

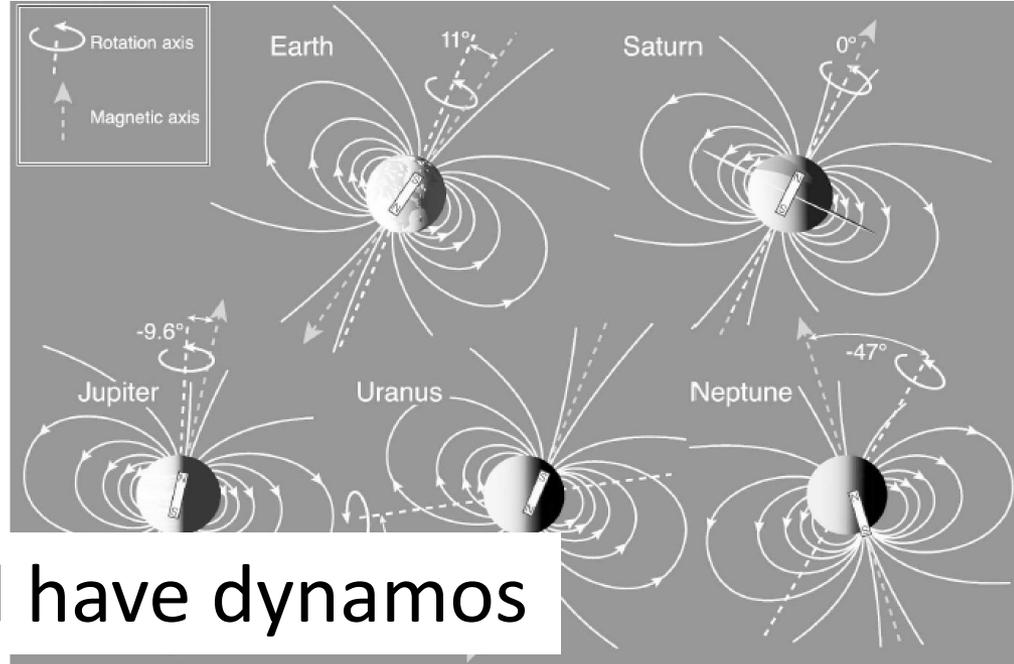
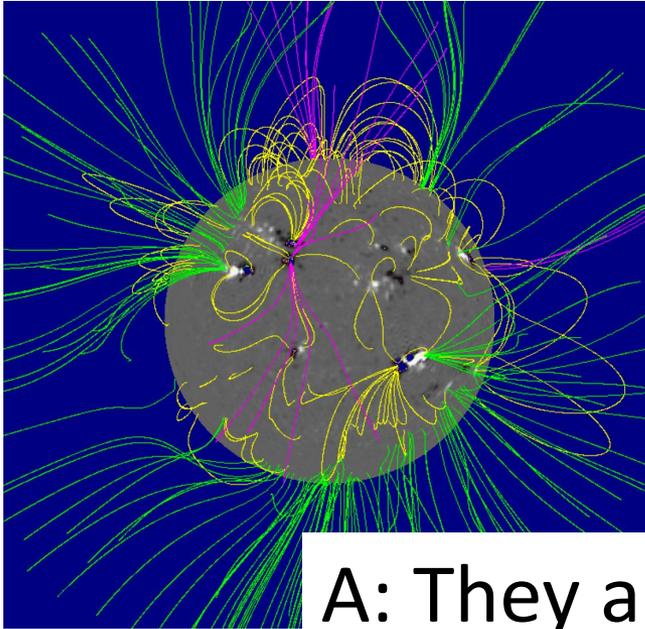
# Q: Why do the Sun and planets have magnetic fields?

Dana Longcope

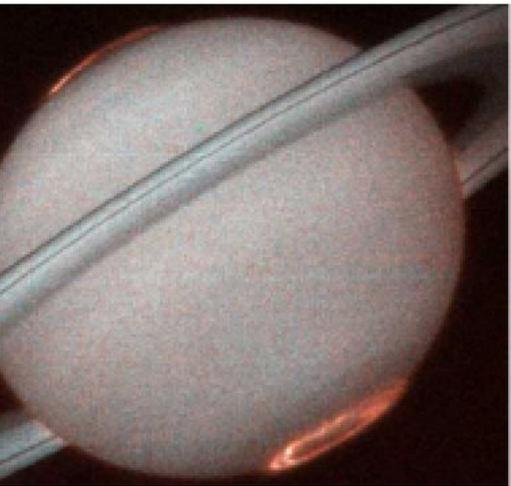
Montana State University

w/ liberal “borrowing” from Bagenal,  
Stanley, Christensen, Schrijver,  
Charbonneau, ...

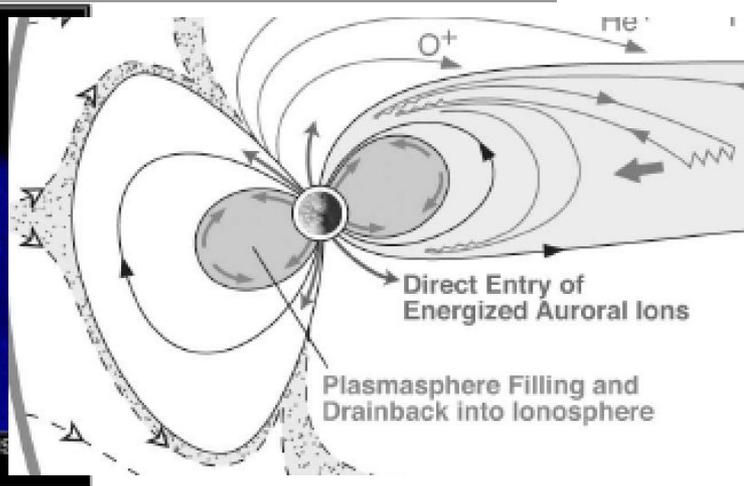
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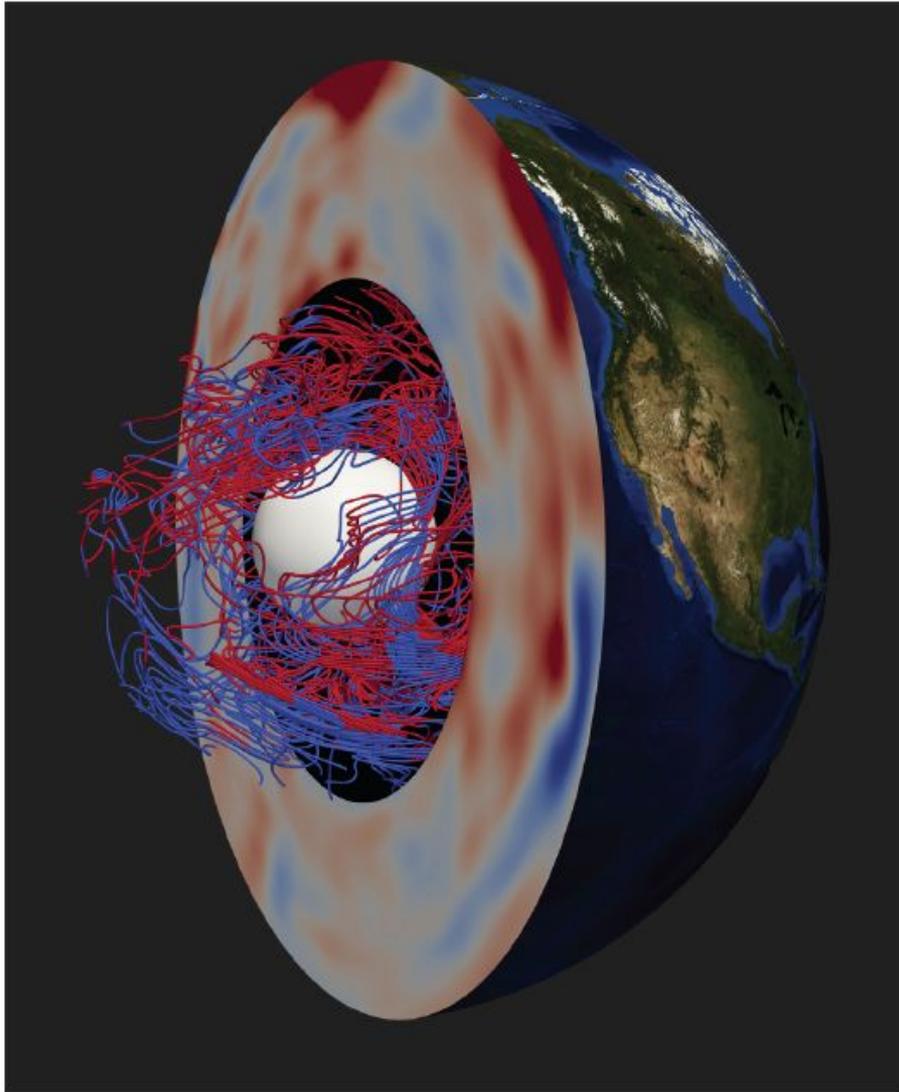
A: They all have dynamos



Jupiter Aurora  
NASA and J. Clarke (University of Michigan) • STScI-PRC00-38



# DYNAMO INGREDIENTS



(1) electrically conducting fluid

- Plasma (stars)
- Liquid iron (terrestrial planets)
- Metallic hydrogen (gas giants)
- Ionized water (ice giants)

(2) fluid must have complex motions

- Complex turbulent flows
- Rotation: breaks mirror-symmetry  
not required, but needed for large-scale, organized fields

(3) motions must be vigorous enough

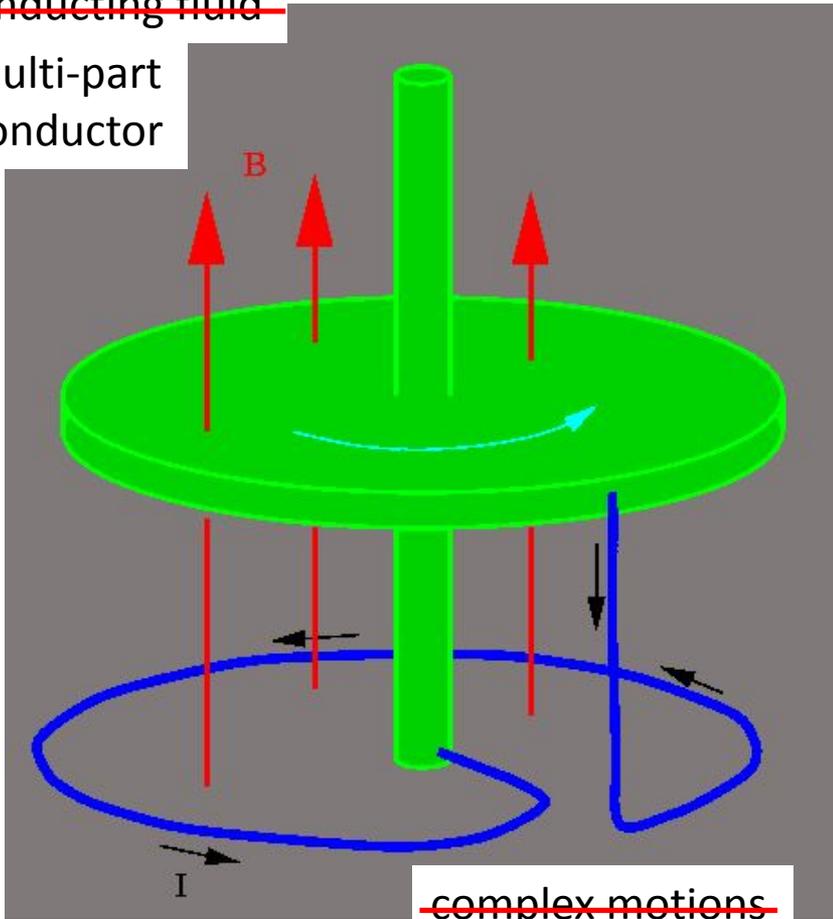
- Figure of merit: Magnetic Reynold's #

$$Rm = \text{velocity} \times \text{size} \times \text{conductivity}$$

# A Toy w/ all ingredients

~~conducting fluid~~

multi-part  
conductor



~~complex motions~~

lack of mirror  
symmetry

differential  
motion of parts

$$V_{\text{disk}} = \int_0^{\square} vB dr = \int_0^{\square} (\Omega r) B dr = \frac{1}{2\pi} \Omega \Phi_{\text{disk}}$$

$$IR = \underbrace{\frac{\Omega}{2\pi} \Phi_{\text{disk}}}_{\text{motional EMF}} - \underbrace{L \frac{dI}{dt}}_{\text{back EMF}} = \frac{\Omega}{2\pi} M_{w,d} I - L \frac{dI}{dt}$$

$$\frac{dI}{dt} = \left( \underbrace{\frac{\Omega}{2\pi} \frac{M_{w,d}}{L}}_{\text{generation}} - \underbrace{\frac{R}{L}}_{\text{dissipation}} \right) I = \gamma I \quad I(t) = I_0 e^{\gamma t}$$

$$\text{Growth: } \gamma > 0 \Leftrightarrow \Omega > 2\pi \frac{R}{M_{w,d}}$$

$$\frac{v}{\square} > 2\pi \frac{1/\sigma \square}{\mu_0 \square} = \frac{2\pi}{\mu_0 \sigma \square^2}$$

Growth:  $Rm = \mu_0 \sigma \square v > 2\pi$

# Reality: Conducting fluid – MHD

If  $\mathbf{B}$  is weak: kinematic equations

Fluid dynamics

$$\left\{ \begin{array}{l} \frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \mathbf{v}) \\ \rho \frac{\partial \mathbf{v}}{\partial t} + \rho (\mathbf{v} \cdot \nabla) \mathbf{v} = -\nabla p + \rho \mathbf{g} + \nabla \cdot \vec{\sigma} \\ \rho c_v \left( \frac{\partial T}{\partial t} + \mathbf{v} \cdot \nabla T \right) = -\frac{2}{3} T \nabla \cdot \mathbf{v} + \nabla \mathbf{v} : \vec{\sigma} + \dot{Q} \end{array} \right.$$

Traditional  
(neutral) fluid –  
solve first

Faraday's  
+ Ohm's laws

$$\frac{\partial \mathbf{B}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{B} = (\mathbf{B} \cdot \nabla) \mathbf{v} - \mathbf{B} (\nabla \cdot \mathbf{v}) + \eta \nabla^2 \mathbf{B}$$

Linear equation for  $\mathbf{B}(\mathbf{x}, t)$  – solve w/ known  $\mathbf{v}(\mathbf{x}, t)$

# Dynamo action in MHD

$$\underbrace{\frac{D\mathbf{B}}{Dt}}_{\frac{\partial\mathbf{B}}{\partial t} + (\mathbf{v} \cdot \nabla)\mathbf{B}} \quad \underbrace{\mathbf{B} \cdot [\nabla\mathbf{v} - \mathbf{I}(\nabla \cdot \mathbf{v})]}_{(\mathbf{B} \cdot \nabla)\mathbf{v} - \mathbf{B}(\nabla \cdot \mathbf{v})} = \mathbf{B} \cdot \mathbf{M}$$

$$\frac{\partial\mathbf{B}}{\partial t} + (\mathbf{v} \cdot \nabla)\mathbf{B} = (\mathbf{B} \cdot \nabla)\mathbf{v} - \mathbf{B}(\nabla \cdot \mathbf{v}) + \eta\nabla^2\mathbf{B}$$

If  $\mathbf{M}$  has a positive eigenvalue  $\lambda > 0$

$\mathbf{B}$  can grow exponentially: **DYNAMO ACTION**

- $\mathbf{B} \square -\mathbf{B}$  : same e-vector  $\square$  same  $\lambda$
- Reverse velocity AND reflect in mirror  $\square \lambda \square \lambda$
- Do one and not the other  $\square \lambda \square -\lambda$

$$\gamma \sim \frac{v}{\square} - \frac{\eta}{\square^2} = \frac{v}{\square} \left( 1 - \frac{\eta}{\square v} \right)$$

λ η∇²

Growth:  $Rm = \frac{\square v}{\eta} = \mu_0 \square v \sigma > 1$

# Q: What kind of flow has $\lambda > 0$ ?

- Turbulent flows have pos. Lyapunov exponent:  $\lambda > 0$

- tend to stretch balls into strands 
- tend to amplify fields

- Conditions for turbulence:

