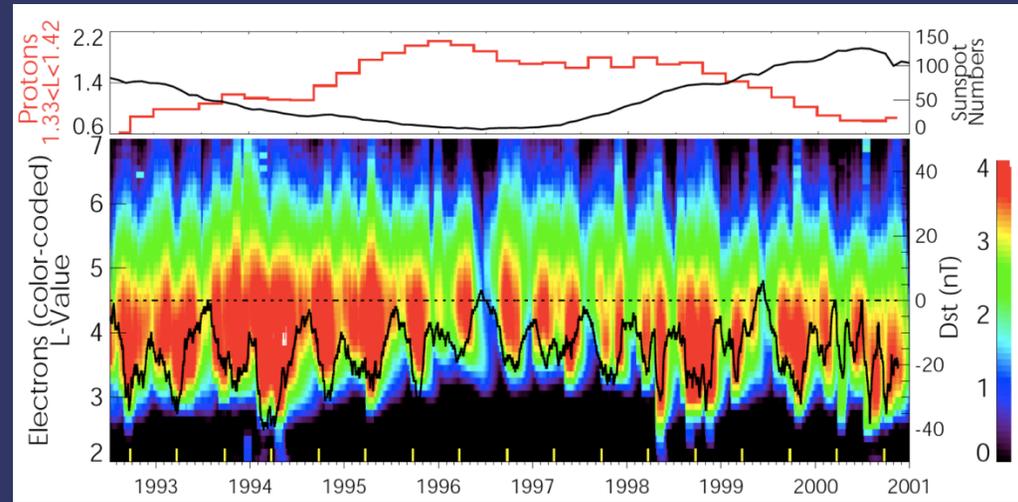
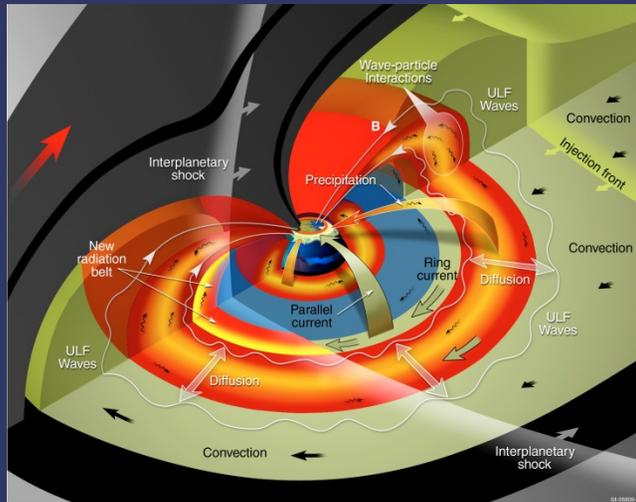


$$\frac{\partial f}{\partial t} = L^2 \frac{\partial}{\partial L} \left[D_{LL} L^{-2} \frac{\partial f}{\partial L} \right] - \frac{f}{\tau_{\text{effective}}}$$

Electron profiles obtained from the modeling of the quiet time inward diffusion
 Theoretical profiles of fluxes at constant energy agree with quiet time observations.

Lyons and Thorne, 1972

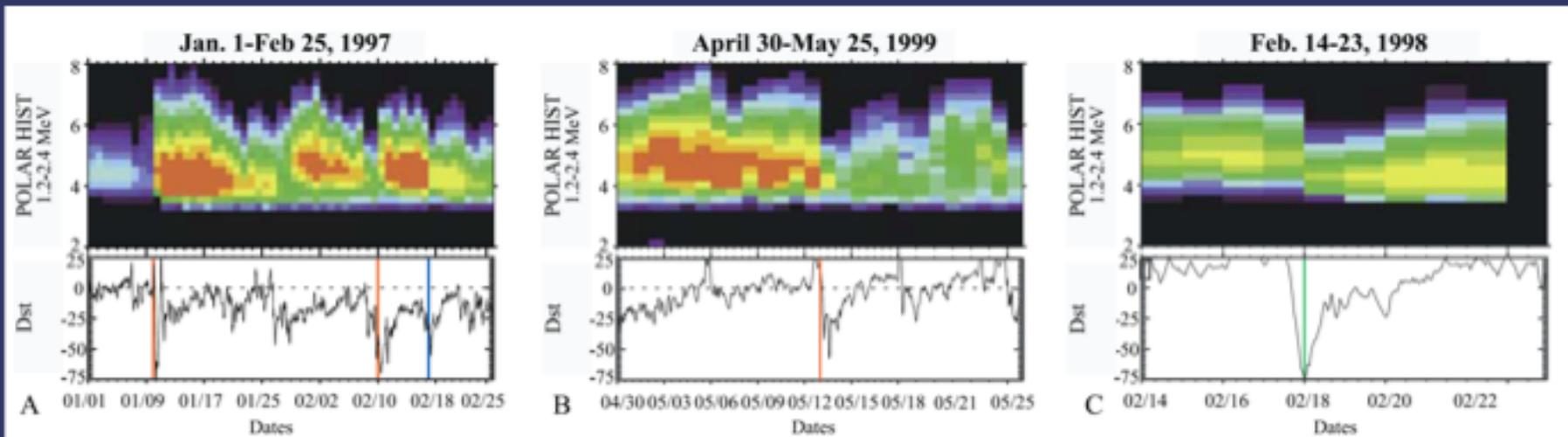
Dynamical Variability of the Radiation Belt Electrons



Creation and variation of radiation populations are produced by a complicated interplay of multiple processes. A broad range of coordinated measurements is needed to sort them out. How processes interact with each other under varying conditions to generate real space environments is unknown. Profound mysteries remain because existing observations are insufficient to resolve the system science.



Competition Between Acceleration and Loss



- (1) Cause dramatic radiation belt enhancement;
- (2) Deplete radiation belt fluxes;
- (3) Cause no substantial effect of flux distributions;

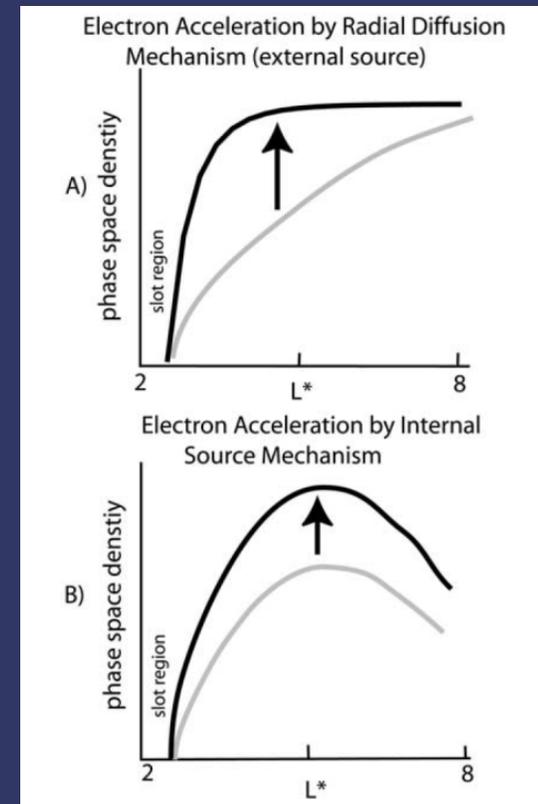
How to Differentiate Between Diffusion from the Outer Source and Local Acceleration processes

(A) Phase space density increase caused by radial

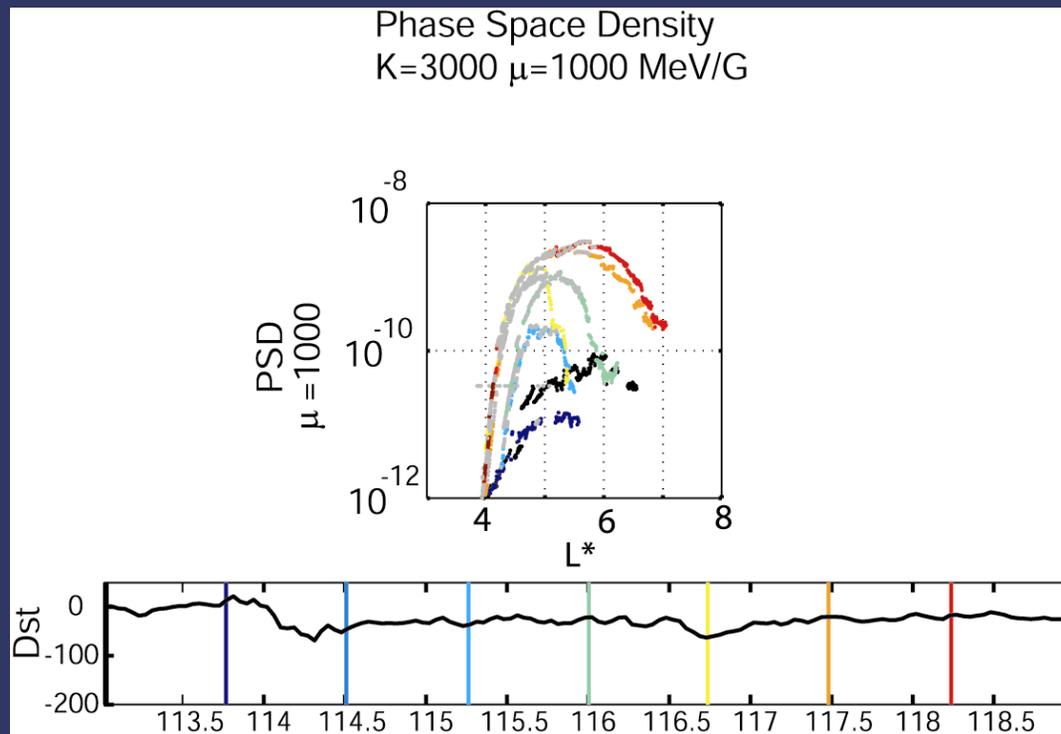
diffusion from an external source.

