

Exploring the Sun and its effects on the Earth's atmosphere and physical environment...

HIGH ALTITUDE OBSERVATORY

Planetary Dynamos

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Outline

I) What is a Dynamo?

- Magnetic field creation vs dissipation
- Conditions for a planetary dynamo

II) A whirlwind tour of the Solar System

- Observed magnetic fields
- Internal structure

III) Convection in Rotating Spheres

- Dynamical balances
- Columns and waves

IV) Numerical Models

- General trends
- Case Studies (Earth, Jupiter...)



What is a (hydromagnetic) Dynamo?

An object (such as a star or a planet or a lab experiment) that converts the kinetic energy of fluid motions into magnetic energy



MHD Magnetic Induction equation

$$\frac{\partial \mathbf{B}}{\partial t} = \boldsymbol{\nabla} \times (\mathbf{v} \times \mathbf{B} - \eta \boldsymbol{\nabla} \times \mathbf{B})$$

Comes from Maxwell's equations (Faraday's Law and Ampere's Law)

$$\frac{1}{c}\frac{\partial \mathbf{B}}{\partial t} = -\boldsymbol{\nabla} \times \mathbf{E} \qquad \qquad \boldsymbol{\nabla} \times \mathbf{B} = \frac{4\pi}{c} \mathbf{J} \quad \text{(Assumes v << c)}$$

And Ohm's Law

 $\tilde{\mathbf{J}} = \sigma \tilde{\mathbf{E}}$

Magnetic diffusivity

$$\eta = \frac{c^2}{4\pi\sigma}$$
 electrical conductivity



How would you demonstrate this?

(Hint: have a sheet handy with lots of vector identities!)

$$E_m = \frac{B^2}{8\pi}$$



$$\mathbf{F}_P = \mathbf{E} \times \mathbf{B} = \left[\frac{\eta}{c} \mathbf{J} - \frac{1}{4\pi} \left(\mathbf{v} \times \mathbf{B}\right)\right] \times \mathbf{B}$$





 $\operatorname{Rm} = \frac{UD}{\eta}$

If Rm >> I the source term is much bigger than the sink term

....Or is it???



δ can get so small that the two terms are comparable

It's not obvious which term will "win" - it depends on the subtleties of the flow, including geometry & boundary conditions



What is a Dynamo? (A corollary)

A dynamo must sustain the magnetic energy (through the conversion of kinetic energy) against Ohmic dissipation

If v = 0 and $\eta = constant$ then the induction equation becomes

$$\frac{\partial \mathbf{B}}{\partial t} = -\eta \boldsymbol{\nabla} \times \boldsymbol{\nabla} \times \mathbf{B} = \eta \nabla^2 \mathbf{B}$$

The field will diffuse away (dissipation of magnetic energy) on a time scale of

A more careful calculation for a planet gives

 $\tau_d \approx \frac{R^2}{\pi^2 \eta}$

Earth: *τ*_d ~ **80,000** yrs

Jupiter: $\tau_d \sim 30$ million yrs

Planetary fields must be maintained by a dynamo or they would have decayed by now!

 $\tau_d \approx \frac{D^2}{n}$

Conditions for a Planetary (or Stellar) Dynamo

n Absolutely necessary

An electrically conducting fluid

- Stars: plasma
- Terrestrial planets: molten metal (mostly iron)
- Jovian planets: metallic hydrogen (maybe molecular H)
- Ice Giants: water/methane/ammonia mixture
- Icy moons: salty water

Fluid motions

Usually generated by buoyancy (convection)

→ Rm >> 1

 Too much ohmic diffusion will kill a dynamo



Conditions for a Planetary (or Stellar) Dynamo

n Not strictly necessary but it (usually) helps

Rotation

- Good: helps to build strong, large-scale fields (promotes magnetic self-organization)
- Bad: can suppress convection (though this is usually not a problem for planets)
- Turbulence (low viscosity / Re >> 1)
 - Good: Chaotic fluid trajectories good at amplifying magnetic fields (chaotic stretching)
 - Bad: can increase ohmic dissipation

Earth

Dynamo!

Field strength ~ 0.4 G

> Dipolarity ~ 0.6 l

> > Tilt ~ 10°

Archetype of a terrestrial planet!