- driving: e.g. Rayleigh-Taylor instability
- viscosity fights driving – must be small

$$Re = \frac{U L}{\nu} \gg 1$$

- Rotation can organize turbulence:

align stretching direction  $\square$  azimuthal

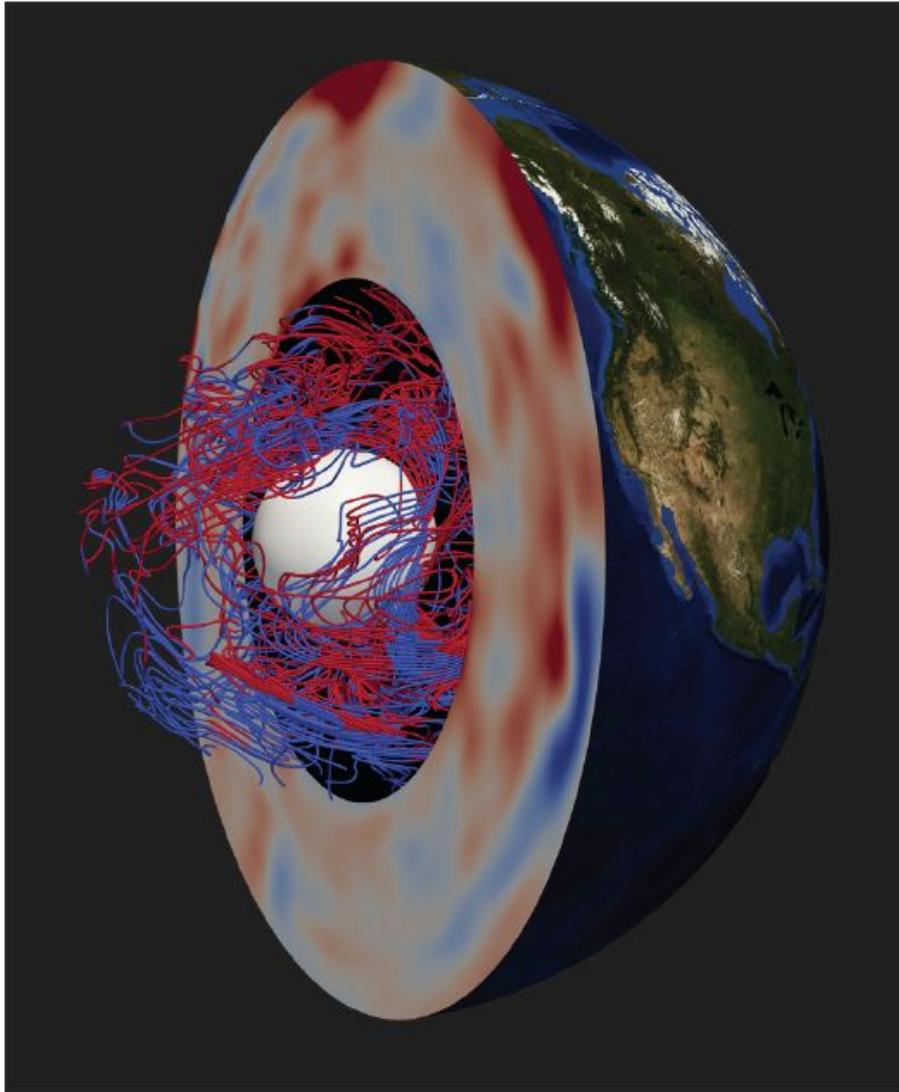
(toroidal) – known as  $\Omega$ -effect

– must be significant w.r.t. fluid motion

$$Ro = \frac{v}{\Omega L} \ll 1$$

	$\eta$ [m <sup>2</sup> /s]	$\nu$ [m <sup>2</sup> /s]	L [m]	v [m/s]	$\Omega$ [rad/s]	Rm	Re	Ro
Sun (CZ)	1	10 <sup>-2</sup>	10 <sup>8</sup>	1	10 <sup>-6</sup>	10 <sup>8</sup>	10 <sup>10</sup>	10 <sup>-2</sup>
Earth (core)	1	10 <sup>-5</sup>	10 <sup>6</sup>	10 <sup>-4</sup>	10 <sup>-4</sup>	10 <sup>2</sup>	10 <sup>7</sup>	10 <sup>-6</sup>

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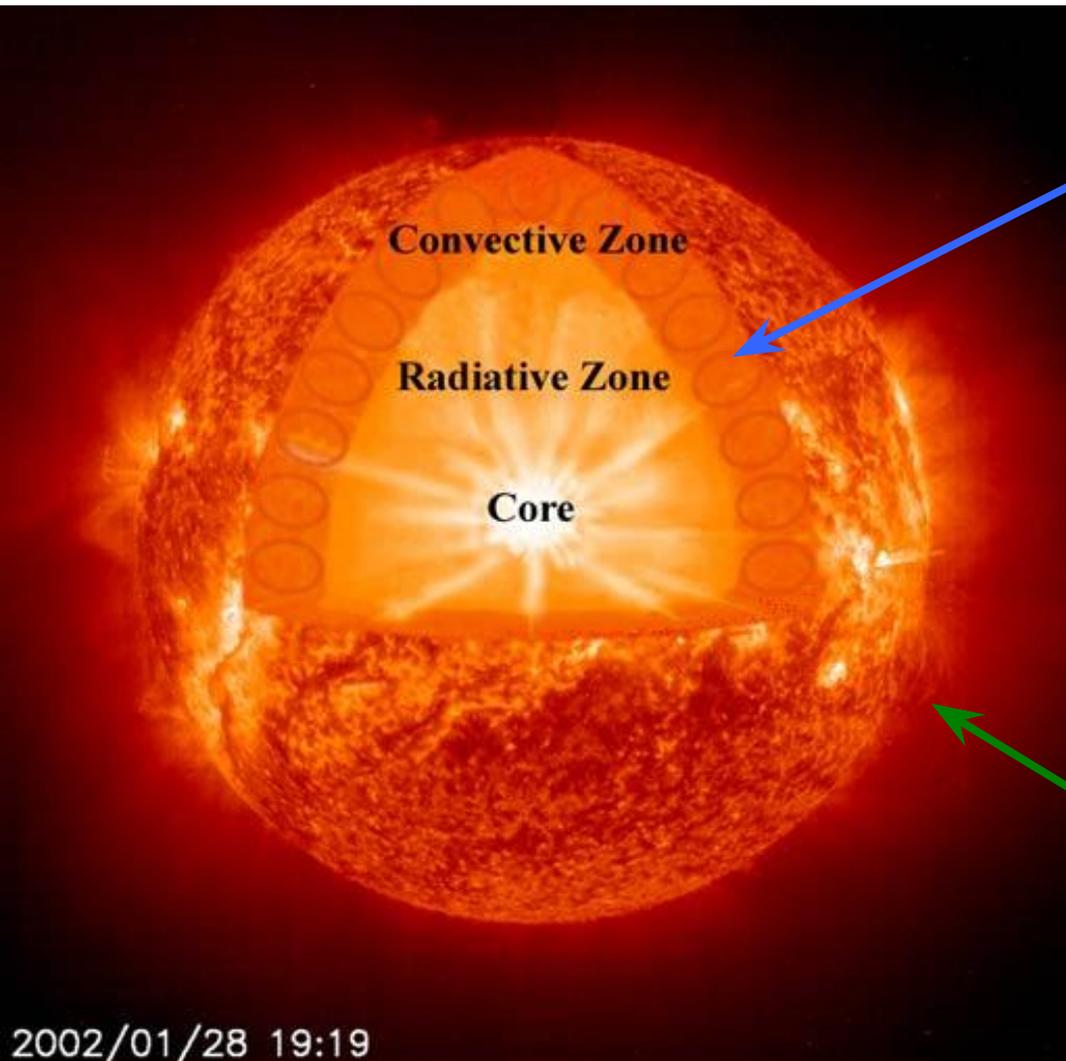
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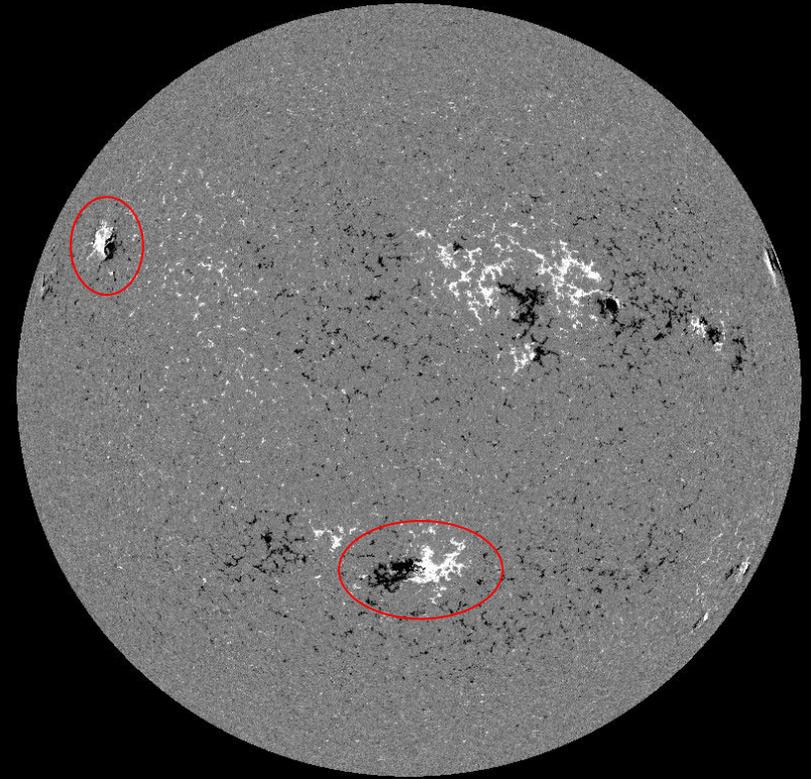
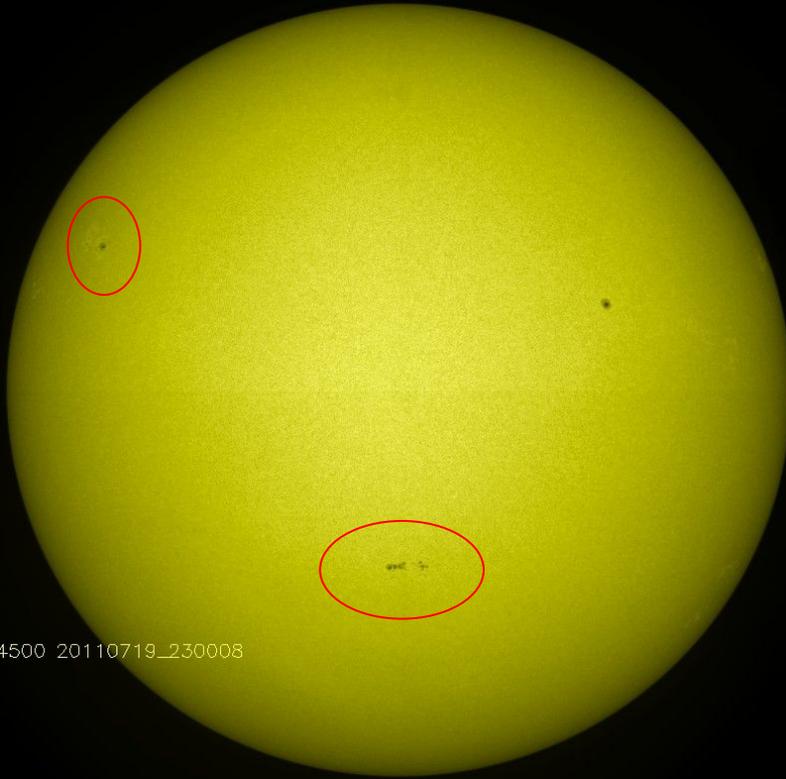
# How it works in the Sun



- Entire Star: H/He plasma
- Convection Zone (CZ)
  - Outer 200,000 km
  - Turbulence:  
 $Re = 10^{10}$
  - Thermally driven
  - Good conductor  
 $Rm = 10^8$
  - Rotation effective  
 $Ro = 10^{-2}$
- Corona – conductive but tenuous:  
J smaller (~0?)

# Evidence of the dynamo

Magnetic field where there are sunspots

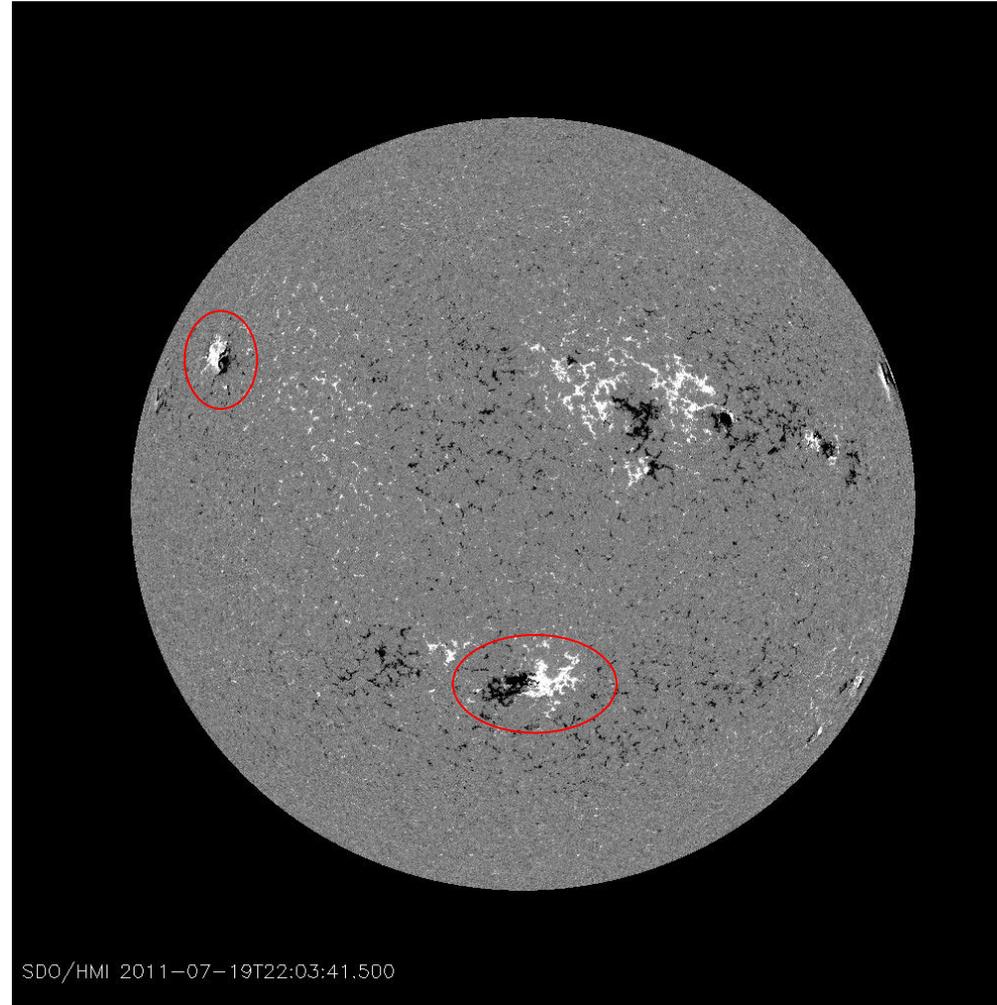
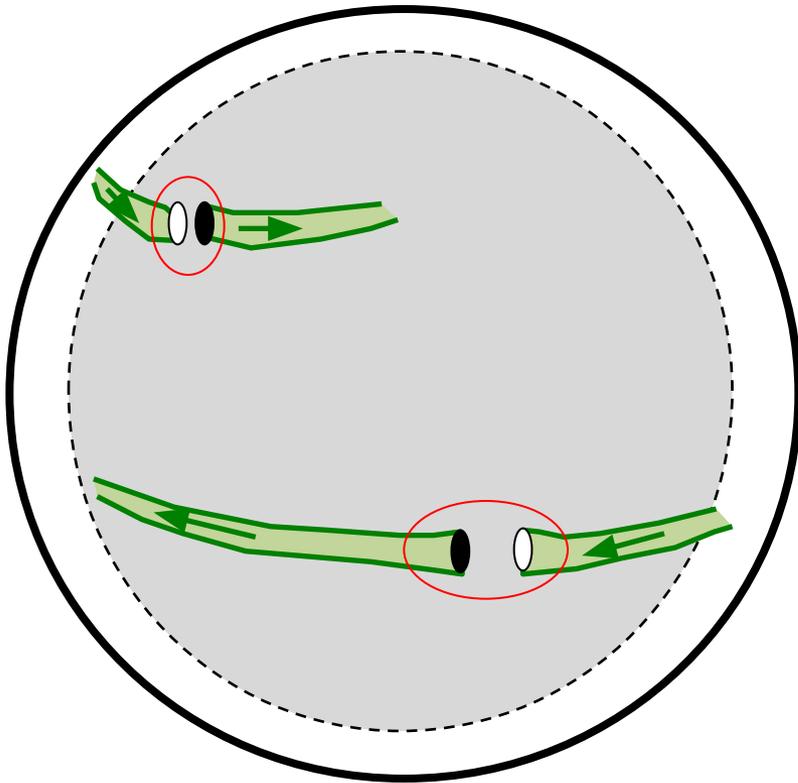