(B) Phase space density increases predicted by local internal source acceleration mechanisms.

L^* is approximately the distance from the Earth measured in Earth radii



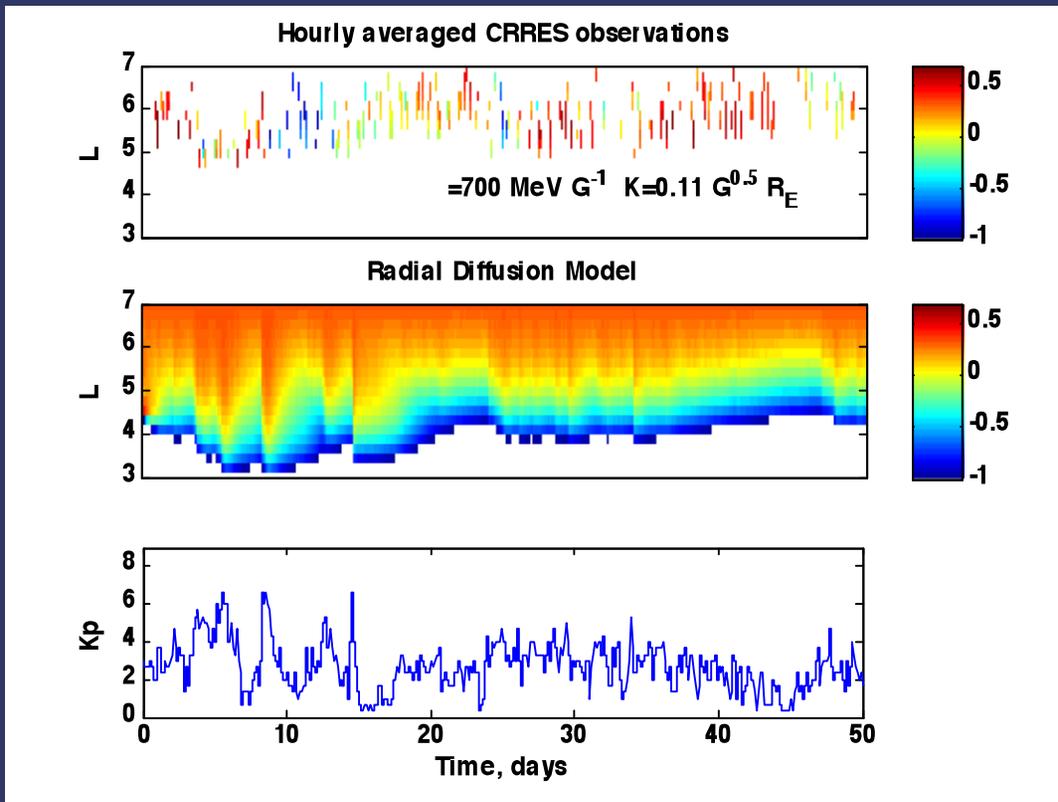
Observations of the Radial Profile of Phase Space Density



Evidence of the presence of the **local acceleration source**

Building up peaks are consistent with the local acceleration source produced by resonant interactions with whistler mode chorus waves.

[Green and Kivelson, 2004]



- Observations are sparse
- Model is continuous but may be missing essential physics.

Data Assimilation with Kalman Filter

Assume initial state and data and model errors

Make a prediction of the state of the system and error covariance matrix, using model dynamics

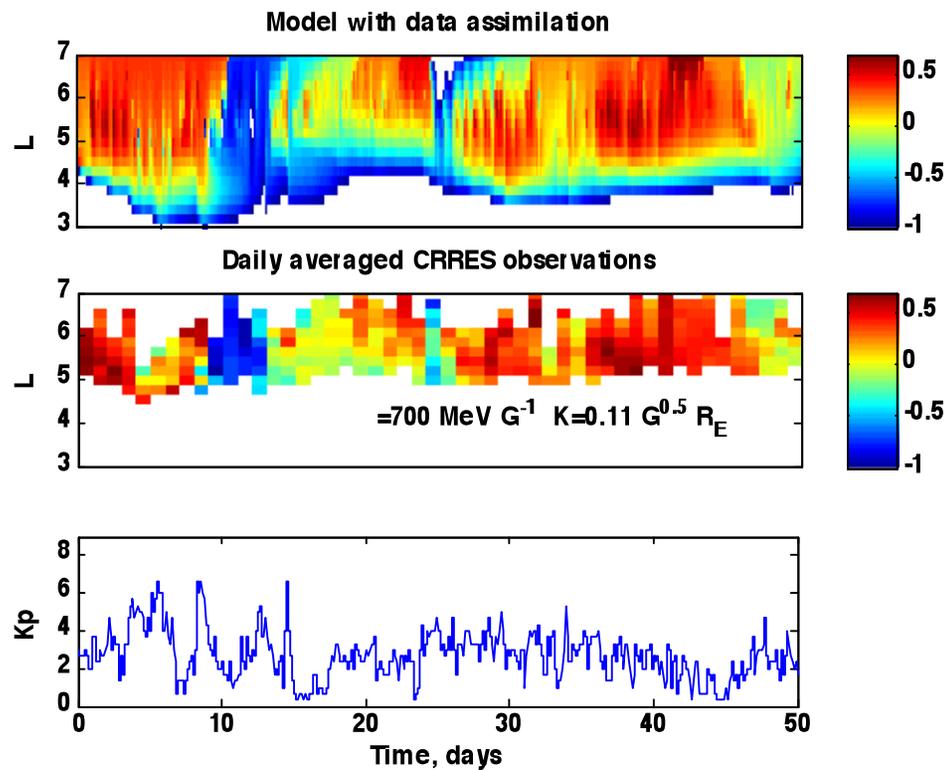
$$w_k^f = \Psi_{k-1} w_{k-1}^f$$

Compute Kalman gain and innovation vector

Update state vector using innovation vector

$$w_k^a = w_k^f + Kd$$

Compute updated error covariance matrix



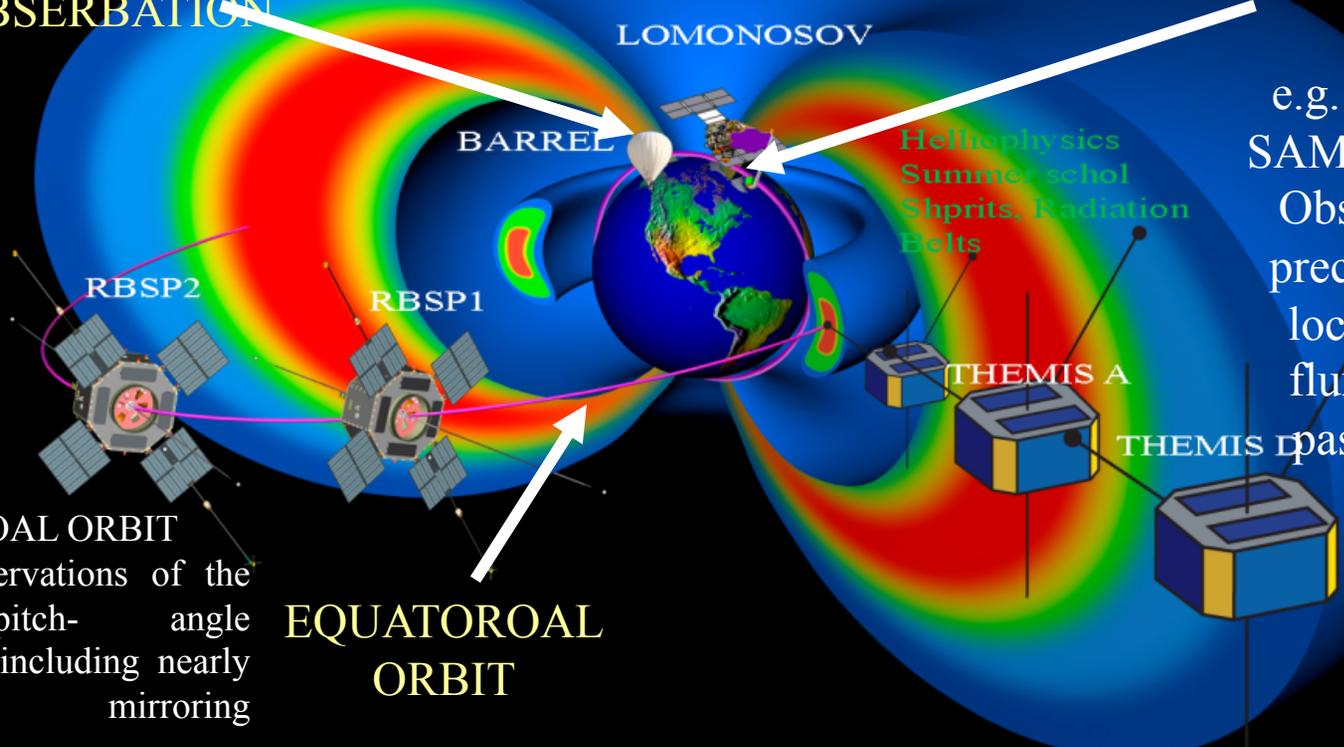
Data assimilation can fill in spatio-temporal gaps.