<u>Earth</u>

Direct measurements of Earth's magnetic field date back to the early 1500's, with a boost in the early 1800's with the Magnetic Crusade led by Sabine in England and Gauss and Weber in Germany

Today we also have satellite measurements



Magnetometer used by Alexander von Humboldt in his Latin America expedition of 1799-1804



Longer time history can be inferred from measurements of magnetic signatures in crustal rocks



<u>Earth</u>



Mantle convection responsible for plate tectonics but not the geodynamo





Mantle non-conducting, slow

Overturning time ~100 million years





Outer Core conducting, fast

Overturning time ~500 years



Rotational influence quantified by

Rossby number

$$\operatorname{Ro} = \frac{U}{2\Omega D} = \frac{1}{4\pi} \; \frac{P_{rot}}{\tau_c}$$



$$\mathrm{Ro} \sim 4 \times 10^{-7}$$

Outer Core conducting, fast

Overturning time ~500 years



Spherical Harmonic expansion of the surface field allows for a backward extrapolation to the core-mantle boundary (CMB)

Assuming no currents in the non-conducting mantle & crust

$$B_r \propto r^{-(\ell+2)}$$



R. Townshend (Wisconsin)

Jones (2011)

<u>Earth</u>

$$B_r \propto r^{-(\ell+2)}$$

Dipole dominates at large distances from the dynamo region ~ r⁻³



Time evolution of surface field can be used to infer flows at the CMB

<u>Earth</u>

- n Energy sources for convective motions
 - Outward heat transport by conduction
 - Cooling of the core over time
 - Proportional to the heat capacity

Latent heat

 Associated with the freezing (phase change) of iron onto the solid core

Gravitational Differentiation

 Redistribution of light and heavy elements, releasing gravitational potential energy

Radioactive Heating

Intersection Energy released by the decay of heavy elements



No Dynamo

No field detected



Core may be liquid and conducting, but it may not be convecting (rigid top may inhibit cooling)



Also - slow rotation

Mars

No Dynamo

Fields patchy, reaching ~ 0.01 G in spots but no dipole

Why?

It had a dynamo in the past (remnant crustal magnetism) but it cooled off fast, freezing out its molten core

Mercury

Dynamo!

Field strength ~ 0.003 G

> Dipolarity ~ 0.71 G

Tilt ~ 3°

Huge iron core relative to size of planet that is still partially molten

But we're still not really sure what's going on!

NASA/ESA

Jupiter

Big Whopping Dynamo!

Field strength ~ 7 G

Dipolarity ~ 0.6 l

Tilt ~ 10°

Jupiter

Jupiter: Internal Structure

French et al (2012)

Jupiter: Internal Structure

French et al. (2012)

Jupiter: Magnetic Field (Pre-Juno)

Initial results from Juno

Stronger and more patchy than expected (higher-order multipoles)

This suggests that dynamo action might exist closer to the surface than previously thought

 $B_r \propto r^{-(\ell+2)}$

Moore et al (2017)

Cowling's Theorem

Why is this a surprise?

0

Assume B is axisymmetric and consider the longitudinally-averaged MHD induction equation:

$$\frac{\partial \mathbf{B}}{\partial t} = \mathbf{\nabla} \times \left(\langle \mathbf{v} \rangle \times \mathbf{B} - \langle \eta \rangle \, \mathbf{\nabla} \times \mathbf{B} \right)$$

Express B as
$$\mathbf{B} = \mathbf{
abla} imes \left(A \hat{oldsymbol{\phi}}
ight) + B \hat{oldsymbol{\phi}}$$

Evolution eqn for A (after some manipulation)

 $\lambda = r\sin\theta$

$$\frac{\partial}{\partial t} \left(\lambda A \right) = -\mathbf{v} \cdot \nabla \left(\lambda A \right) + \eta \lambda \left(\nabla^2 A - \lambda^{-2} \right) A$$

Multiply by λA and integrate over volume: if $\nabla \cdot v = 0$ then the first term on the RHS is zero and the second term is negative

Cowling's Theorem (cont.)

$$\frac{\partial}{\partial t} \left(\lambda A \right) = -\mathbf{v} \cdot \nabla \left(\lambda A \right) + \eta \lambda \left(\nabla^2 A - \lambda^{-2} \right) A$$

A decays with time

If A decays with time, then B will decay with time too (Work it out!)