Field outside sunspots and elsewhere too

SDO/AIA-4500 20110719\_230008

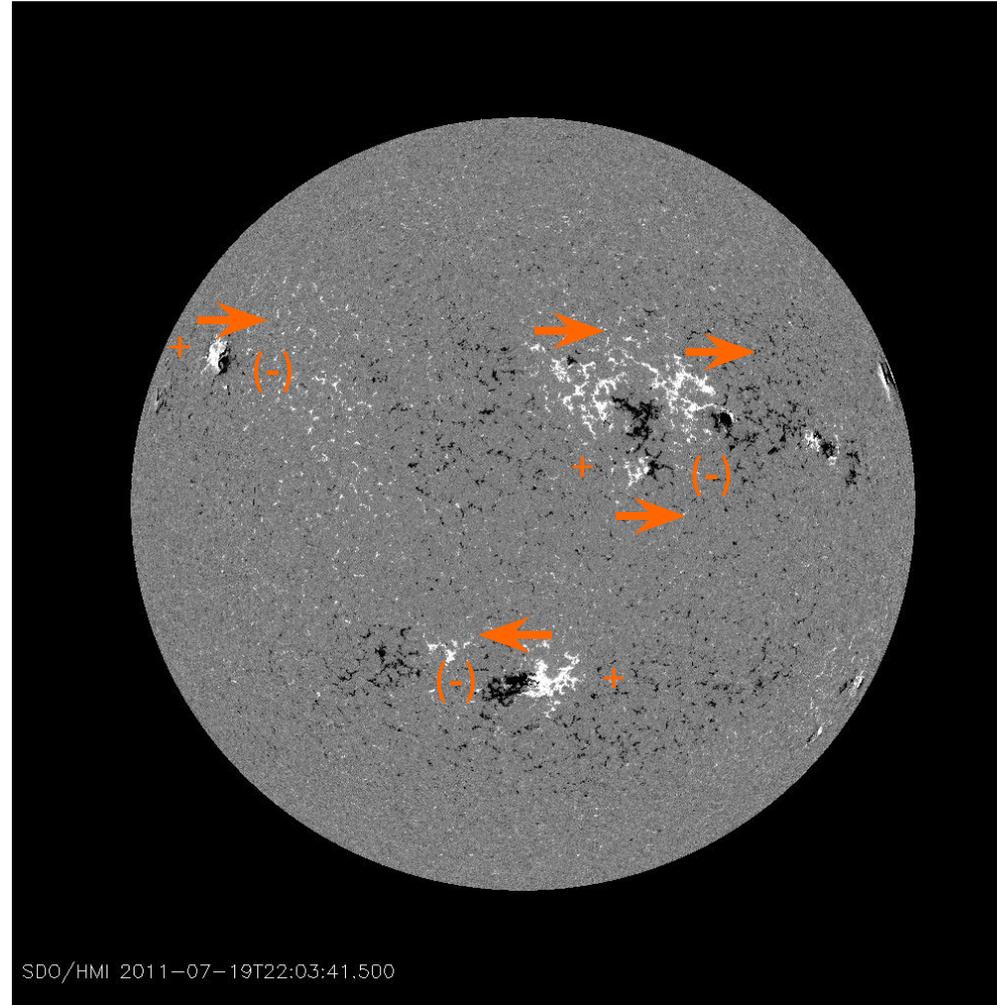
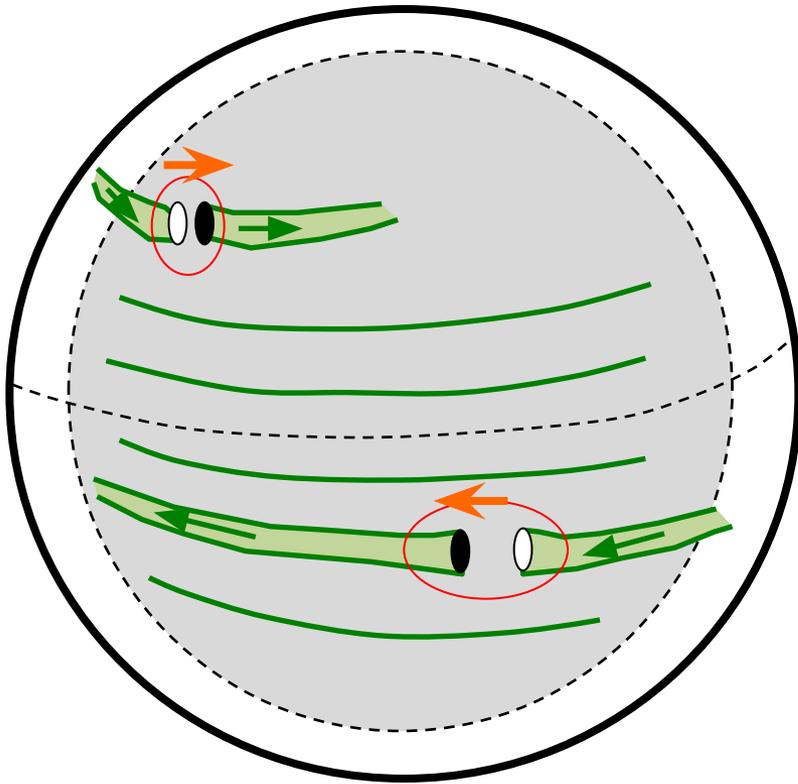
# Evidence of the dynamo

Field is **fibril**

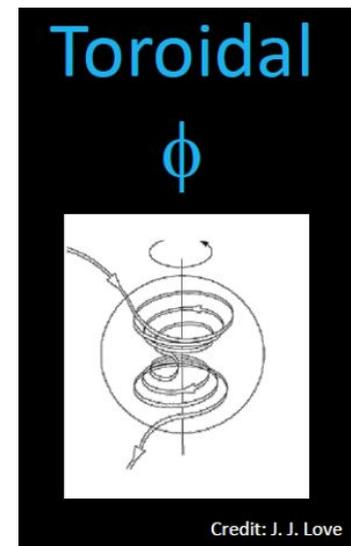
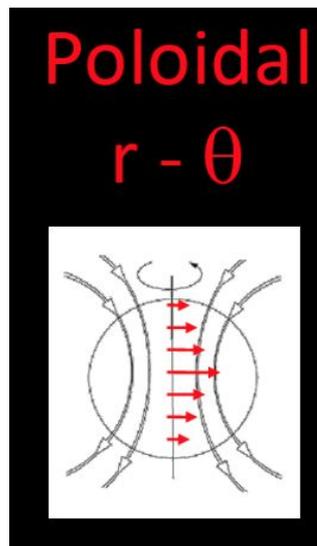


# Evidence of the dynamo

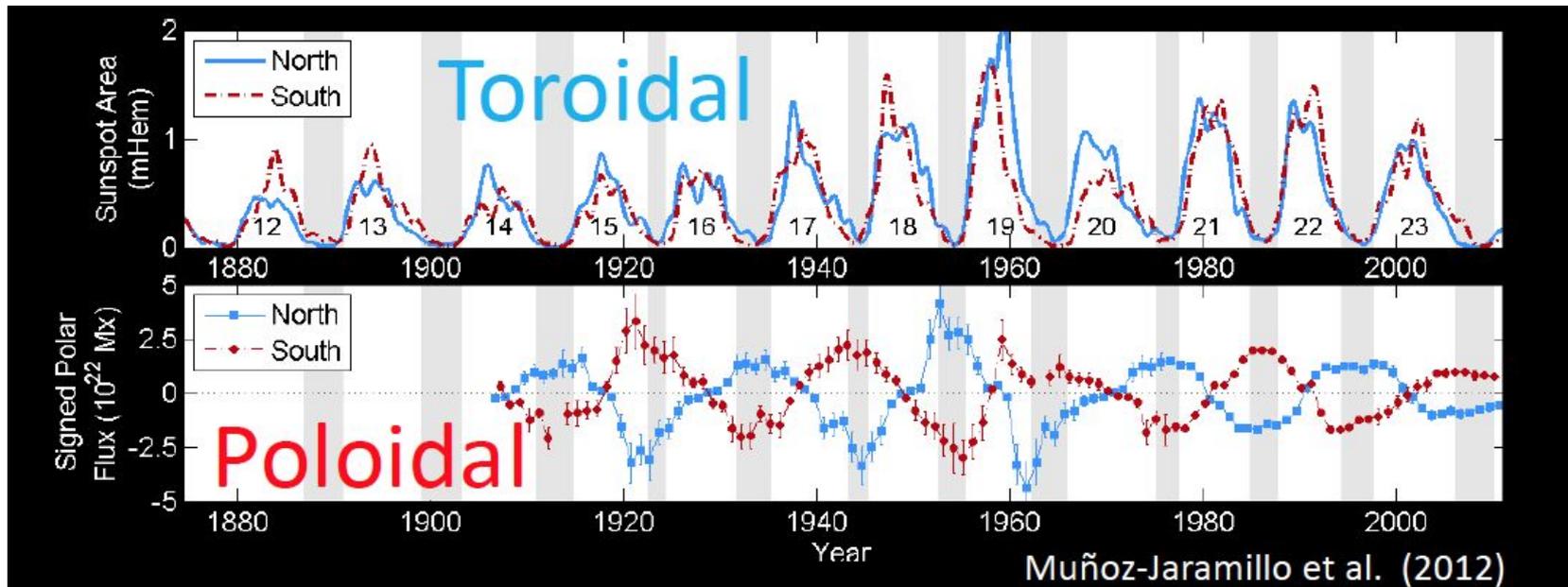
Field orientation: mostly  
**toroidal (E-W)**



An oscillator  
at work:



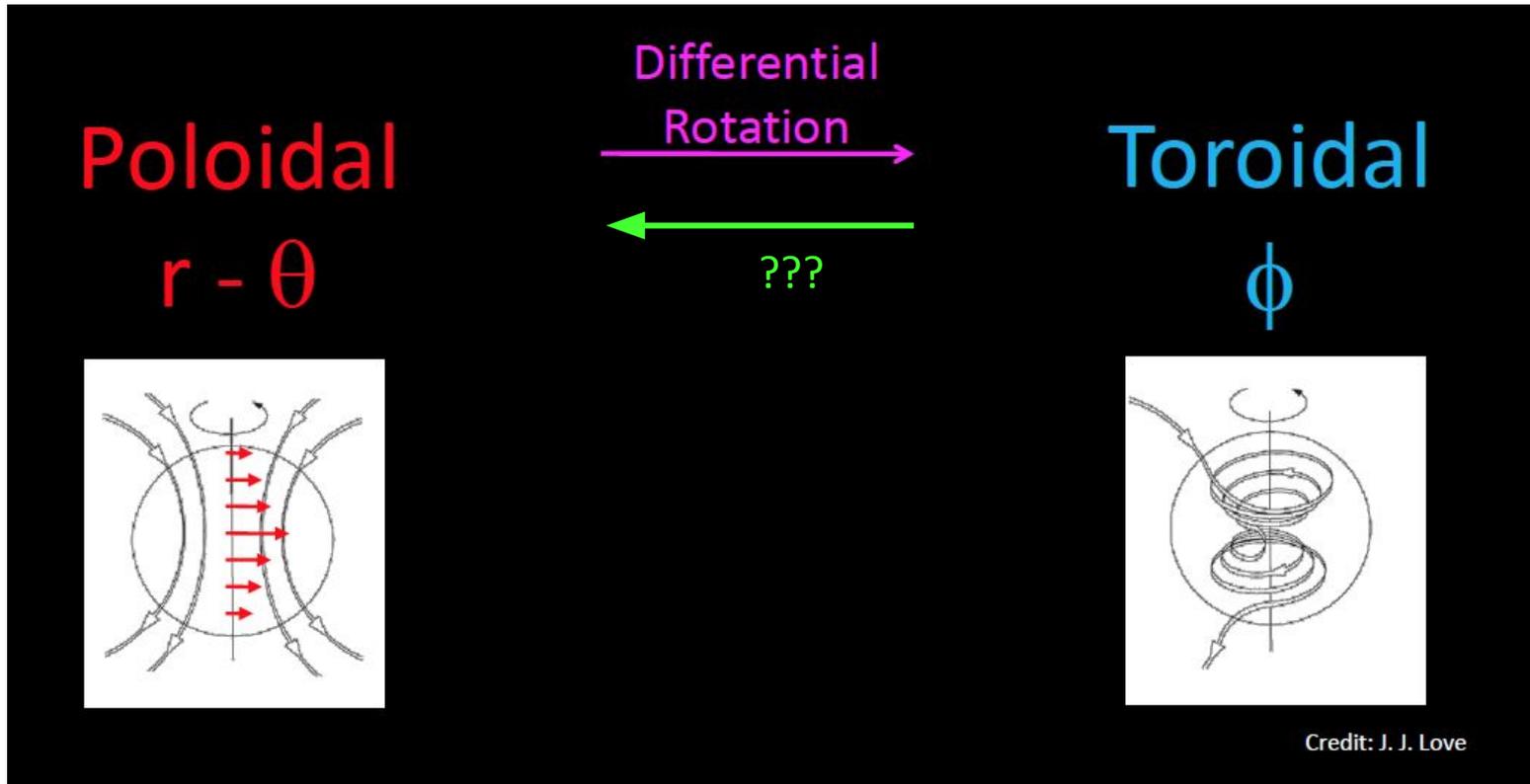
Credit: J. J. Love



Munoz-Jaramillo lecture (2015)

NB:  $B_{pol}$  &  $B_{tor}$  out of phase like Q & I in LC circuit (oscillator)

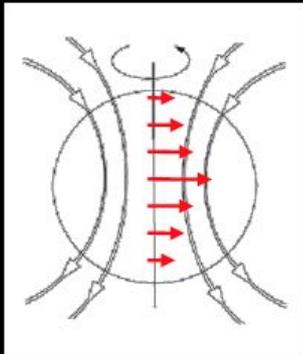
# Solar dynamo as an oscillator



Half of the  
oscillation...