Data assimilation is consistent with daily average values but allows to minimize errors and can be used with multiple spacecraft

Multi Point Observations

BALLOON
OBSERVATION

LEO ORBIT



e.g. Lomonosov,
SAMPEX, NOAA.
Observations of
precipitating and
locally trapped
fluxes, Several
passes per day

EQUATORIAL ORBIT
Allows observations of the
whole pitch-angle
distribution including nearly
equatorially mirroring
particles.

EQUATORIAL
ORBIT

Heliophysics

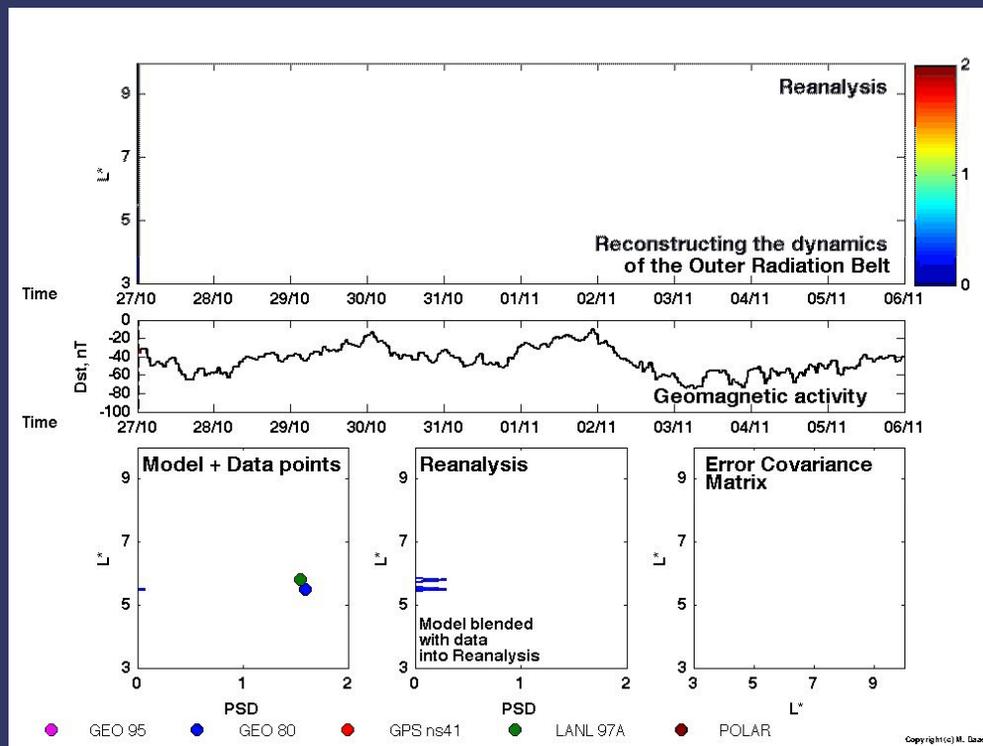
2012 Summer School
Heliophysical Exploration

31 May - 7 June 2012
Boulder, Colorado

Application Deadline: 18 December 2011

Radiation Belts

6/4/12

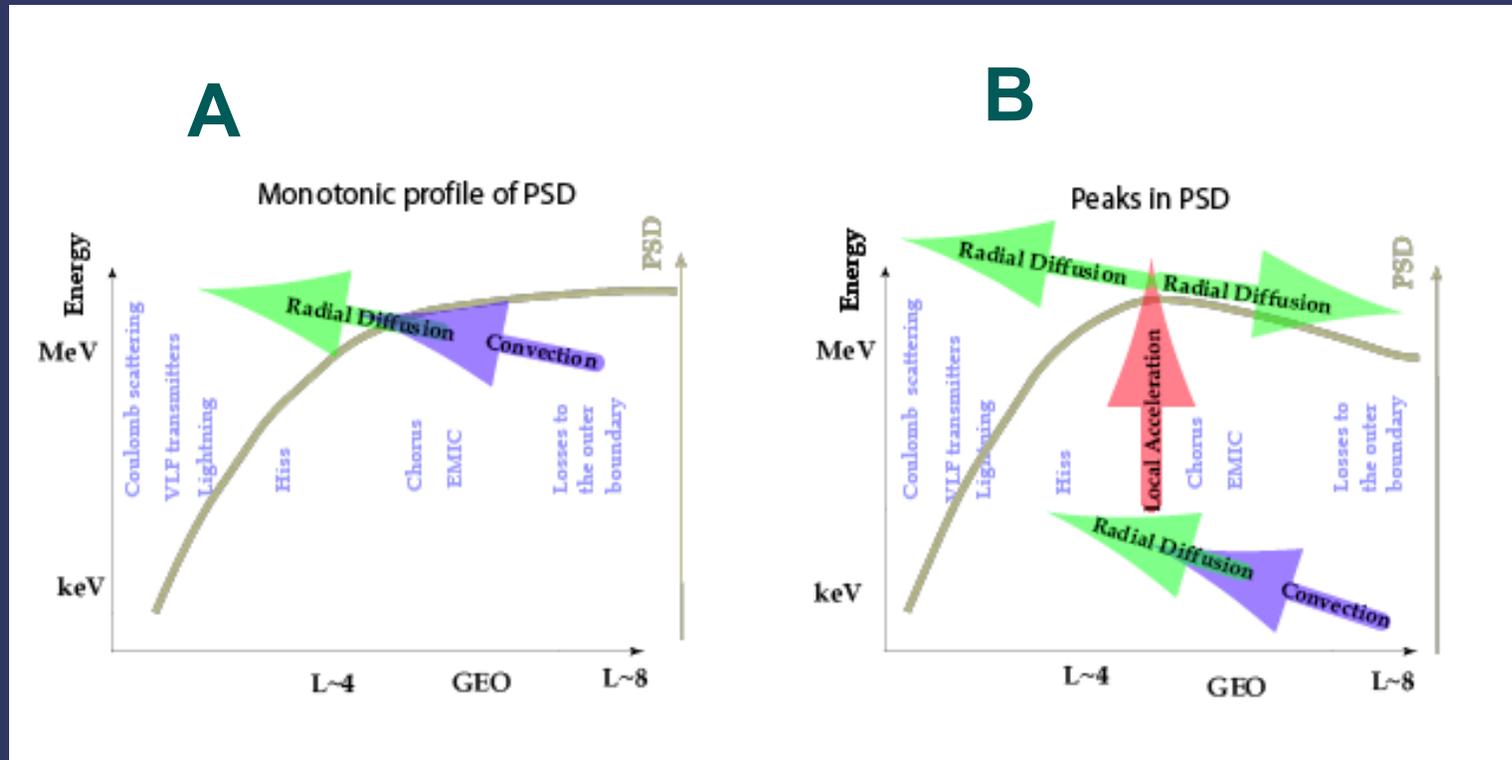


Data is blended with the model according to the underlying structure of data and model errors.