Even if $\nabla \cdot v \neq 0$ you can show that a steady field ($\partial A/\partial t = 0$) cannot be maintained

Conclusion: it is not possible to sustain a steady axisymmetric B field against ohmic dissipation

Corollary: It is not possible for a dynamo to produce a steady axisymmetric field!!

Understanding the Dynamics

Conservation of momentum in MHD

But rotation exerts an overwhelming influence Coriolis accelerations happen quickly (days) compared to convection and dynamo time scales (hundreds to thousands of years)

$$\operatorname{Ro} = \frac{U}{2\Omega D} << 1 \qquad \qquad \operatorname{Ek} = \frac{\nu}{2\Omega D^2} << 1$$

$$\rho \frac{\partial \mathbf{v}}{\partial t} = -\left(\rho \mathbf{v} \cdot \boldsymbol{\nabla}\right) \mathbf{v} - 2\rho \left(\boldsymbol{\Omega} \times \mathbf{v}\right) - \boldsymbol{\nabla} P + \rho \mathbf{g} + c^{-1} \mathbf{J} \times \mathbf{B} - \boldsymbol{\nabla} \cdot \boldsymbol{\mathcal{D}}$$

Result: Flows evolve quasi-statically in so-called Magnetostrophic (MAC) Balance

$$c^{-1}\mathbf{J} \times \mathbf{B} \approx 2\rho \left(\mathbf{\Omega} \times \mathbf{v}\right) + \mathbf{\nabla} P - \rho \mathbf{g}$$

Conservation of mass

Anelastic approximation (valid for small Ma)

Boussinesq approximation (valid for small Ma, $H_{\rho} >> D$)

hydrostatic background

$$\boldsymbol{\nabla} \boldsymbol{\cdot} \left(\hat{\rho} \mathbf{v} \right) = 0$$

$$\nabla \cdot \mathbf{v} = 0$$

$$c^{-1}\mathbf{J} \times \mathbf{B} \approx 2\rho \left(\mathbf{\Omega} \times \mathbf{v}\right) + \mathbf{\nabla} P - \rho \mathbf{g}$$

Now set B = 0 and assume that $\nabla \rho$ is mainly radial

Then the ϕ component of the curl gives (anelastic approximation):

$$oldsymbol{\Omega} \cdot oldsymbol{
abla} (
ho \mathbf{v}) = rac{\partial}{\partial z} \left(
ho \mathbf{v}
ight) = 0$$
 Taylor-Proudman Theorem

Boussinesq version:

$$\frac{\partial \mathbf{v}}{\partial z} = 0$$

Rapidly rotating flows tend to align with the rotation axis

Work with a partner to draw what you think <u>convective</u> motions might look like in a rapidly-rotating spherical shell

How can you get the heat out while still satisfying the Taylor-Proudman theorem

$$\frac{\partial \mathbf{v}}{\partial z} = 0$$

Linear Theory

The most unstable convective modes in a rapidly-rotating, weakly-stratified shell are **Busse columns** aka **Banana Cells**

The preferred longitudinal wavenumber (m) scales as

Ek^{-1/3}

Coriolis vs viscous diffusion

Linear Theory

The **Tangent Cylinder**

Delineates two distinct dynamical regimes

Implication of the Taylor-Proudman theorem

Linear Theory: Traveling Waves

Prograde propagation

(thermal Rossby waves)

Induced by curvature of outer boundary and/or density stratification

 $\frac{\omega_z}{H} = \text{constant}$

Simplest example: Boussinesq fluid, centrifugal gravity, local, linear perturbations, small boundary curvature (Busse 2002)

$$v_p = \frac{4\Omega}{L} \frac{\tan \chi}{(1+P_r)(k_y^2 + k_x^2)}$$

Nonlinear Regimes require Numerical Models

Solve the MHD equations in a rotating spherical shell Anelastic or Boussinesq approximation ρ, T, P, S are linear perturbations about a <u>hydrostatic</u>, <u>adiabatic</u> background state