Poloidal

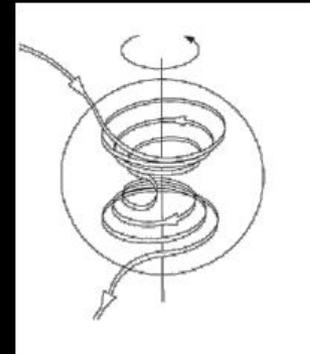
$r - \theta$



Differential  
Rotation →

Toroidal

$\phi$



Credit: J. J. Love

Activity 44: Estimate the time it takes for the solar equator to execute one more full rotation than the poles in the same time.

$$P = 25 \text{ d @ equator: } f = 1/25 = 0.040 \text{ d}^{-1}$$

$$P = 35 \text{ d @ poles: } f = 1/35 = 0.029 \text{ d}^{-1}$$

$$\Delta f = 0.011 \text{ d}^{-1} \quad \square \quad 1/\Delta f = 87 \text{ d}$$

Check:

$$N_{\text{eq}} = 87/25 = 3.5 \text{ rotations of equator}$$

$$N_{\text{pole}} = 87/35 = 2.5 \text{ rotations of equator}$$

Activity 47: If we take the Sun's polar field – averaging at cycle minimum at about 5 Gauss – how long would it take to wind that field into a strength of some  $10^5$  G? Hint: remember the field line stretch-and-fold from Fig. 4.6, look at the illustration in Fig. 4.9, and consider 'compound interest'. [i.e. exponentials]

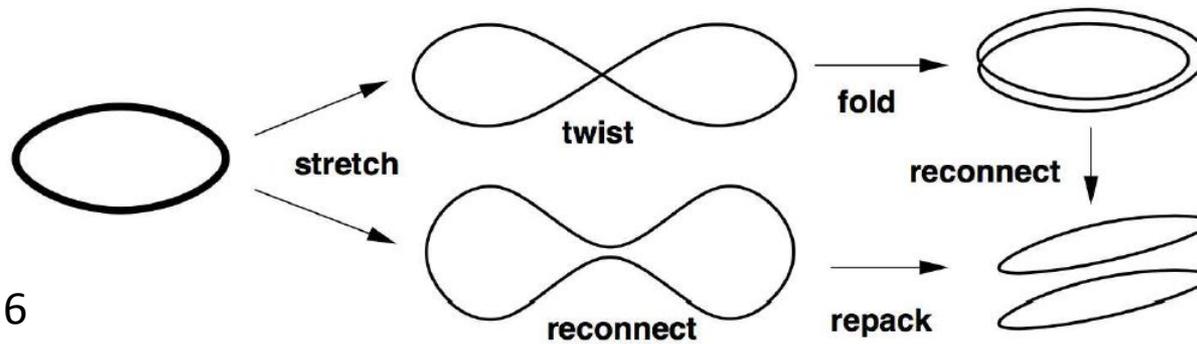


Fig. 4.6

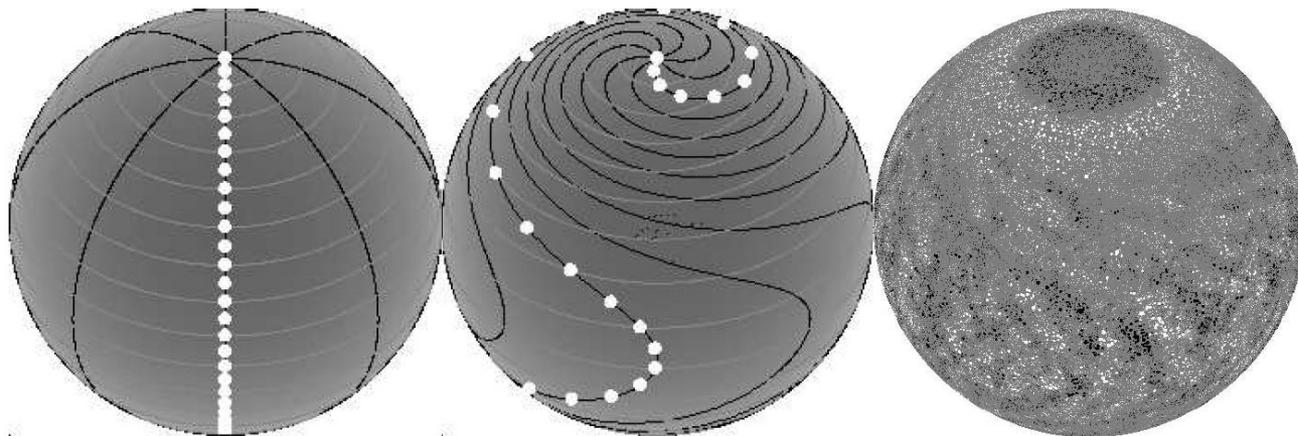


Fig. 4.9

Activity 47: p. 86: If we take the Sun's polar field – averaging at cycle minimum at about 5 Gauss – how long would it take to wind that field into a strength of some  $10^5$  G? Hint: remember the field line stretch-and-fold from Fig. 4.6, look at the illustration in Fig. 4.9, and consider 'compound interest'. [i.e. exponentials]

1 stretch+fold requires **87 days**

(one rotation of pole w.r.t. equator)

'compound interest': Field strength **doubles** every 87 days

□ e-folding time =  $87/\ln(2) = 125$  days

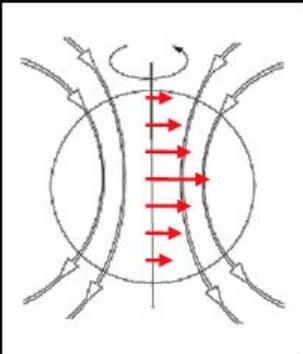
$$B = 5 \text{ G} \times \exp( t/125 ) = 10^5 \text{ G}$$

$$t = 125 \text{ d} \times \ln( 10^5/5 ) = 1,240 \text{ days} = 3.4 \text{ years}$$

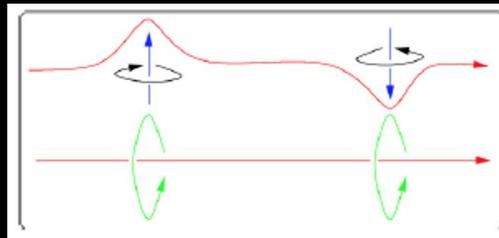
# The other half ... mean field theory

(see lecture by Prof. Bhattacharjee)

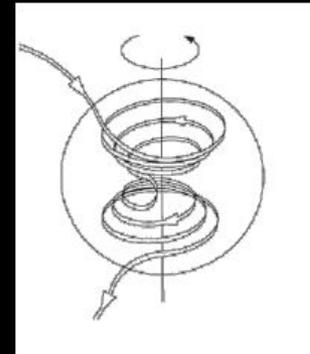
**Poloidal**  
 $r - \theta$



Differential  
Rotation  
 $\alpha$  effect



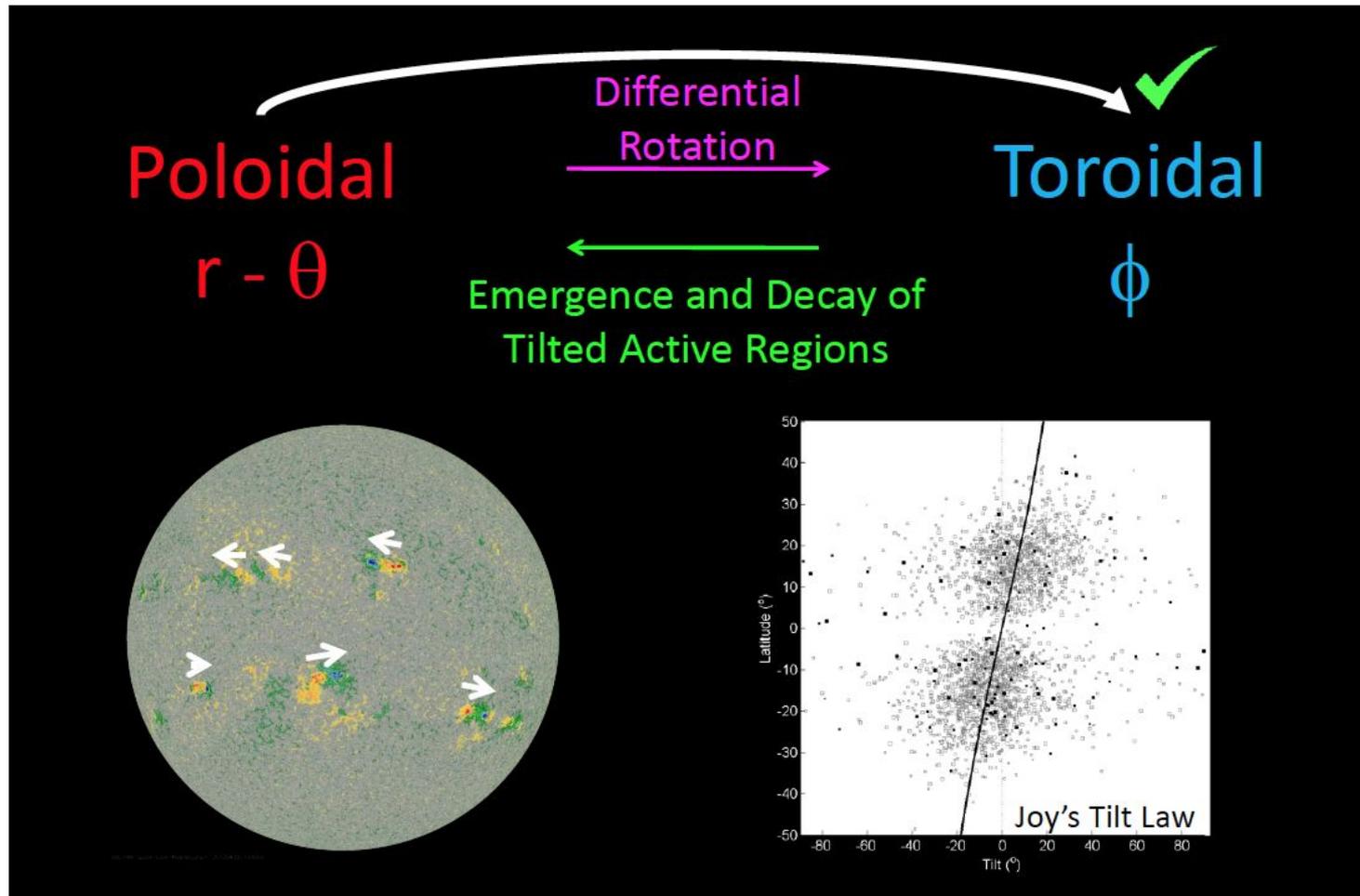
**Toroidal**  
 $\phi$



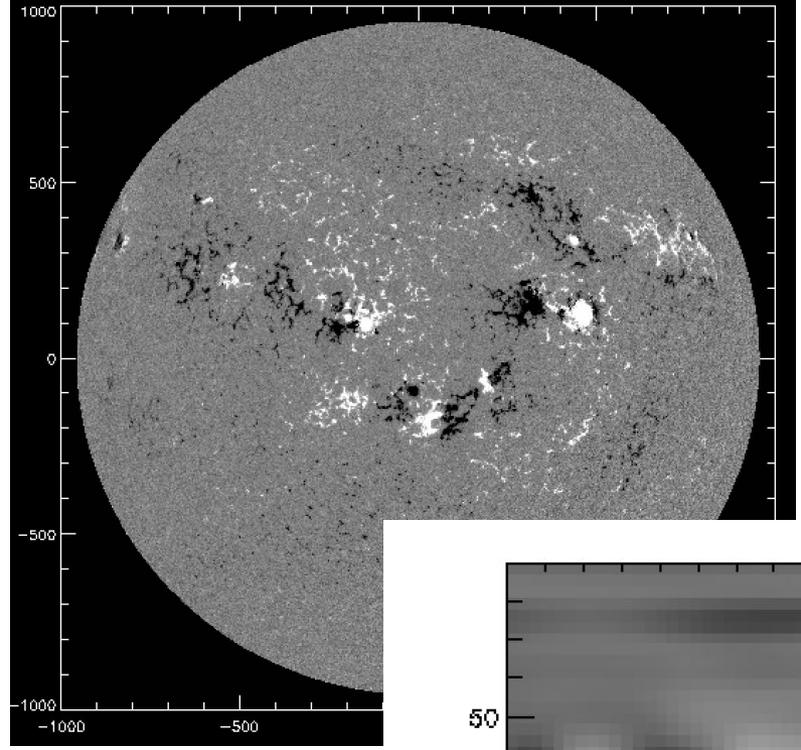
Credit: J. J. Love

# The other half...

## Babcock-Leighton model

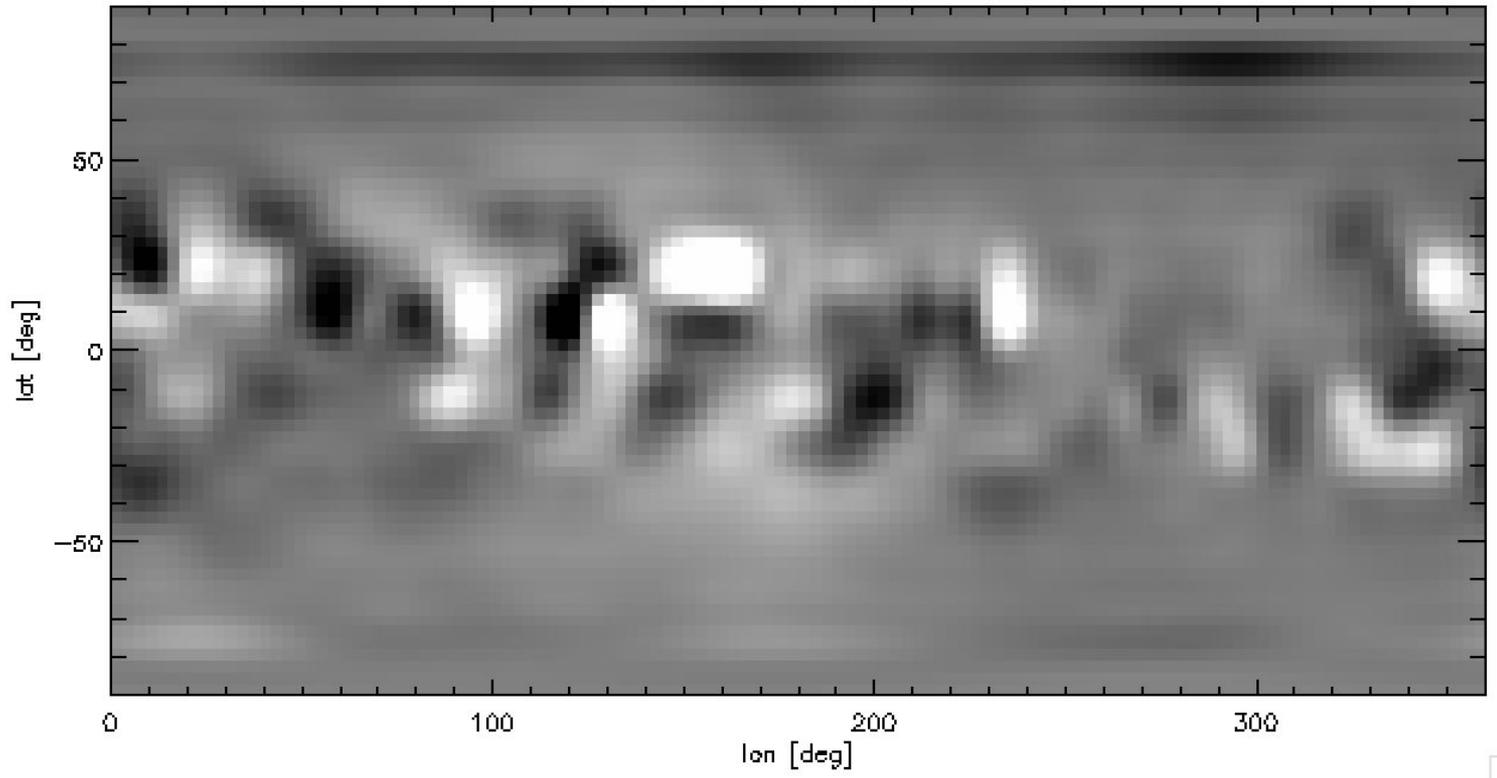


2001-05-26T00:00:02.169Z



Synoptic plot: unwrapped  
view built up over time

Sun @ 2001-05-19T20:26:15.000Z

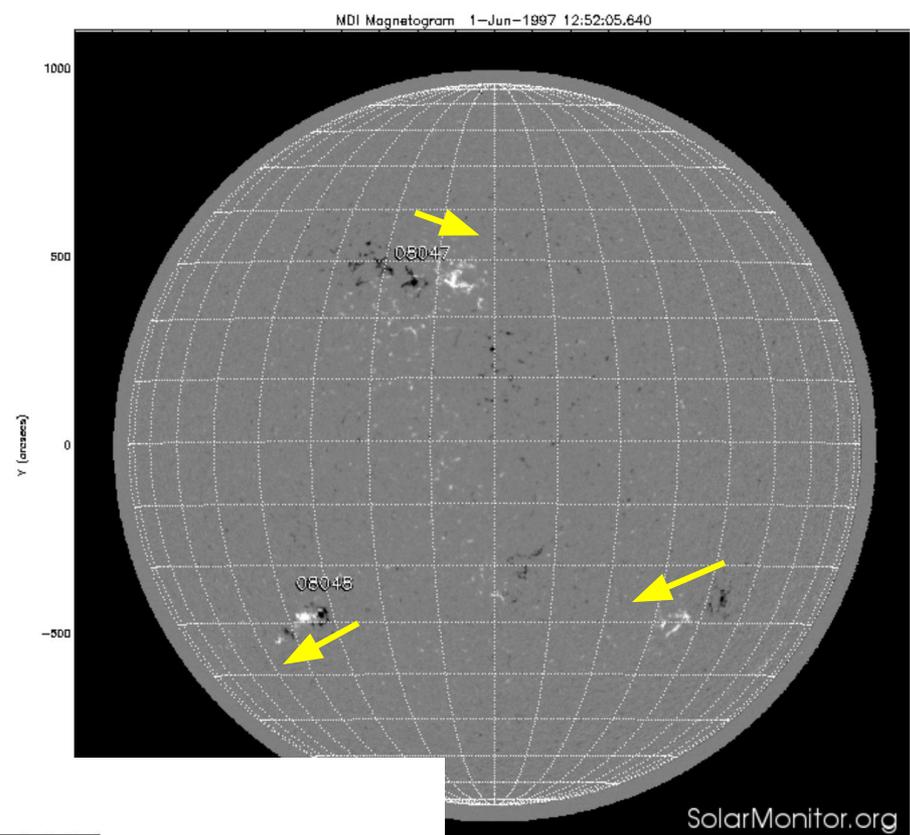


Tilted active regions contribute to magnetic dipole moment – in sense to reverse it