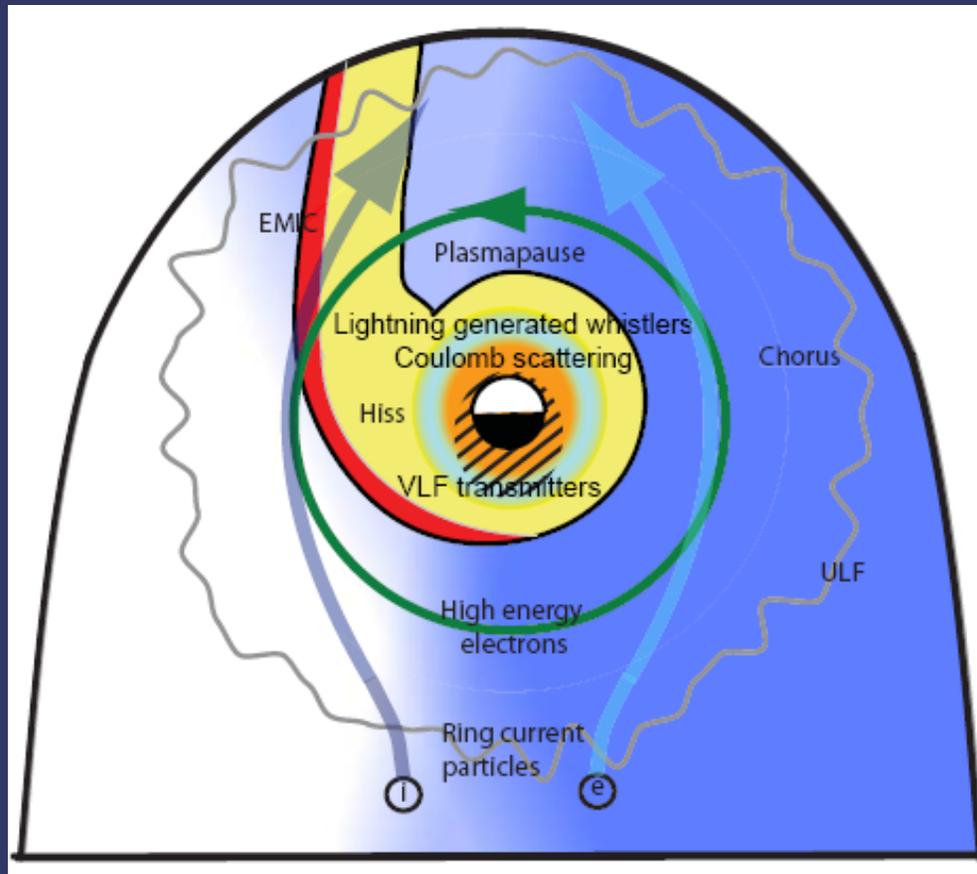
Data from 4 spacecraft is assimilated and radial profile of PSD is dynamically reconstructed

Provided by Marianne Daae

Acceleration of electrons to relativistic energies



Storm Time Acceleration Mechanisms



Inward radial diffusion driven by ULF magnetic field.

Energy diffusion due to resonance interactions with whistler mode chorus waves

Shock induced acceleration

Adiabatic acceleration

Resonant Wave particle interactions

$$\omega - k_{\parallel} v_{\parallel} = n \Omega_e / \gamma, \quad n = 0, \pm 1, 2, 3, \dots$$

$$\gamma = (1 - v^2 / c^2)^{-1/2}$$

where ω is the wave frequency, k_{\parallel} and v_{\parallel} are the components of the wave vector and electron velocity parallel to the ambient magnetic field, respectively, Ω_e is the electron gyrofrequency

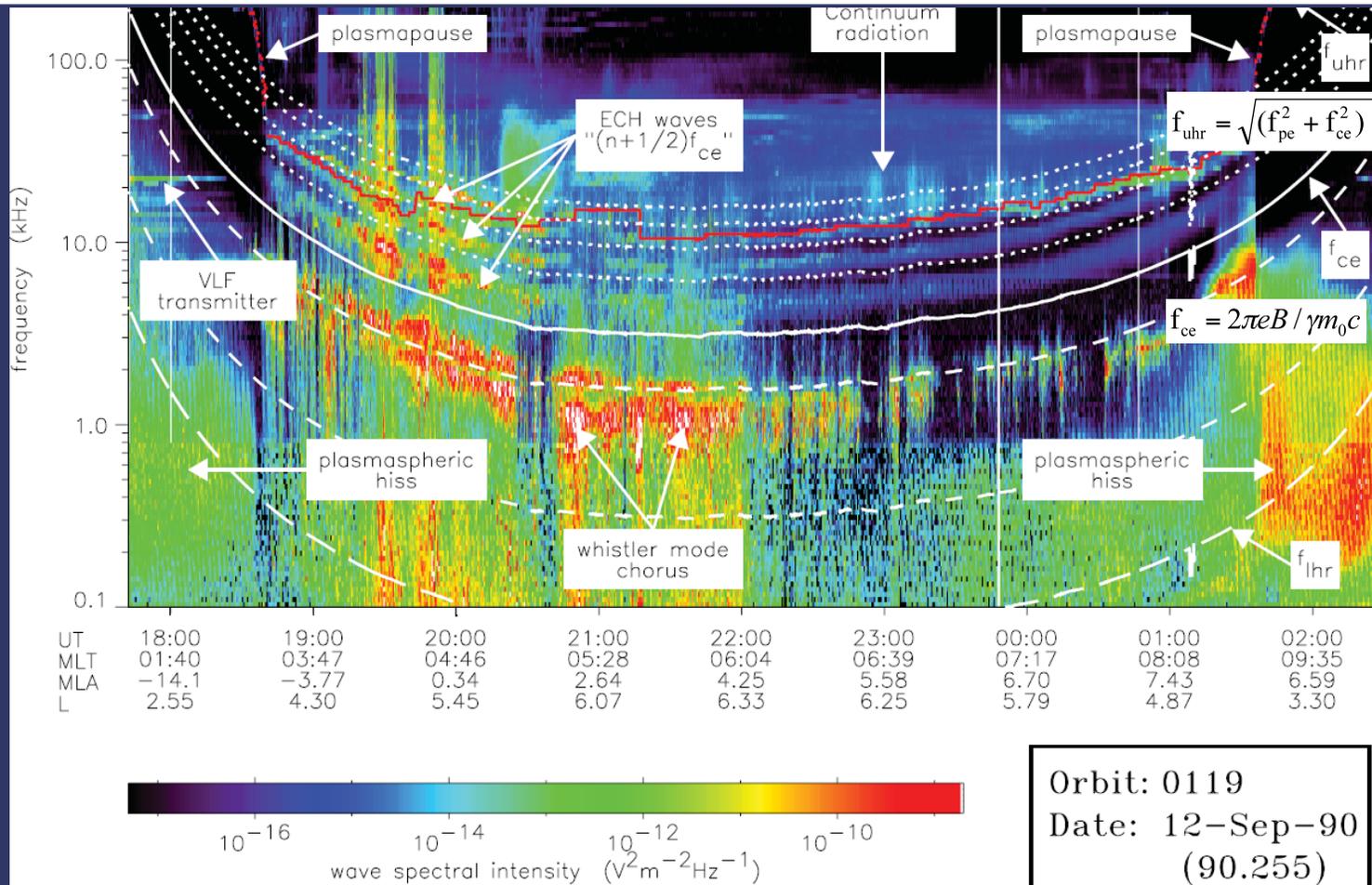
Resonant interactions occur when multiples of gyrofrequency equal the wave frequency

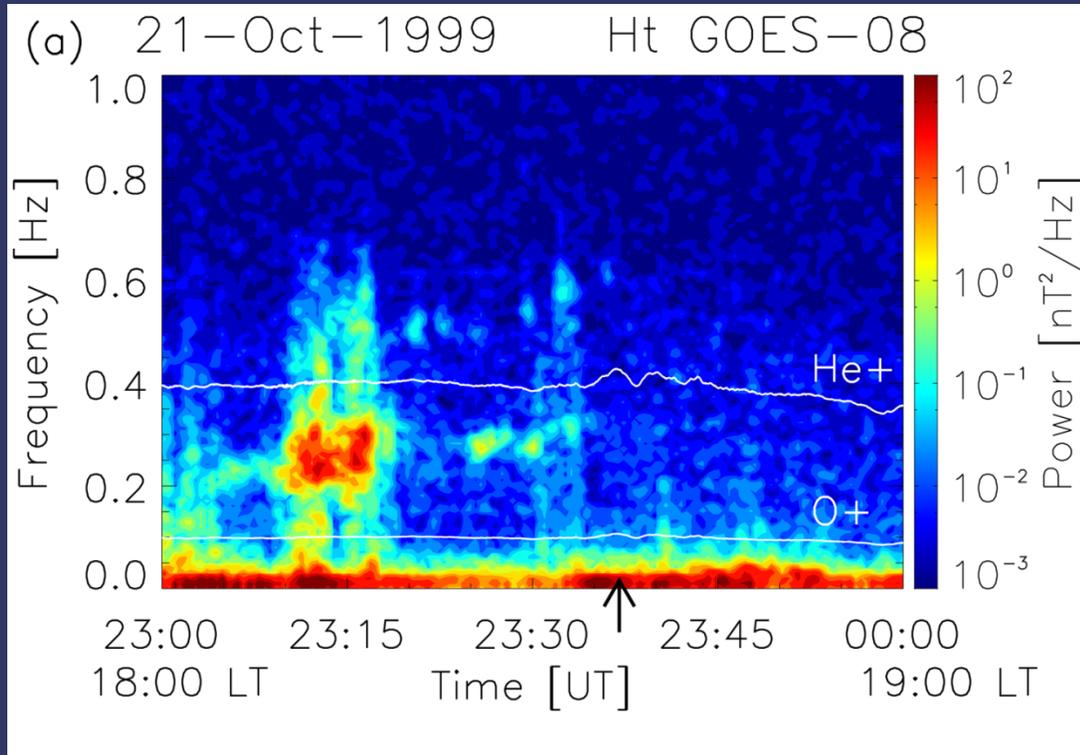
During resonant interactions, ELF and VLF waves can violate the first and second adiabatic invariants and diffuse electrons in pitch-angle and energy.