Convection simulations: heating from below, cooling from above

Axial alignment persists even in turbulent parameter regimes

Kageyama et al (2008)

Axial vorticity $\omega \cdot \Omega$

 $Ek = 2.3 \times 10^{-7}$ $Ek = 2.6 \times 10^{-6}$

Busse columns give way to vortex sheets but the flow is still approximately 2D

 $Ek = \frac{\nu}{20 P^2}$

General trends

Complexity of magnetic field depends mainly on the rotational influence

Rapid rotators tend to be more dipolar

Christensen & Aubert (2006)

$$c^{-1}\mathbf{J} \times \mathbf{B} \approx 2\rho \left(\mathbf{\Omega} \times \mathbf{v}\right) + \mathbf{\nabla} P - \rho \mathbf{g}$$

Assuming MAC balance, compute the ratio of ME/KE How does it scale with Ro?

$$abla imes \mathbf{B} = rac{4\pi}{c} \mathbf{J}$$

$$ME = rac{B^2}{8\pi}$$

$$Ro = rac{U}{2\Omega D}$$

$$KE = rac{1}{2}\rho v^2$$

 $c^{-1}\mathbf{J} \times \mathbf{B} \approx 2\rho \left(\mathbf{\Omega} \times \mathbf{v} \right) + \mathbf{\nabla} P - \rho \mathbf{g}$

Assuming MAC balance, compute the ratio of ME/KE How does it scale with Ro?

$$\frac{ME}{KE} \sim \mathrm{Ro}^{-1}$$

>>1 *if* Ro << 1!

But how do KE and Ro (and thus, ME) depend on observable^{*} global parameters like Ω and F_c?

*in principle

General trends

Field strength scales with the heat flux through the shell (independent of Ω!)

Rapid rotators seem to operate at maximum efficiency, tapping all the energy they can

General trends

Rapid rotators seem to operate at maximum efficiency, tapping all the energy they can

This may apply to rapidly-rotating stars as well as planets!

	Earth	Jupiter	Simulations
Ra	10 ³¹	10 ³⁷	10⁶-10 ⁷
Ek	3×10 ⁻¹⁵	10 -9	10 ⁻⁶ - 10 ⁻⁷
Rm	300-1000	400-3×10 ⁴	50-3000
Pm	5-6 ×10 ⁻⁷	6×10 ⁻⁷	0.1-0.01

Numerical Models: The Hope

Realistic simulations might be possible if you can achieve the right dynamical balances (e.g. MAC balance)

- n The most important parameters to get right (or as right as possible)
 - Ro
 - Appropriate rotational influence on the convection
 - ► Rm
 - Reasonable estimate of the ohmic dissipation

▶ Ek

 At least get it small enough that viscosity isn't part of the force balance

Example: The Geodynamo

Points of comparison: Field strength, morphology (spectrum, symmetry, etc), Reversal timescale

Christensen et al (2010) Best matches are those with Ek < 10⁻⁴ and Rm "large enough"

Example: The Geodynamo

But be careful! They could be right for the wrong reasons! For example, both c and d have a higher Ra and lower Ek than b they <u>should</u> be more realistic, right?

Example: The Geodynamo

Coupling to inner core needed to get the reversal time scale right (Glatzmaier & Roberts 1995; Glatzmaier et al 1999)

But "dipole solutions are not easy to find" for the "best" parameters

Stanley & Glatzmaier (2009)

So what's going on with Saturn?

Maybe the field is "axisymmeterized" by an overlying stable layer that has differential rotation but no convection (Stevenson 1982, Stanley 2010)

Or, maybe it's running a different type of dynamo, driven more by shear than buoyancy (Cao et al 2012)

Numerical Models: Summary

n Lessons Learned

- Rapid Rotation has a profound influence on the dynamics
- Success attributed to correct dynamical balances and (when possible) realistic Rm

n Future challenges

- What happens at <u>really</u> low Ek (tiny v)?
- Peculiarities of particular planets (Saturn, Mercury, Uranus, Neptune...)
 - Boundary conditions (adjacent layers)
 - Rapid variations of η
 - Energy sources
 - Compositional convection
- Moving to more realistic parameters doesn't always improve the fidelity of the model
- Exoplanets!

Featherstone & Heimpel 2017