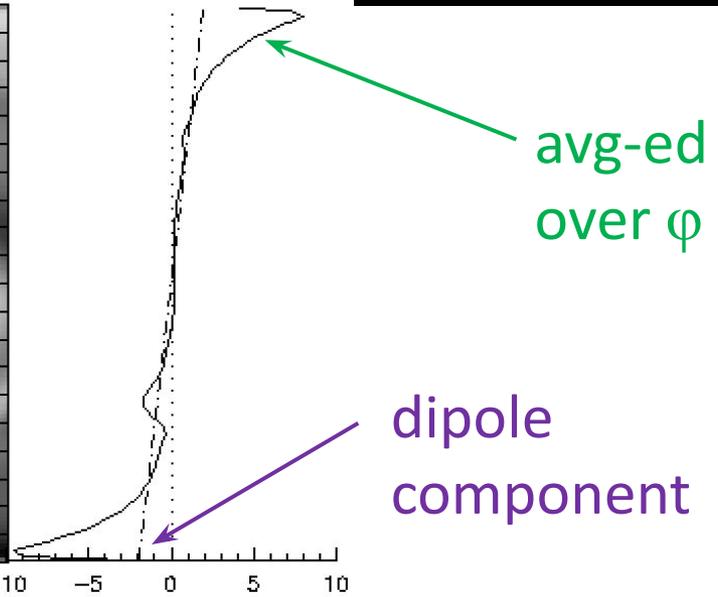
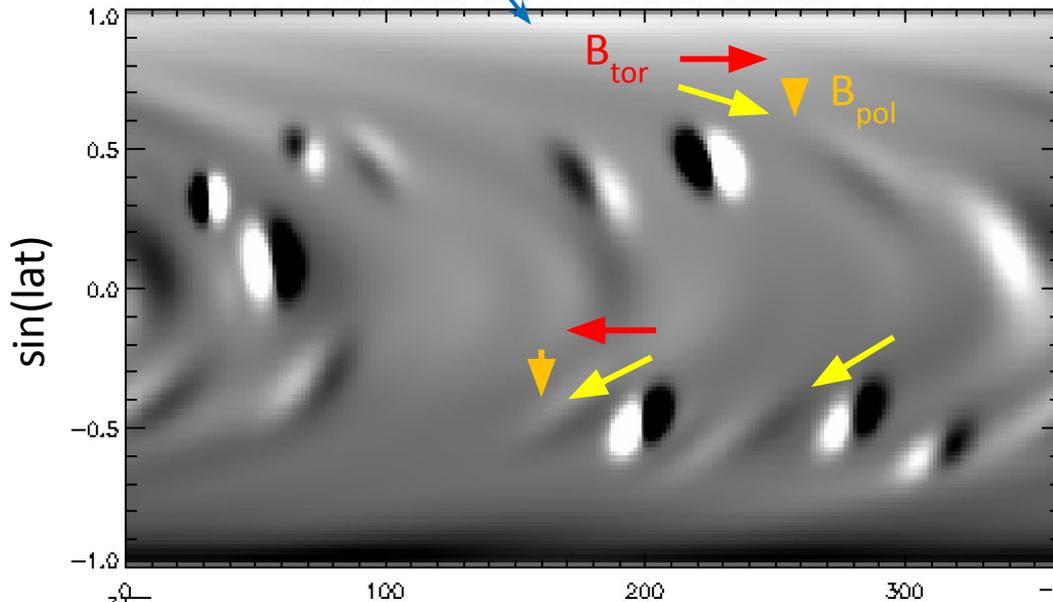
**Hale's Law:** toroidal field from diff. rotation

**Joy's Law:** tilt of bipole  
□ poloidal field

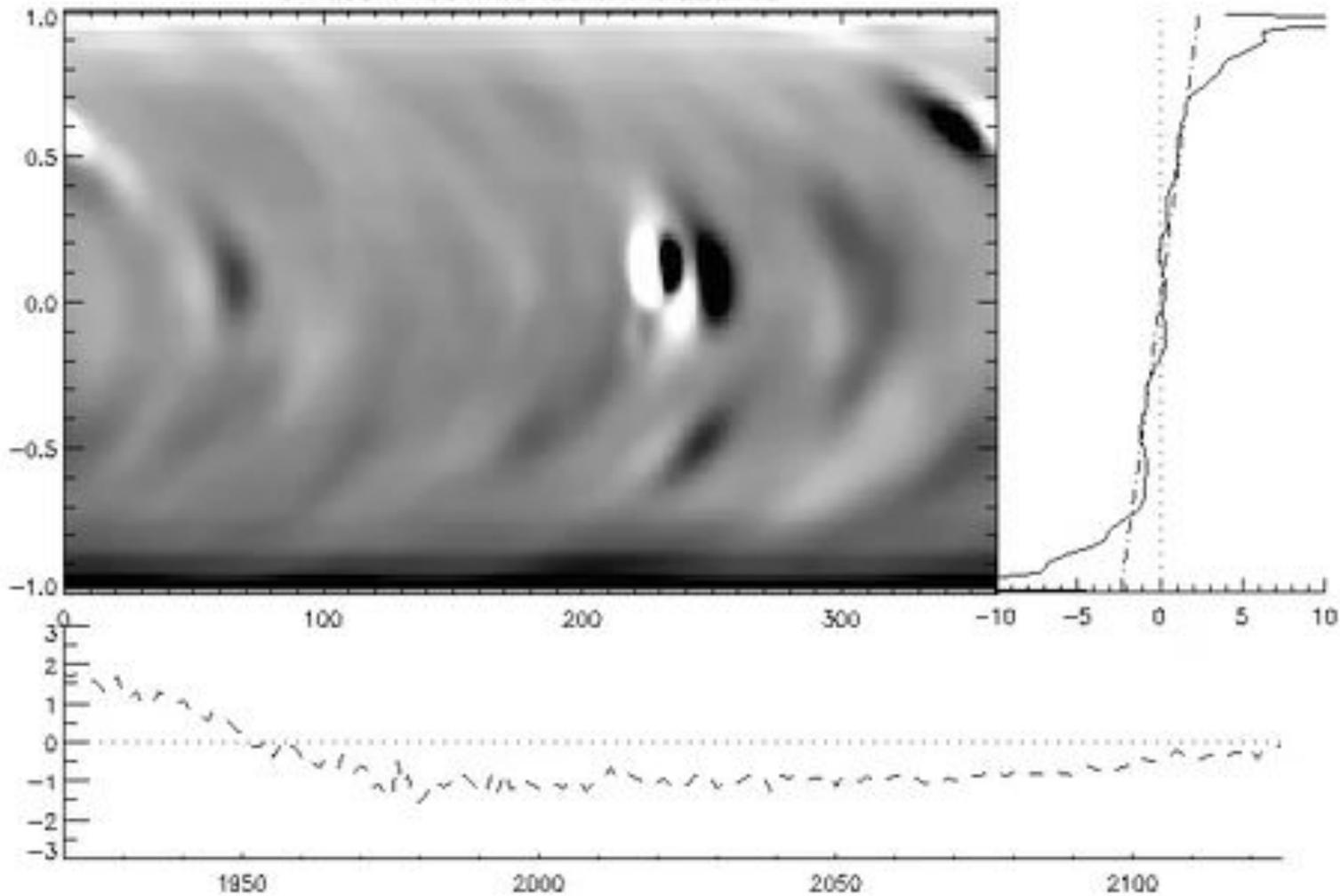
remnant of last cycle

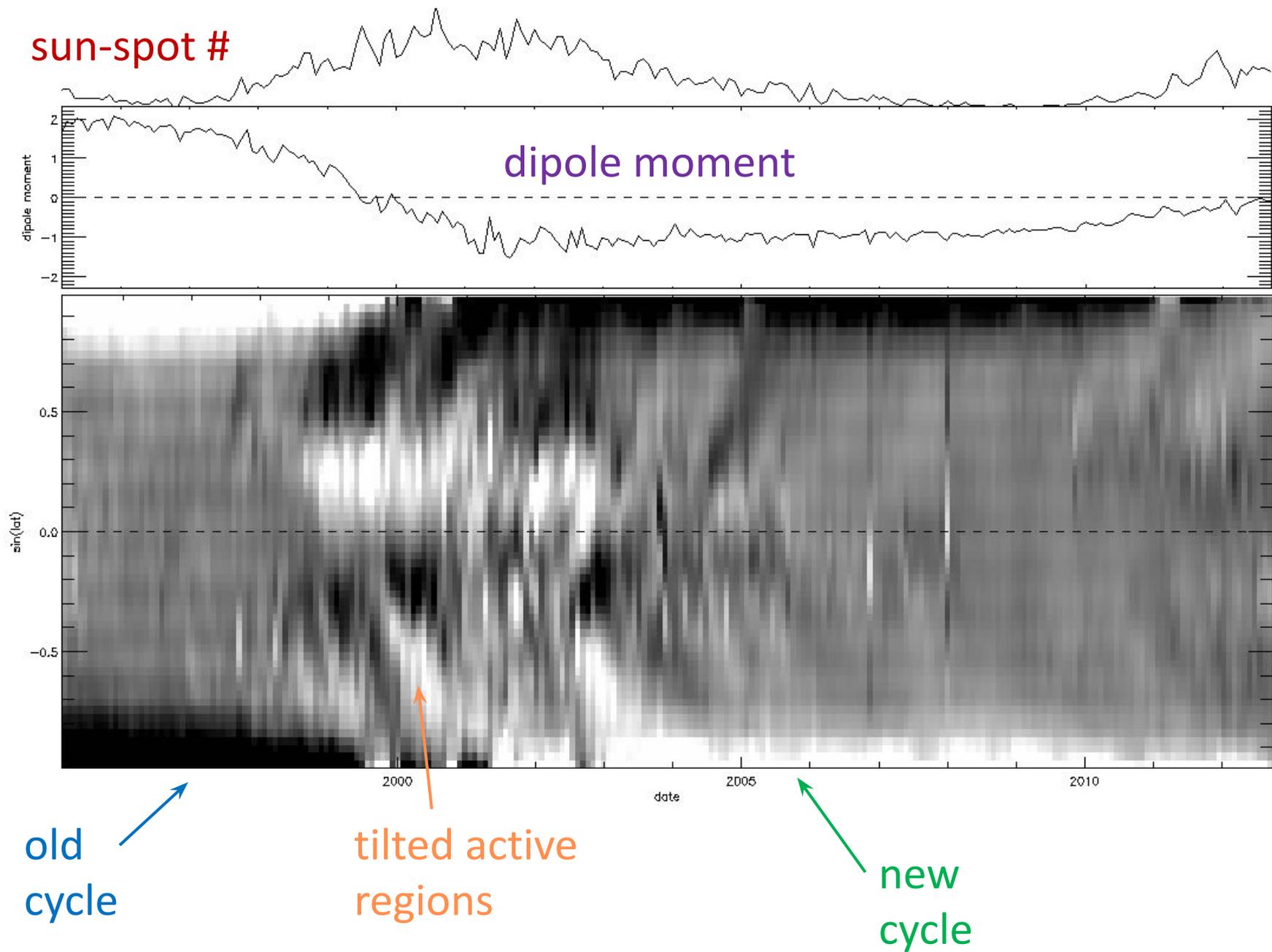


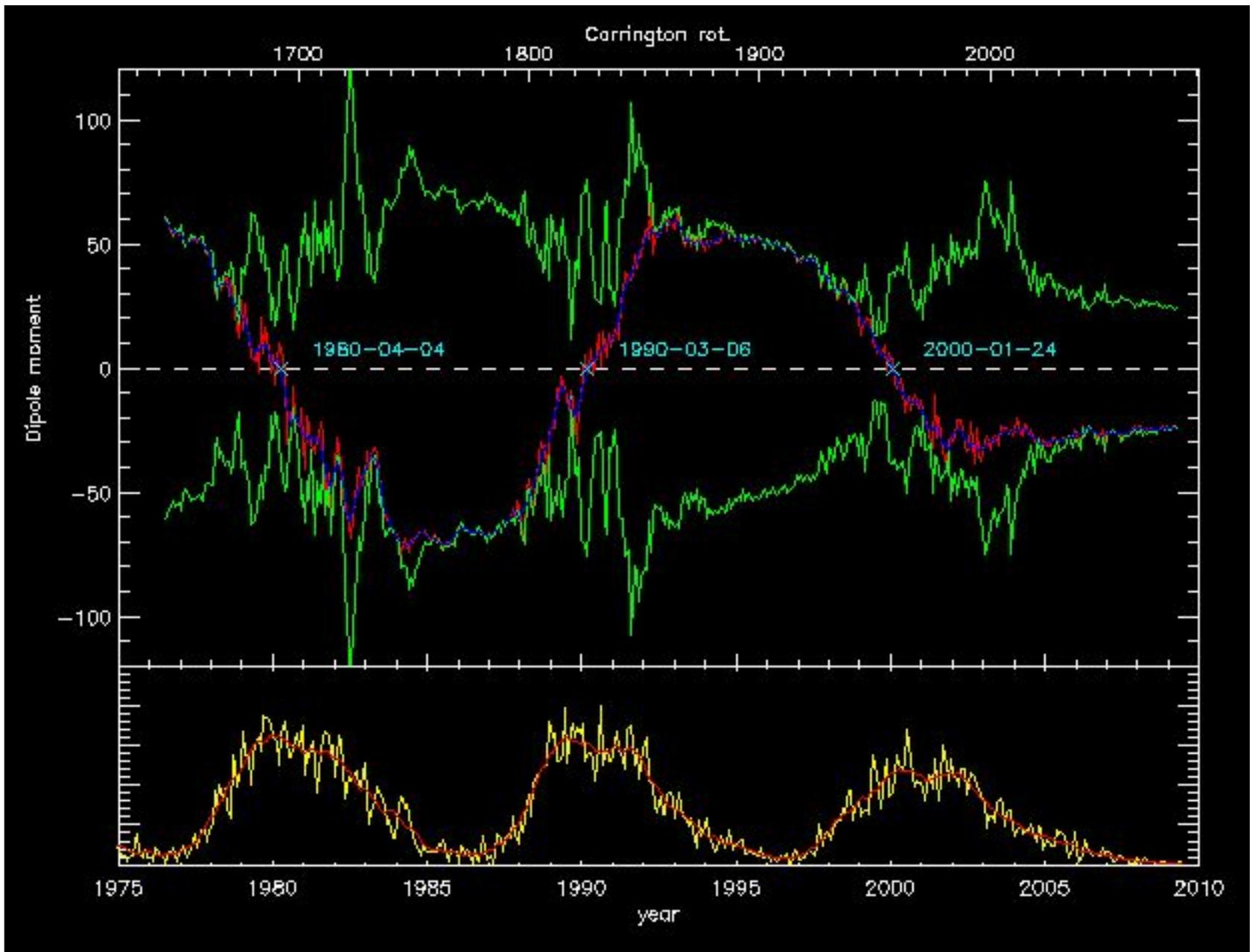
CR 1924: 1997-06-17T21:36:10.000Z



CR 1921: 1997-03-28T01:46:52.000Z

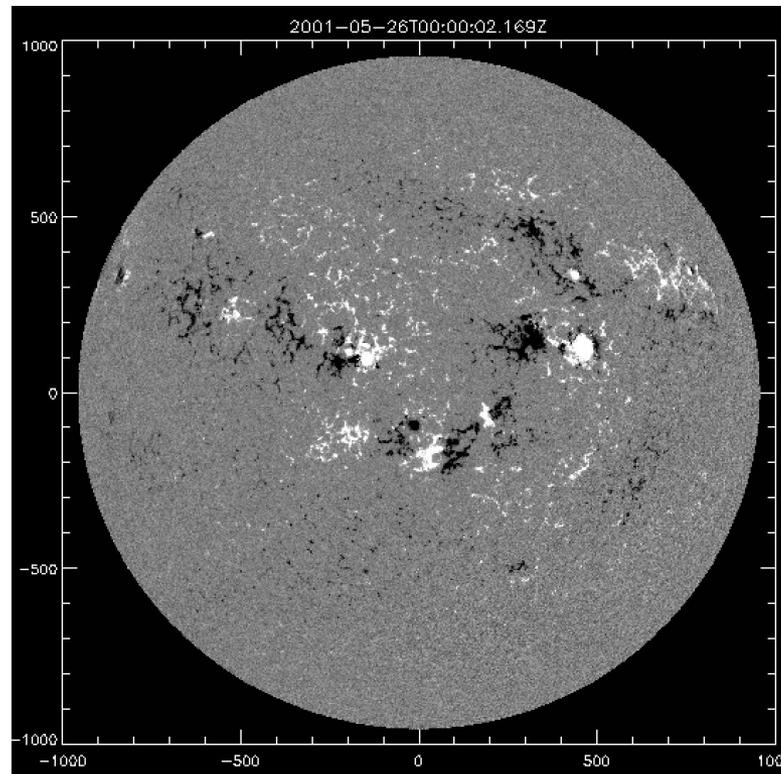






Activity 51: Question: with this value of  $D$  [=250 km<sup>2</sup>/s], what is the characteristic time scale for flux to disperse over the solar surface (hint: Eq. 3.20)? [H]ow important is the meridional advection from equator to pole (with a characteristic velocity of 10 m/s) in transporting the field within the duration of a solar cycle?