$$f_{\text{uhr}} = \sqrt{(f_{\text{pe}}^2 + f_{\text{ce}}^2)}$$

Dynamical Spectrogram



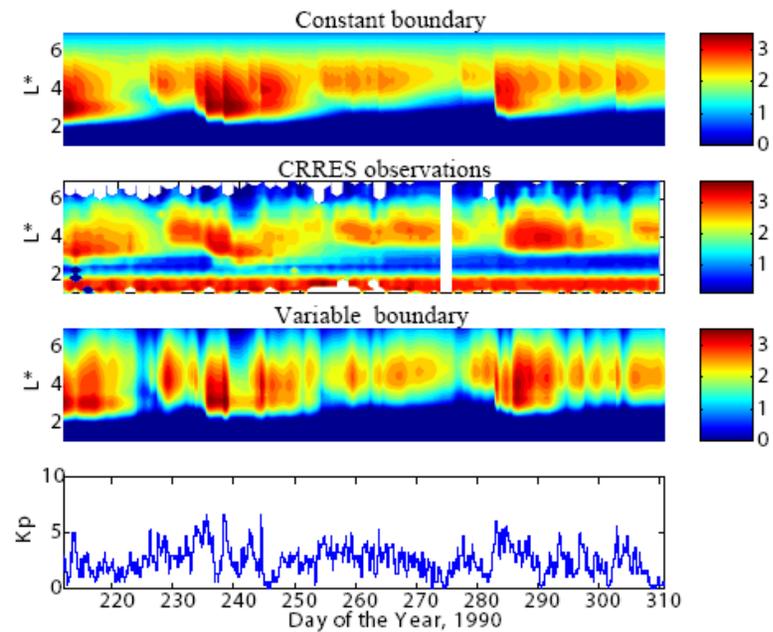
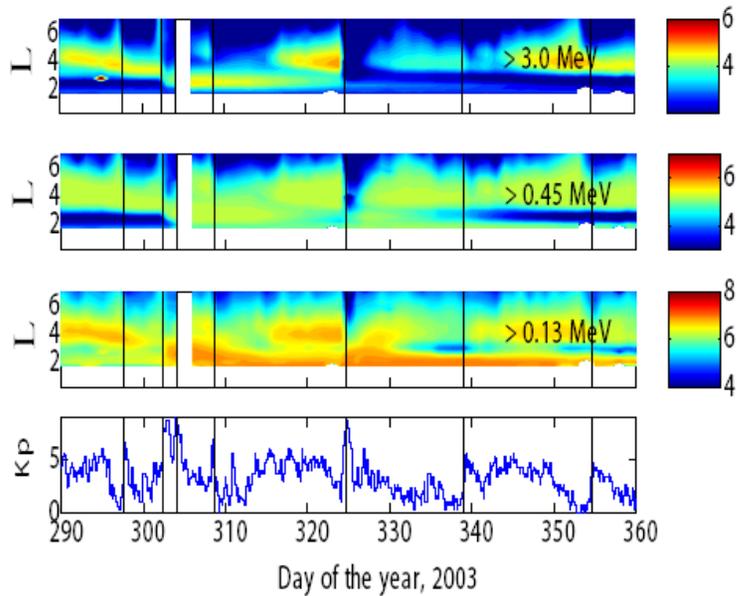


Very intense waves are observed in He band.

EMIC waves can produce scattering of relativistic electrons on very short time scales.

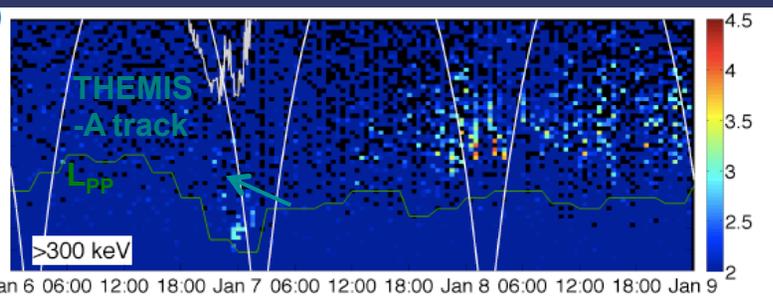
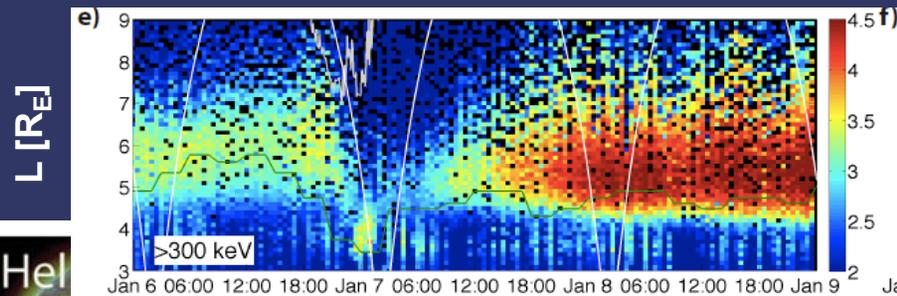
Fraser et al., 2010

Loss of Relativistic electrons in the Radiation Belts



POES Trapped: $\alpha = 90$ deg at 850 km alt.

POES Precipitating: $\alpha = 0$ deg at 850 km alt.