$$\tau_d \sim \frac{L_t^2}{\eta} . \quad (3.20)$$



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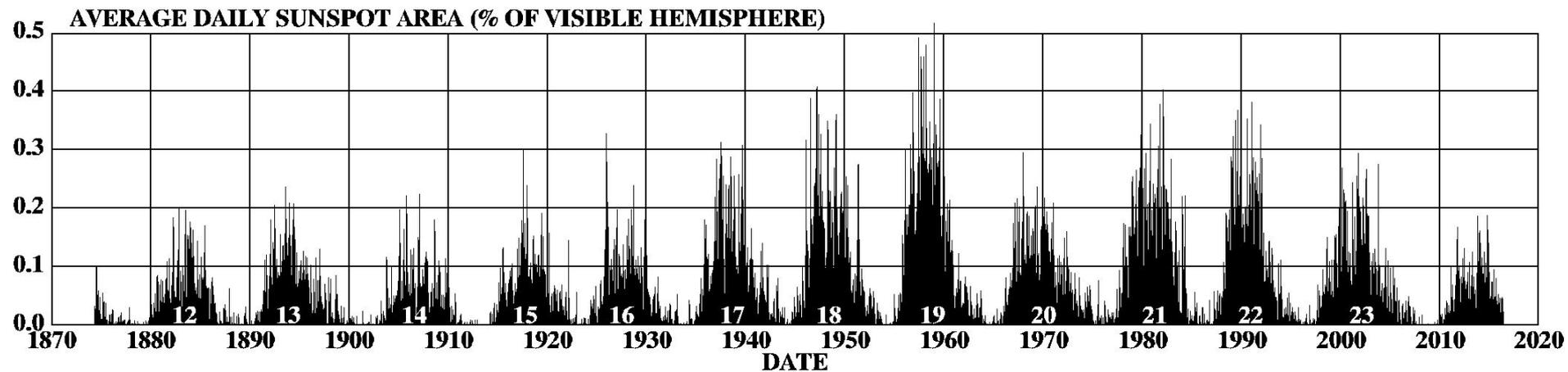
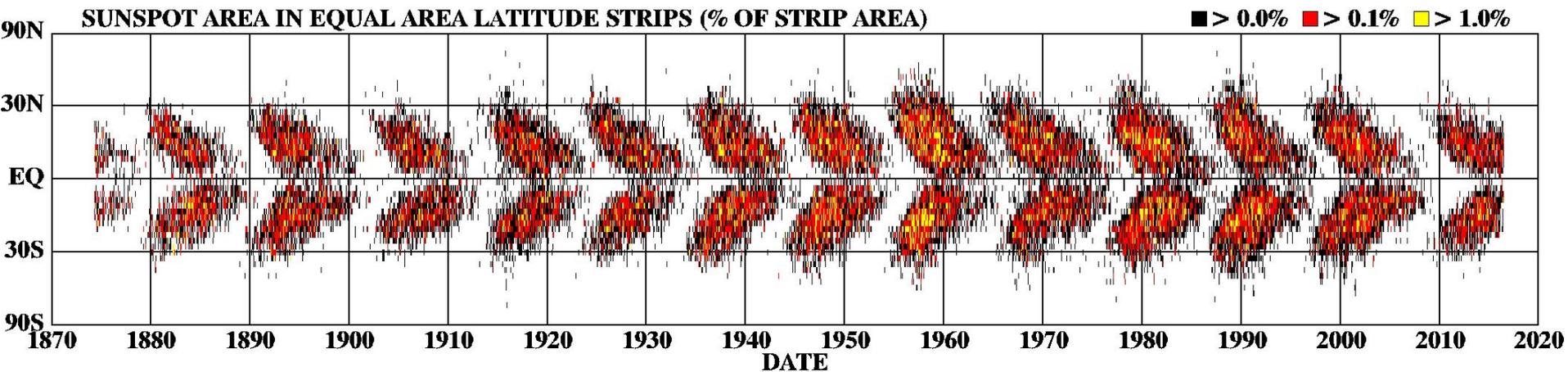
$$\tau_d \sim \frac{L_t^2}{\eta} . \quad (3.20)$$

$$L_t \sim R_\odot = 7 \times 10^5 \text{ km}$$

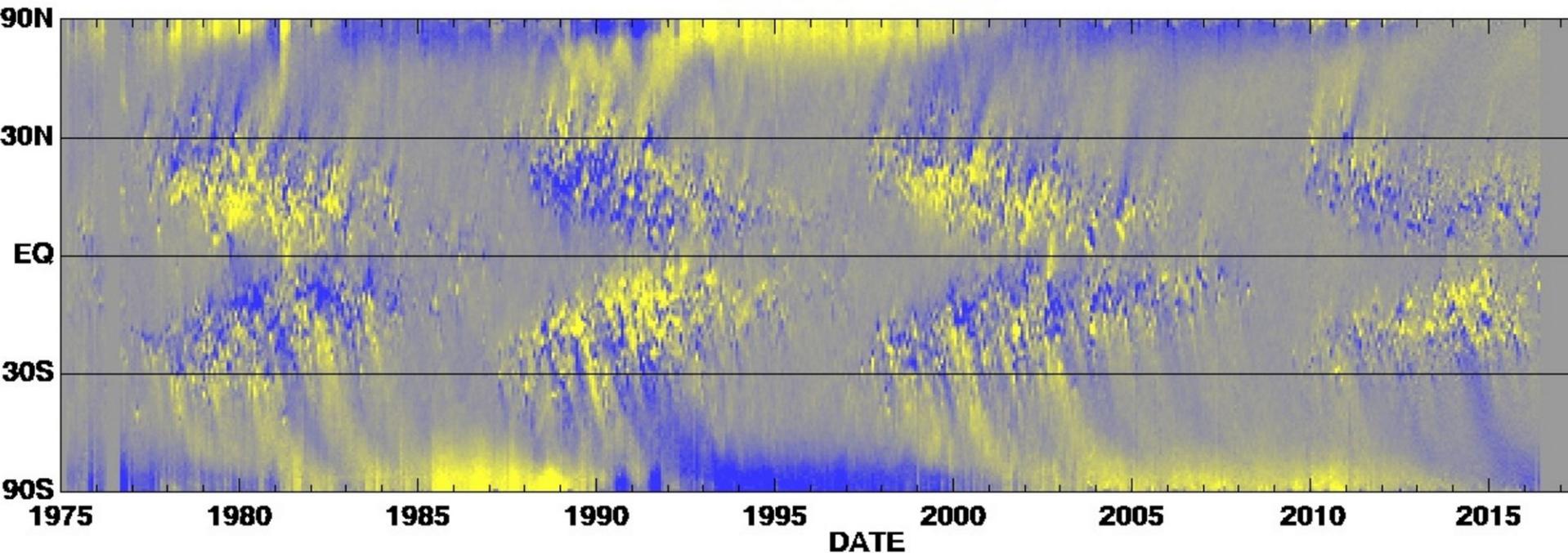
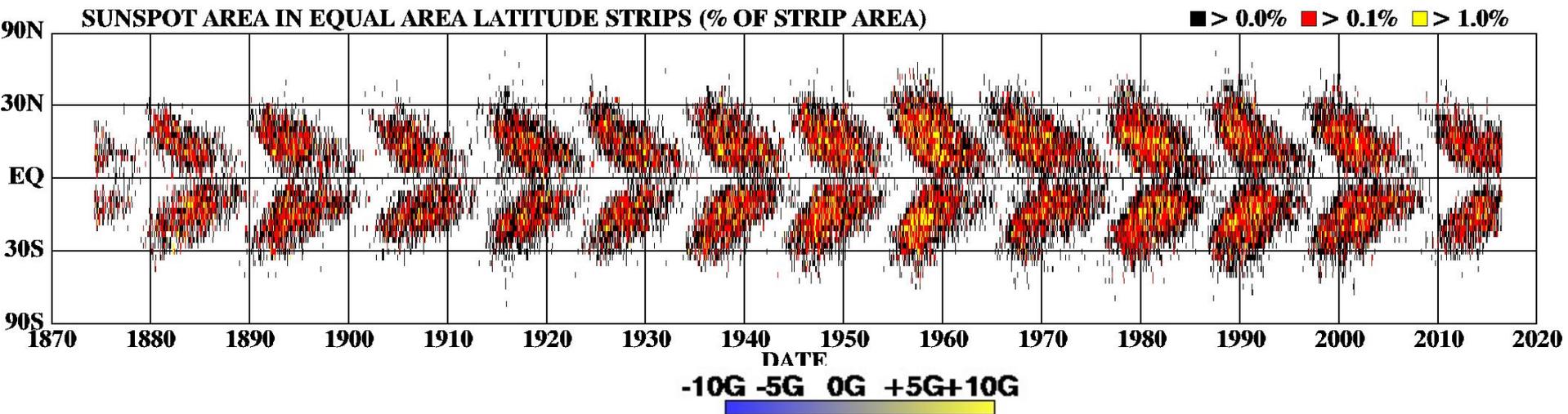
$$\tau_d \sim L_t^2/D = 2 \times 10^9 \text{ s} = 65 \text{ y}$$

$$\begin{aligned} \text{vs. } R_\odot/v &= (7 \times 10^5 \text{ km})/(10^{-2} \text{ km/s}) \\ &= 7 \times 10^7 \text{ s} = 2.3 \text{ y} \end{aligned}$$

# DAILY SUNSPOT AREA AVERAGED OVER INDIVIDUAL SOLAR ROTATIONS



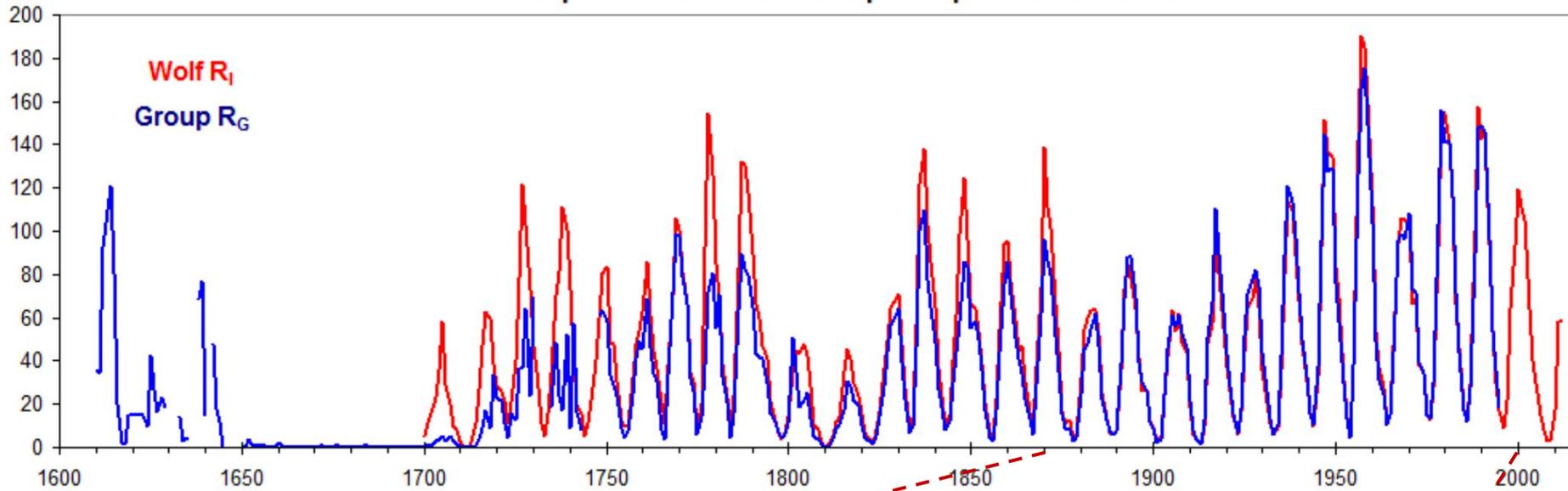
# DAILY SUNSPOT AREA AVERAGED OVER INDIVIDUAL SOLAR ROTATIONS



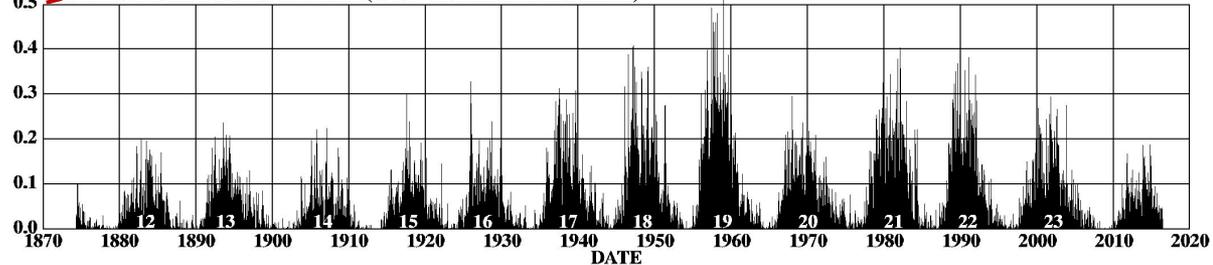
# Long(er)-term variation on the solar dynamo:

Cliver *et. al* 2013

Comparison of Wolf and Group Sunspot Number Series



AVERAGE DAILY SUNSPOT AREA (% OF VISIBLE HEMISPHERE)

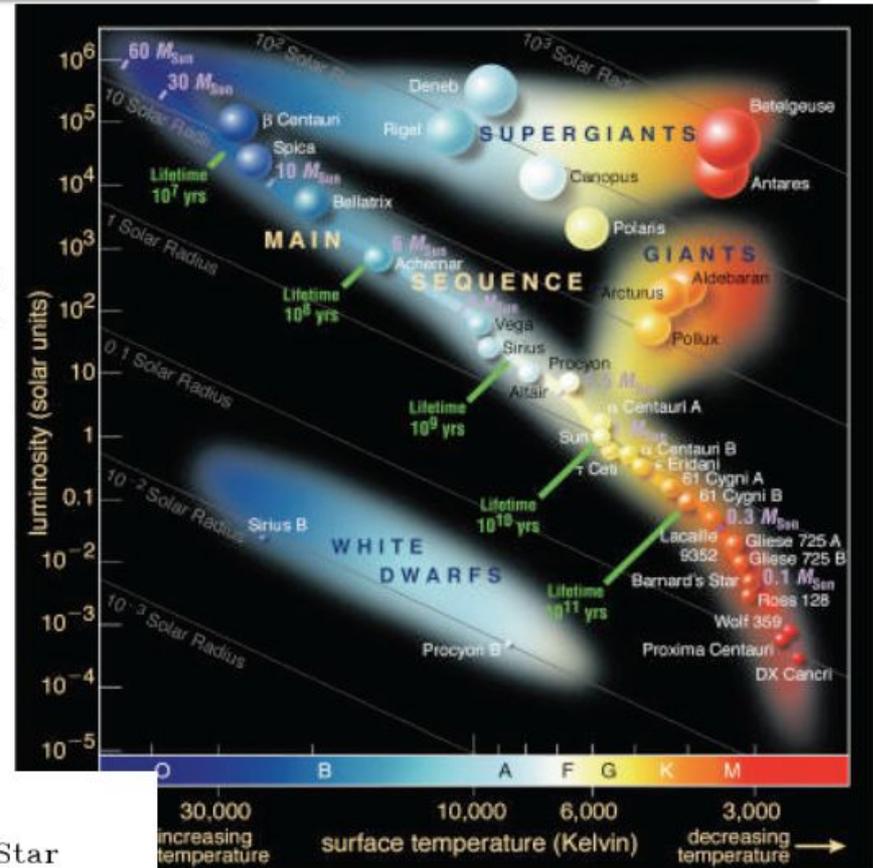


<http://solarscience.msfc.nasa.gov/>

HATHAWAY NASA/ARC 2016/07

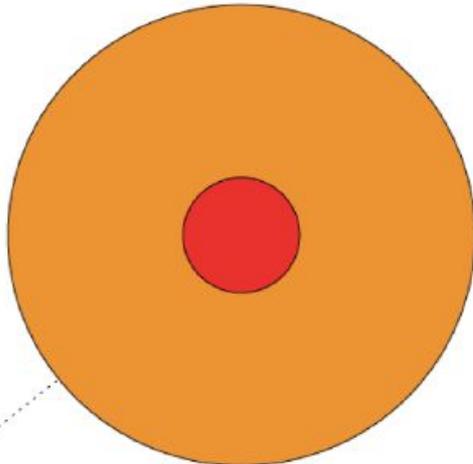
# STELLAR MAGNETIC FIELDS

- correlations between stellar types and magnetic field properties, probably due to geometry of convection zones
- stars with outer convection zones (late-type stars) have observed magnetic fields whose strength tends to increase with their angular velocity
- Cyclic variations are known to exist only for spectral types between G0 and K7).