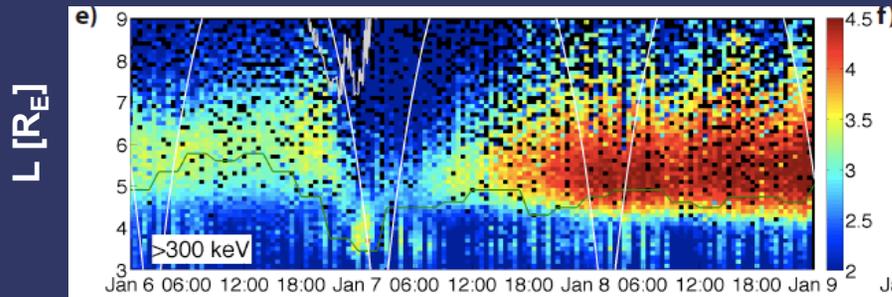


Radiation Belts

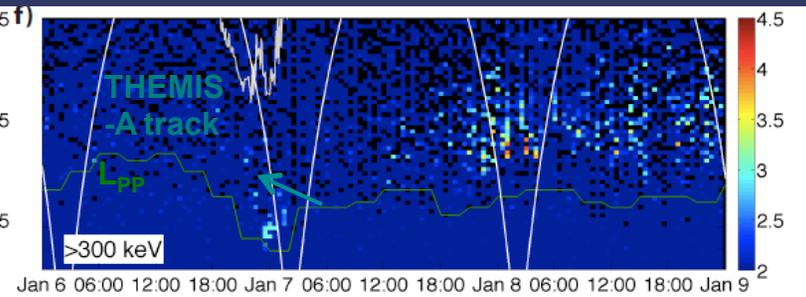
6/4/12



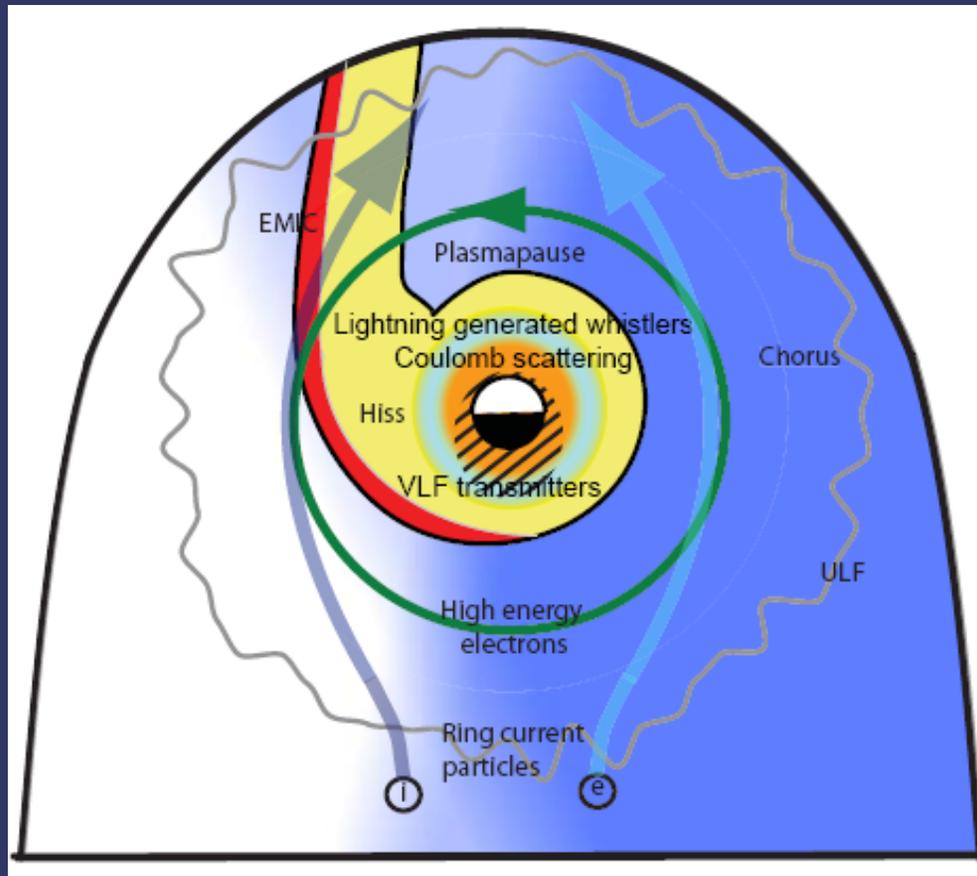
POES Trapped: $\alpha = 90$ deg at 850 km alt.



POES Precipitating: $\alpha = 0$ deg at 850 km alt.



Storm Time Loss Mechanisms



Inward radial diffusion driven by ULF magnetic field.

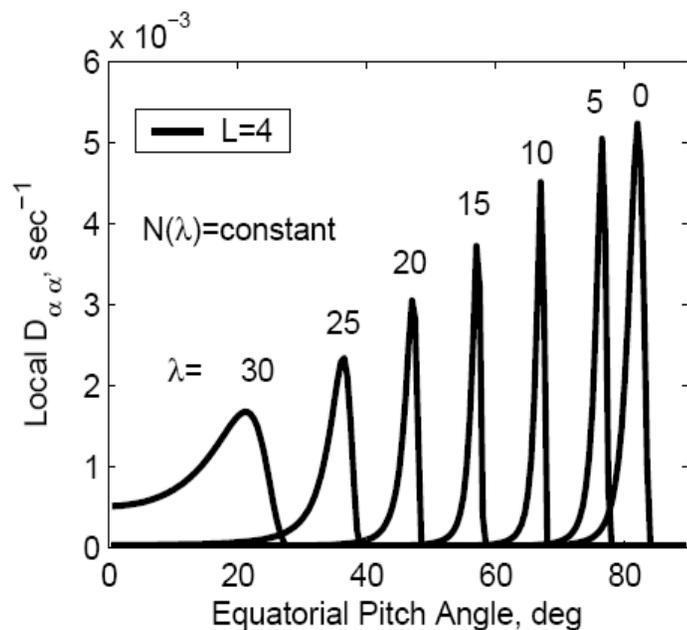
Energy and Pitch angle scattering due to resonance interactions with different waves

Losses to magnetopause and outward radial diffusion.

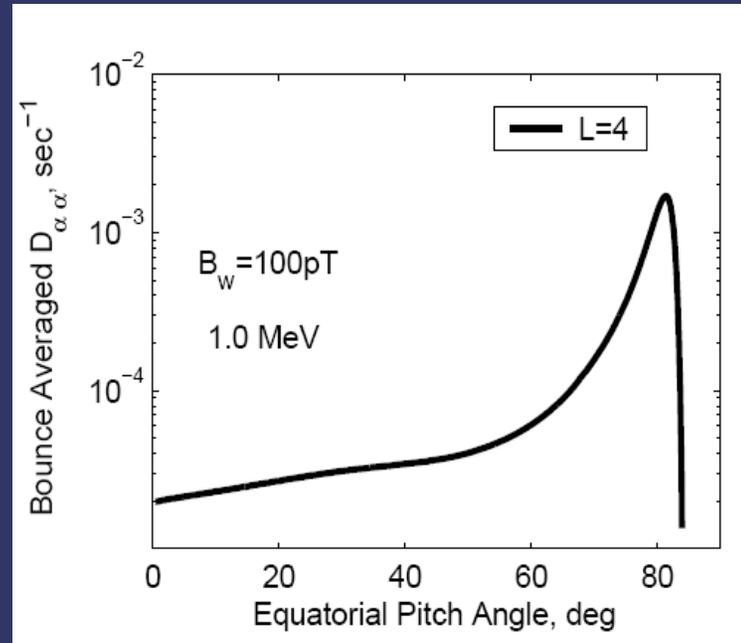
Local and Bounce Average Diffusion Coefficients

$$\omega - k_{\parallel} v_{\parallel} = n \Omega_{\text{gyro}} / \gamma.$$

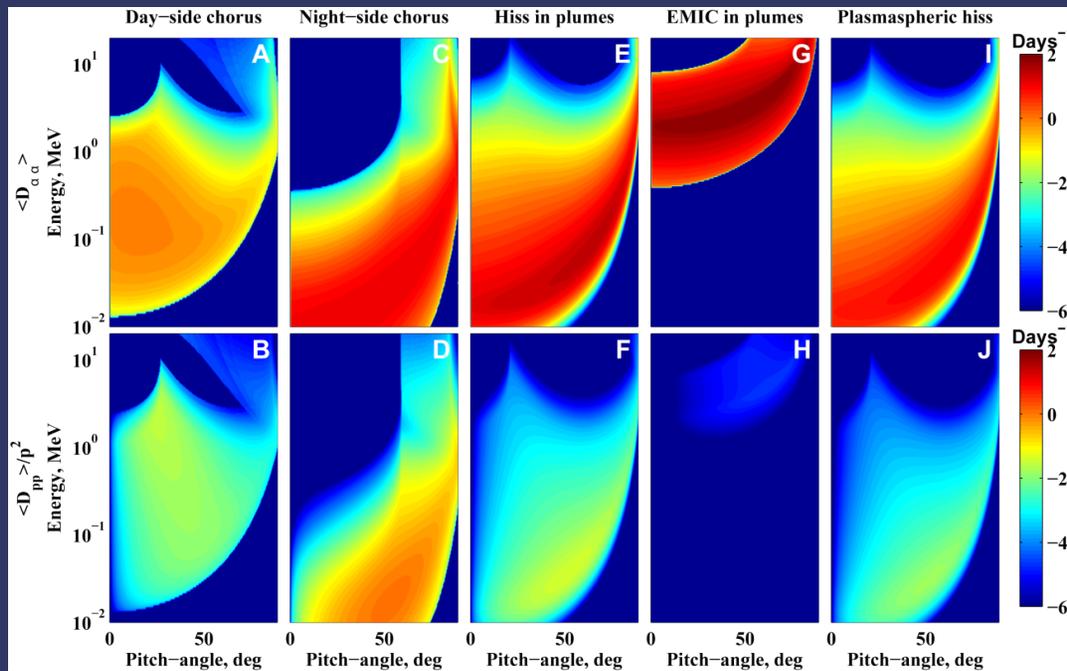
Local Diffusion Rates



Bounce Averaged Diffusion Rates



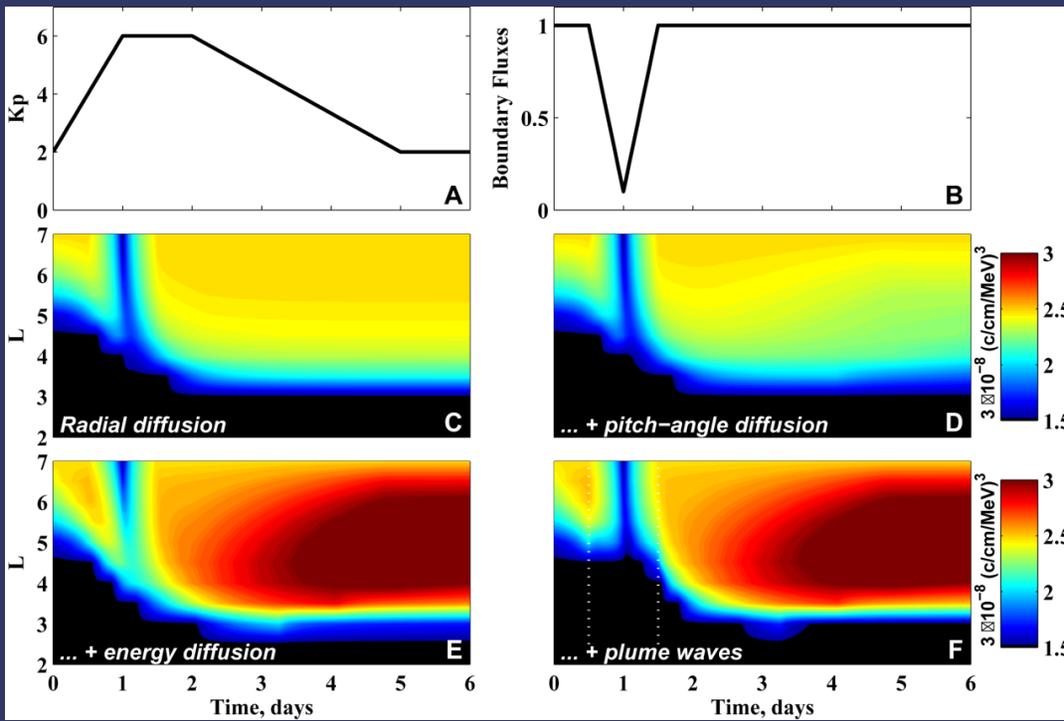
Pitch-angle and energy diffusion coefficients due to resonant wave-particle interactions

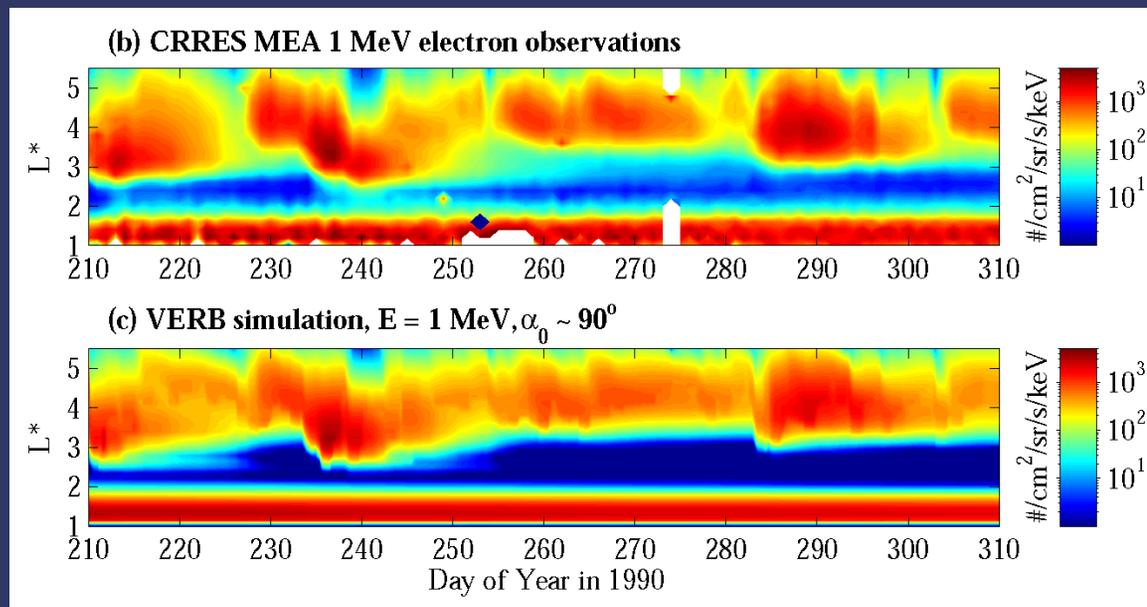


Pitch-angle diffusion results in a loss to the atmosphere

Energy diffusion can accelerate electrons locally

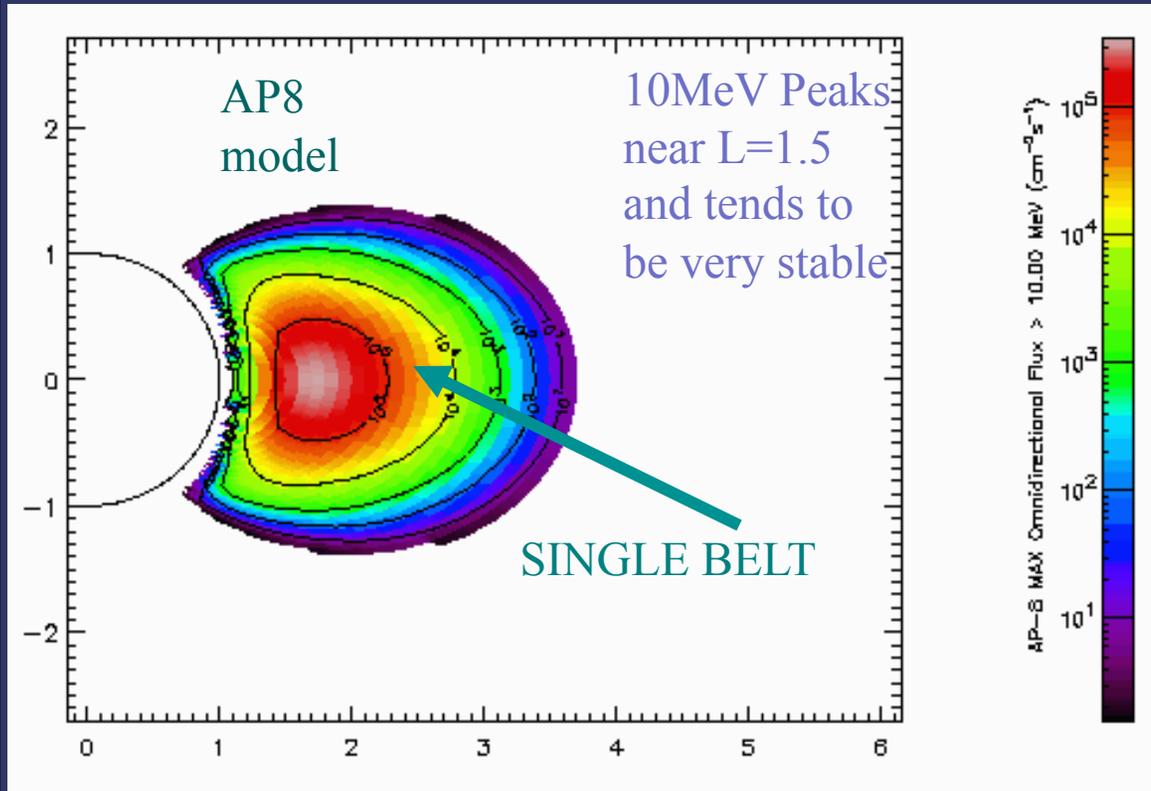
Diffusion rates as due to resonance wave particle interactions with various plasma waves.





- VERB predicts the instantaneous location of the upper boundary of the slot region, the empty slot region, the stable inner belts, the location of peak of fluxes and amplitude of fluxes.

Proton Radiation Belt: Acceleration Mechanisms

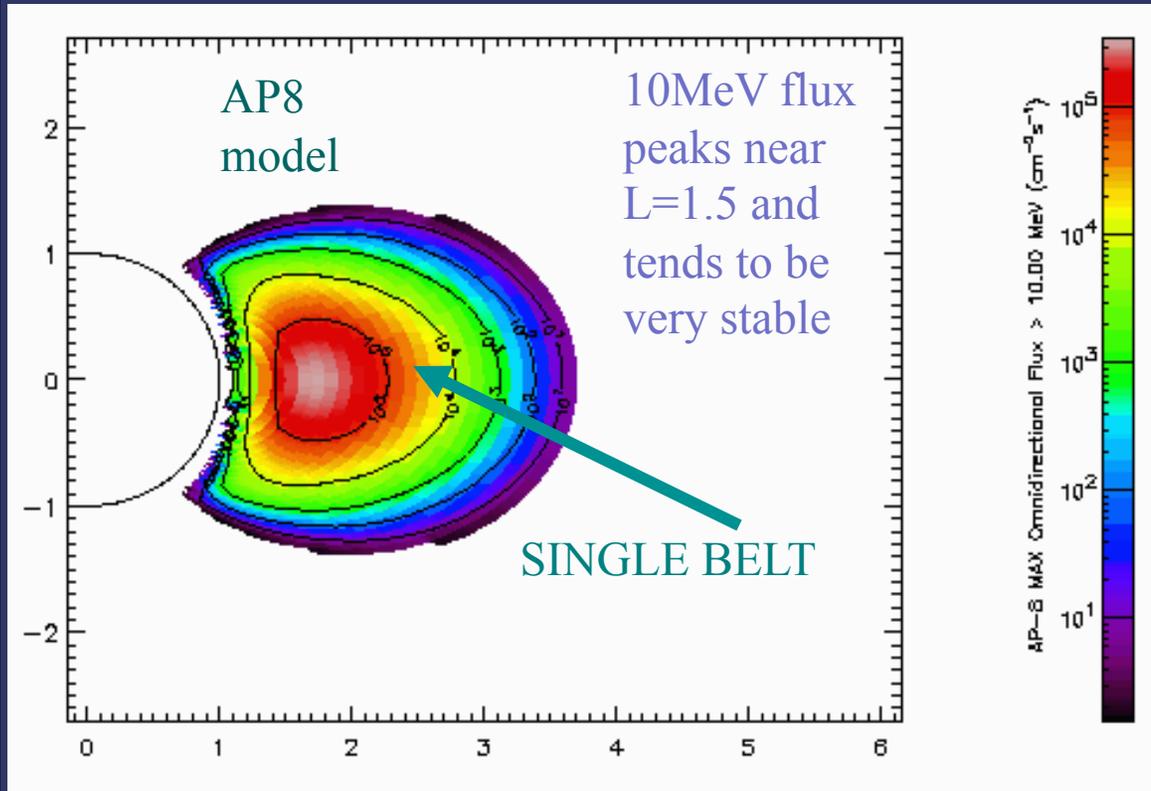


Higher energies peak at lower L-shells

Protons are transported inwards by the large induced electric fields that arise as a large shock passes the Earth.

Can be also transported inwards by radial diffusion

Proton Radiation Belt: Loss Mechanisms

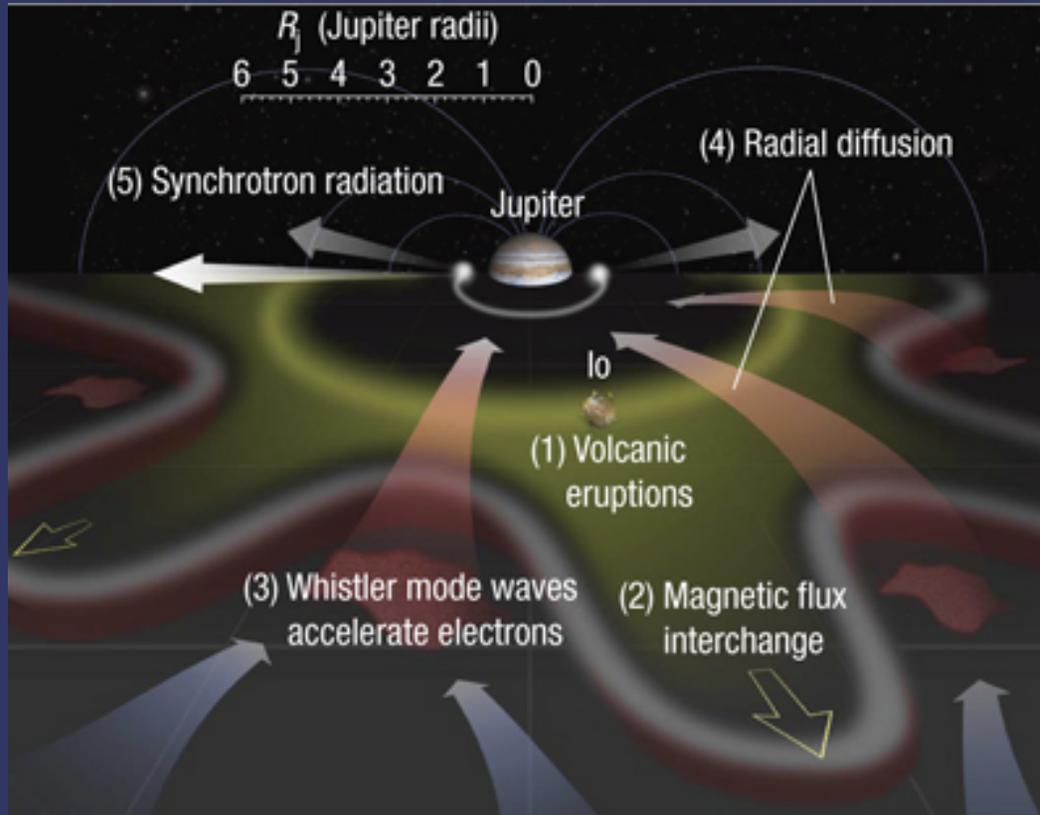


Scattering by EMIC waves

Field line curvature scattering

Adiabatic changes in the inner belt.

Radiation Belts on Jupiter



Inward radial diffusion, interchange instability, and energy diffusion chorus waves accelerate electrons to ultra-relativistic energies.

Losses :interactions with dusty plasmas, planets, synchrotron loss and atmospheric loss.

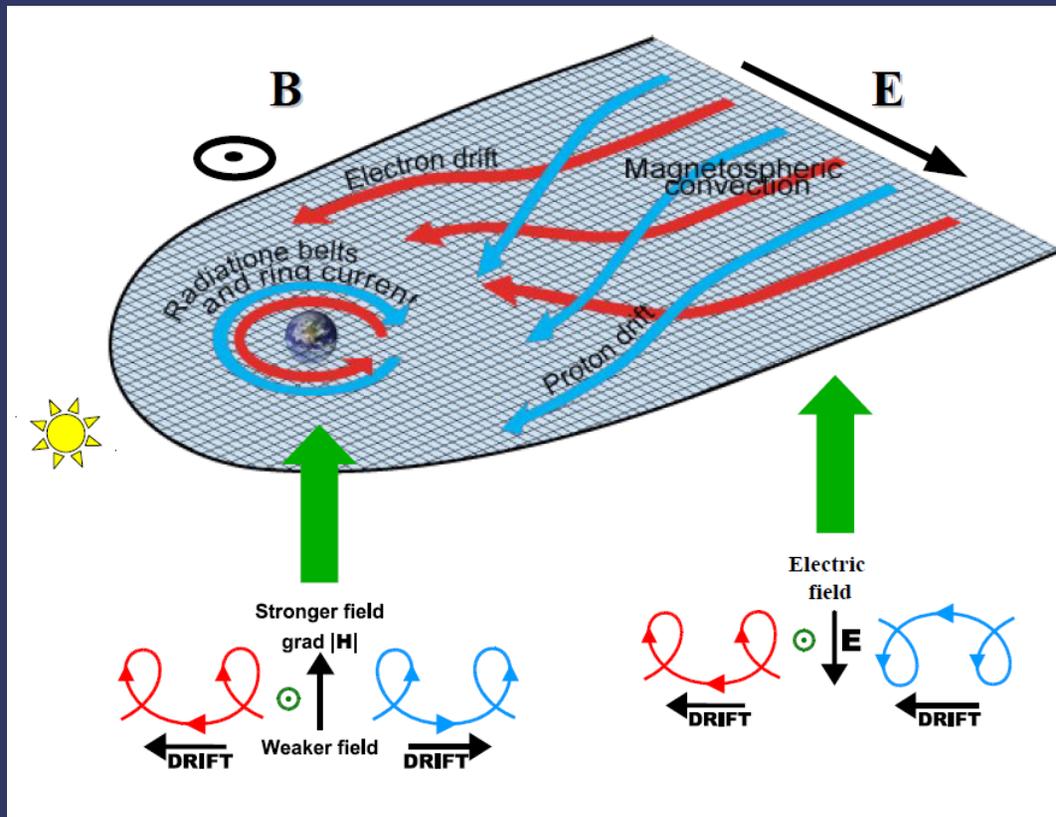
Horne et al., 2008

Acceleration and Loss of Radiatino Belts on Jupiter

- (1) Volcanic gases from Io are ionized and form a cold dense plasma torus around Jupiter.
- (2) Jupiter's rapid rotation drives magnetic flux interchange and excites whistler-mode waves.
- (3) Gyro-resonant wave–particle interactions accelerate electrons to relativistic energies.
- (4) Radial diffusion transports electrons towards the planet and accelerates them to even higher energies via betatron and Fermi processes.
- (5) Intense synchrotron radiation is emitted from ultra-relativistic electrons close to the planet ($1.4R_j$).



Particle Trajectories of Ring Current and Radiation Belt Particles



Drift of lower energy particles is dominated by $E \times B$ drift.

Radiation Belt particles are subject to the gradient and curvature drifts and will drift around the Earth.

Electrons –eastward,
Ions-westward.

Subbotin, et al., 2011