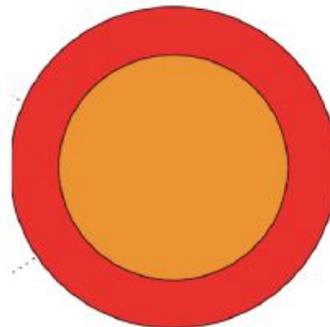


Stanley

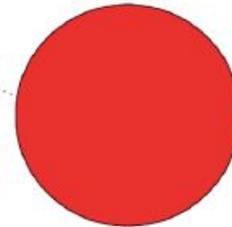
B Star



Sun

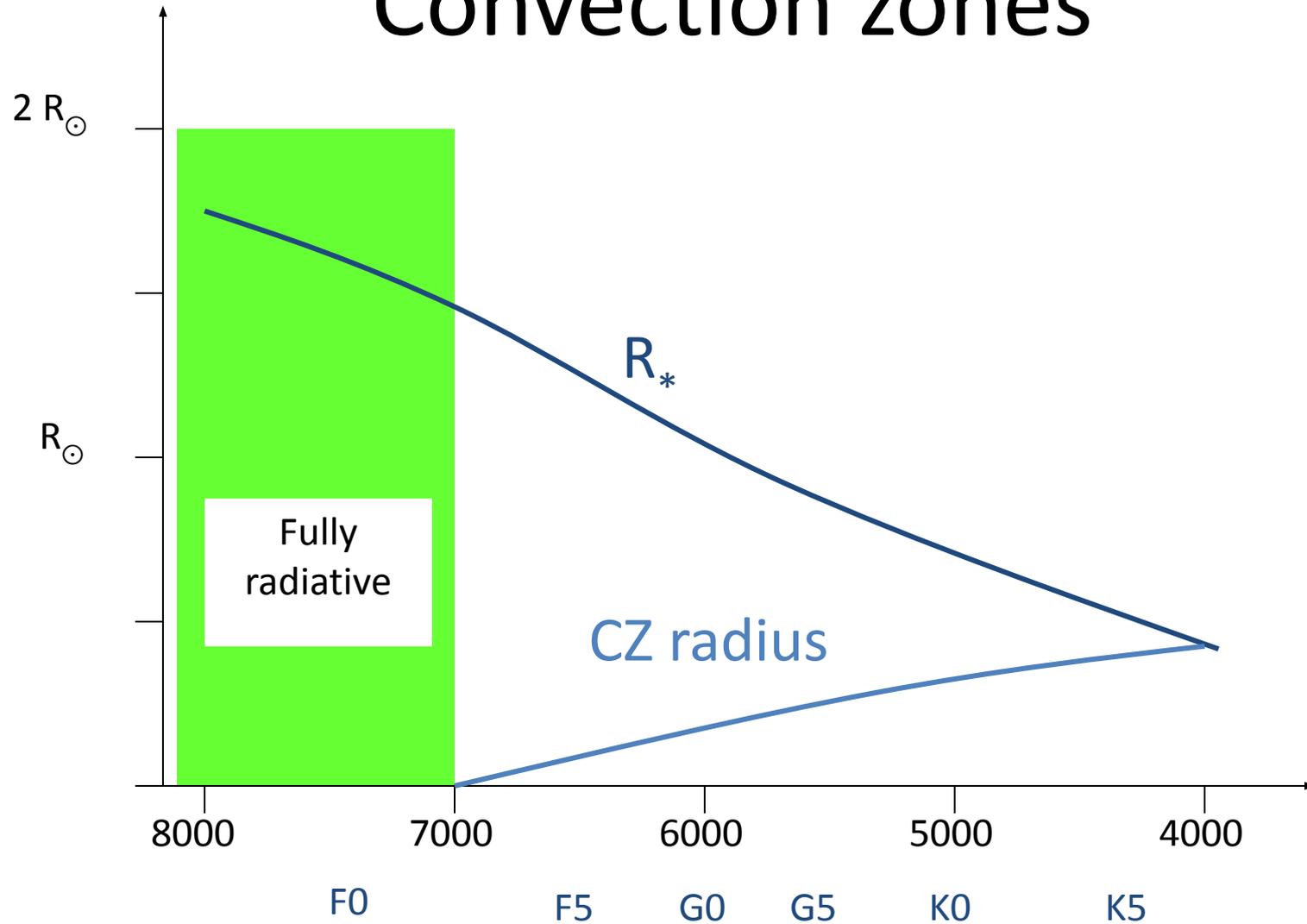


M Star



Convective  
 Radiative

# Convection zones

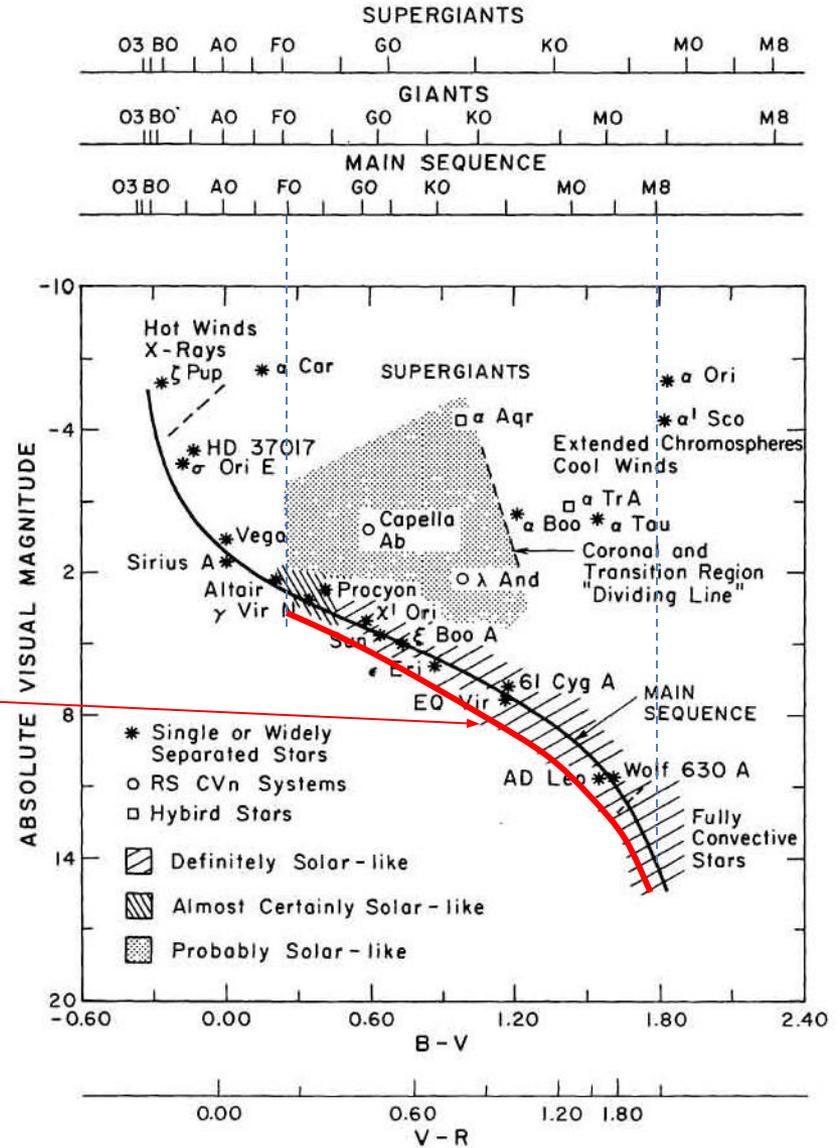


# Other stars

Evidence of magnetic activity

Activity on main sequence:  
types **F**   **M**

$$B-V > 0.4$$



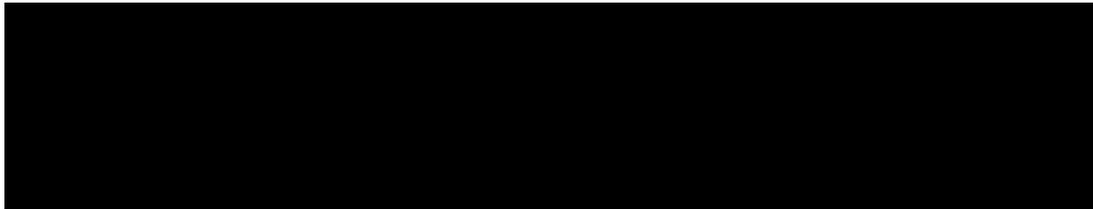
(From Linsky 1985)

Activity 50: Relate the Rossby number in Eq. (4.2) to the dynamo number  $C_\Omega$  and the magnetic Reynolds number  $R_m$  in Eq. (4.20):

$$N_R = R_m^2 / C_\Omega.$$

$$N_R = \frac{v_t}{\Omega L_t}, \quad (4.2)$$

$$C_\alpha = \frac{\alpha_t R}{\eta}, \quad C_\Omega = \frac{\Omega_t R^2}{\eta}, \quad R_m = \frac{u_t R}{\eta}, \quad (4.20)$$



# The Dynamo Number

Dynamo is linear instability for  $D > D_{\text{crit}}$

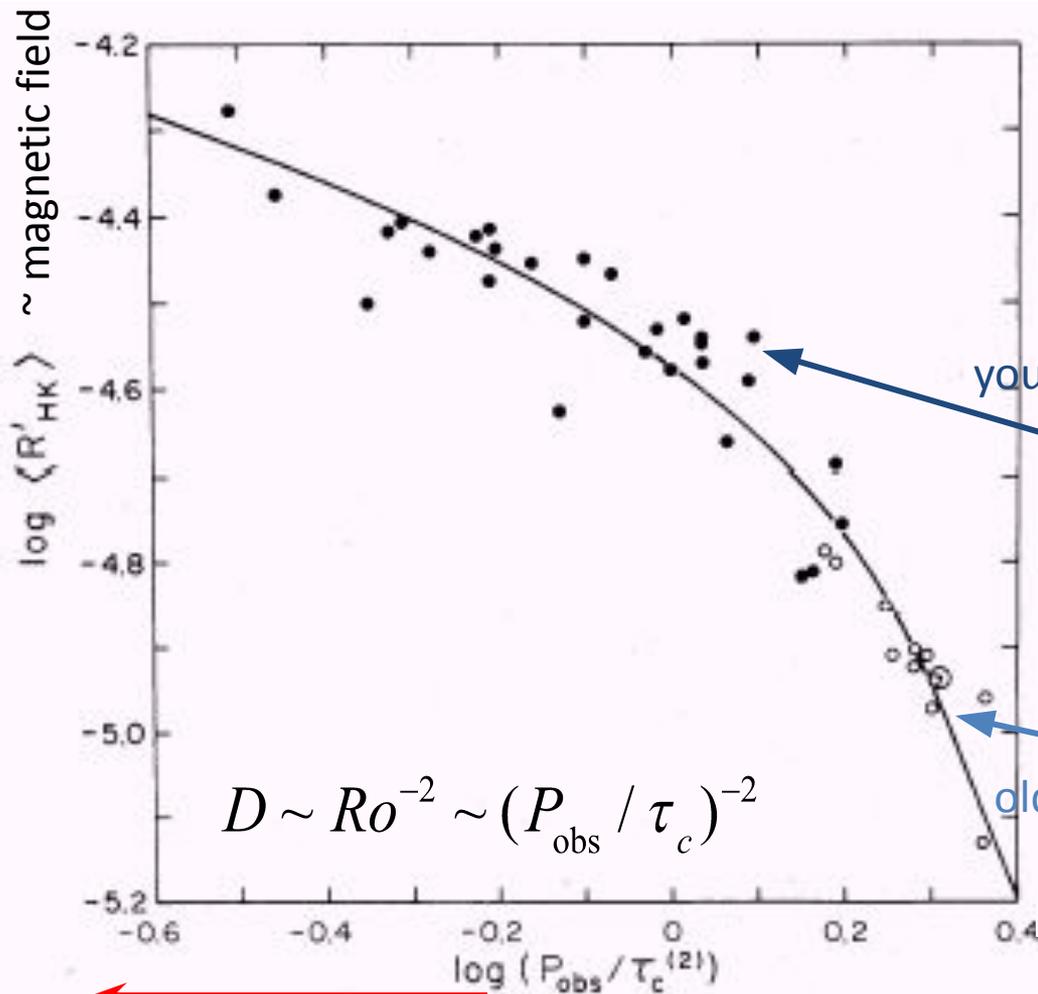
$$D \equiv C_\alpha \times C_\Omega = \frac{\alpha_t \Omega_t R^3}{\eta_{\text{CZ}}^2} . \quad \blacksquare$$

$$C_\alpha = \frac{\alpha_t R}{\eta} , \quad C_\Omega = \frac{\Omega_t R^2}{\eta} ,$$

$$\eta_{\text{CZ}} \sim \frac{R^2}{\tau_c} \quad \alpha_t \equiv \tau \langle \mathbf{v} \cdot (\nabla \times \mathbf{v}) \rangle \sim \Omega_t R$$

$$D \sim (\Omega_t \tau_c)^2 \sim Ro^{-2}$$

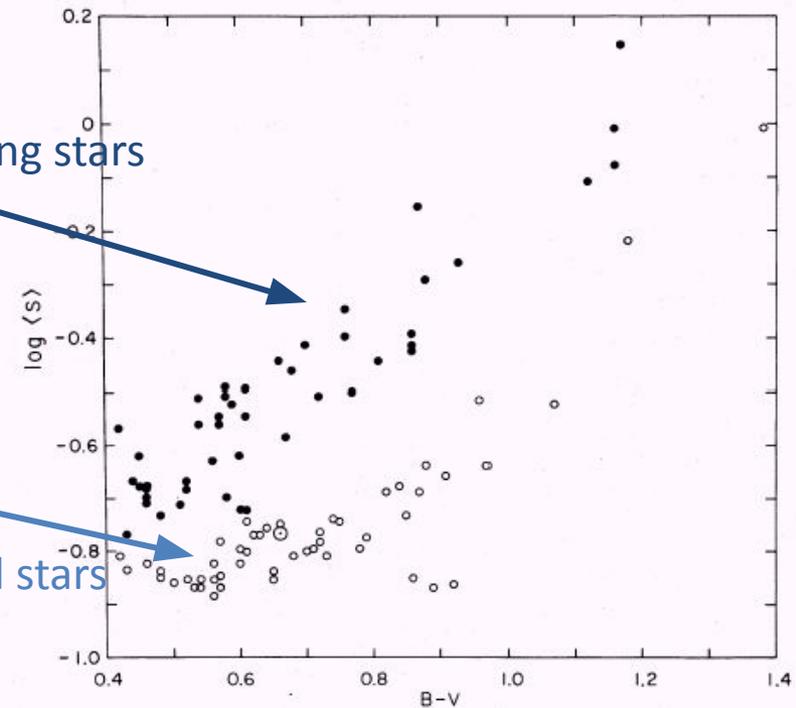
# Activity vs. Rossby Number



Stars spin down over time:

$$\Omega_{\text{rot}} \propto t^{-1/2}. \quad (10.3)$$

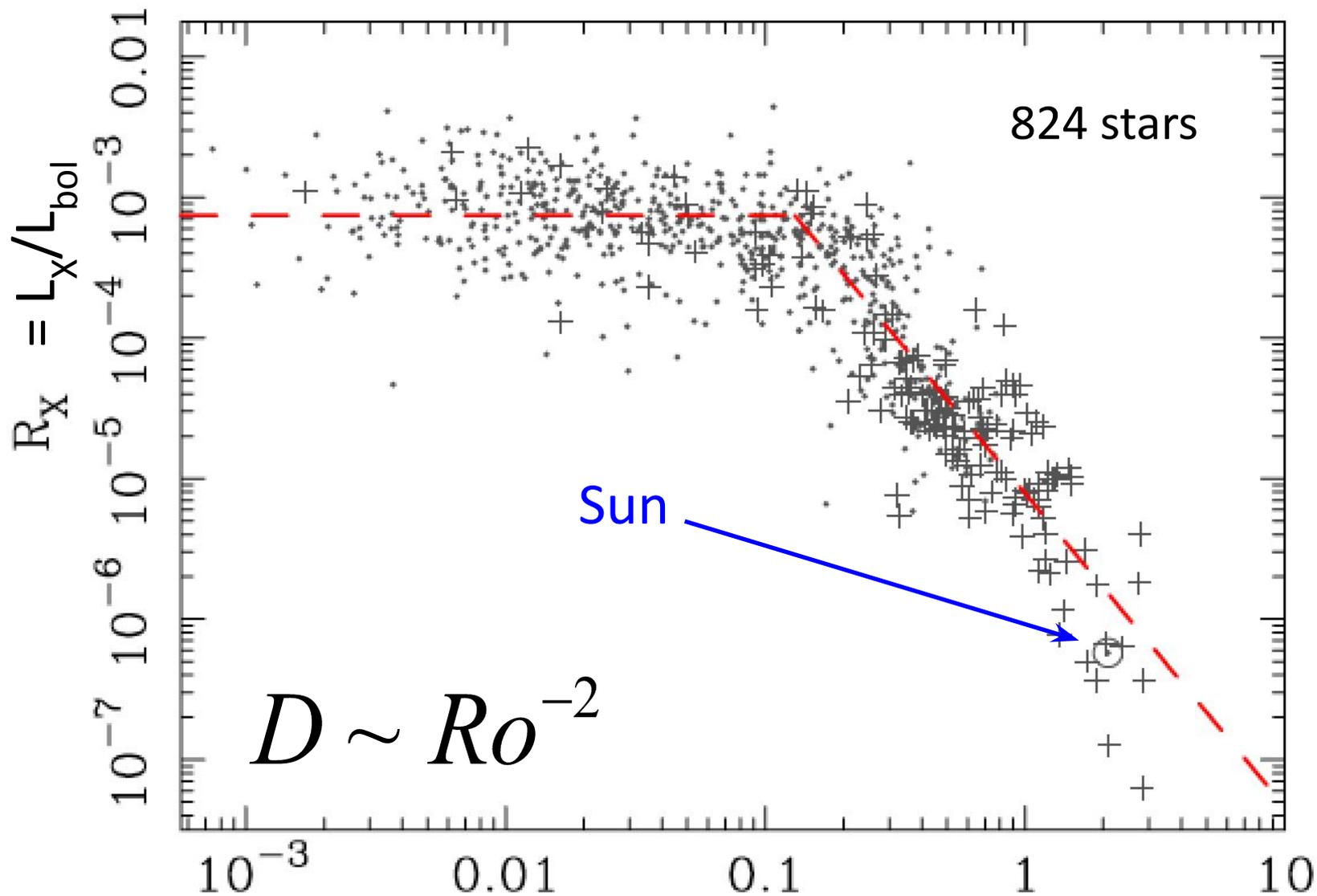
(Skumanich law)



young stars

old stars

← **Increasing D** (from Noyes et al. 1984)



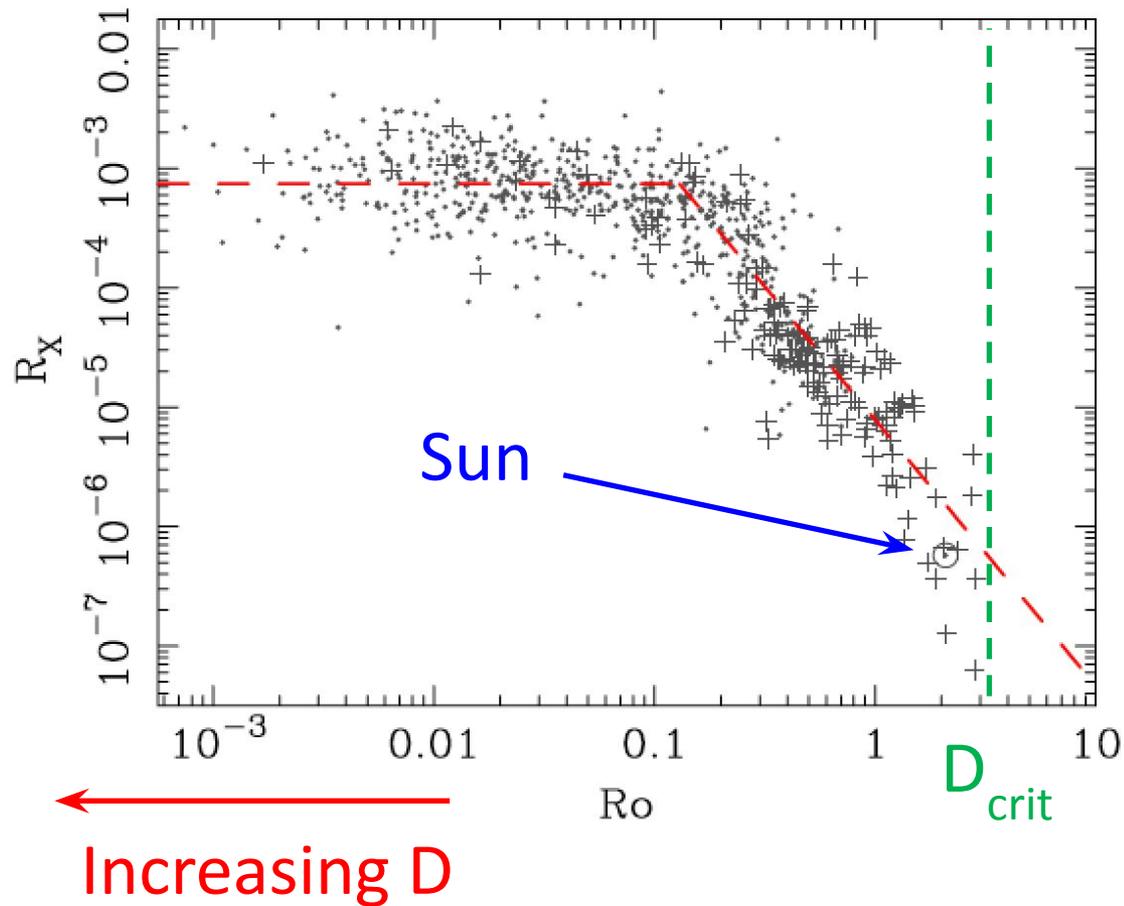
←  
Increasing D

$$Ro = P_{\text{rot}} / \tau_c$$

Wright et al. 2011

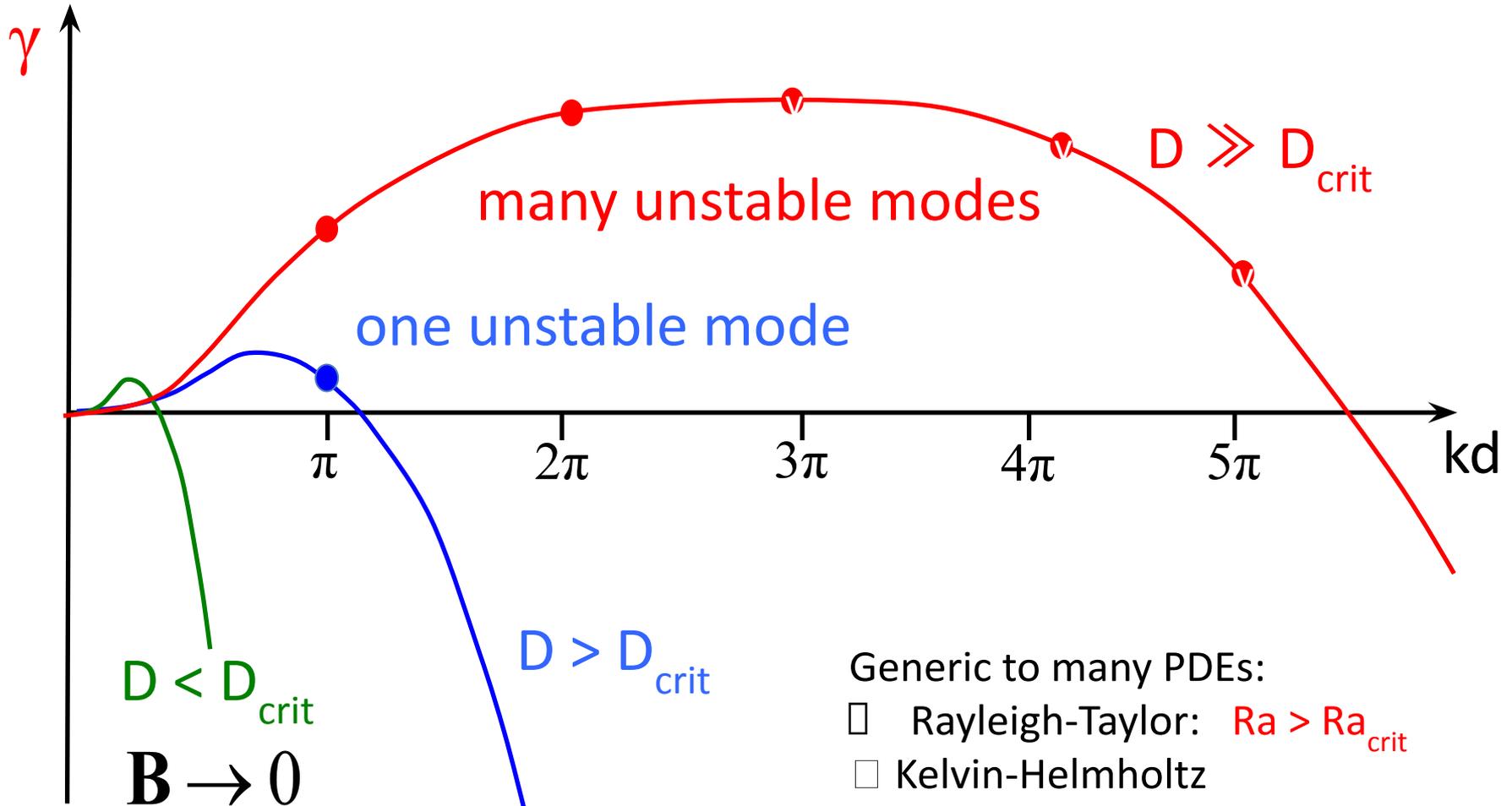
From Cameron & Schussler 2017:

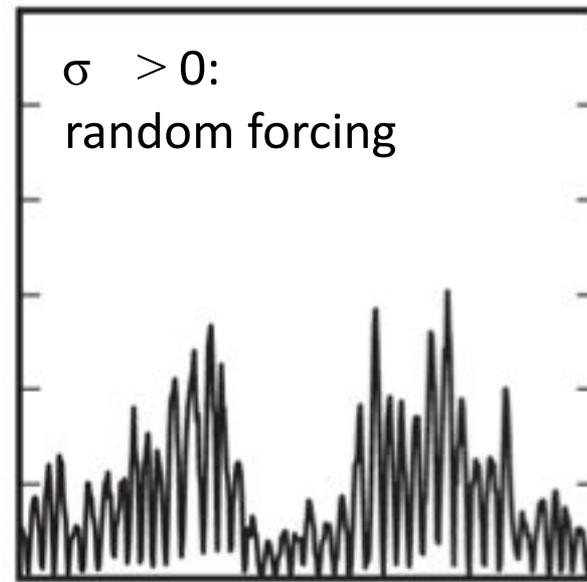
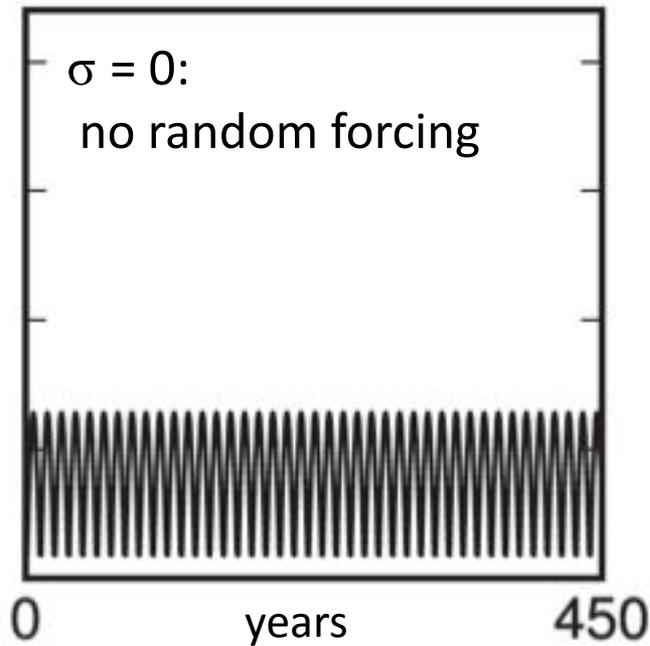
- The solar dynamo is operating near  $D \approx D_{\text{crit}}$
- Only single normal mode is unstable
- Long term variation from fluctuations



$\gamma > 0$

$$D = \frac{\alpha \Omega d^3}{\eta^2} > 2(kd)^3 \geq 2\pi^3 = D_{\text{crit}}$$



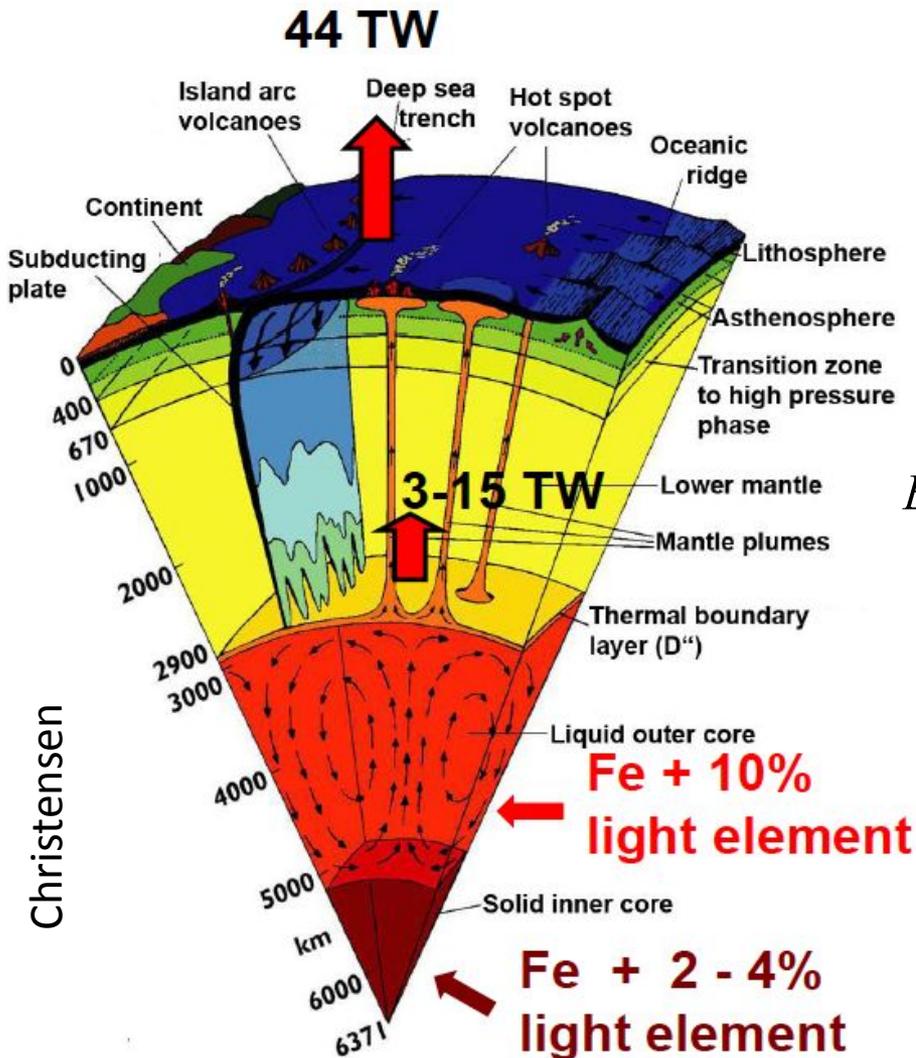


Cameron &  
Schussler 2017

one normal mode:

$$dX = \underbrace{(\beta + i\omega_0)}_{\text{linear instability}} - \underbrace{(\gamma_r + i\gamma_i)|X|^2}_{\text{non-linear saturation}} X dt + \underbrace{\sigma X dW_c}_{\text{random forcing (Wiener process)}} = 0,$$

# How the dynamo works for Earth



Non-conducting mantle

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} = 0$$

$$\mathbf{B} = -\nabla \chi \quad \nabla \cdot \mathbf{B} = -\nabla^2 \chi = 0$$

$$\chi(r, \theta, \phi) = \sum_{\ell, m} g_{\ell, m} Y_{\ell}^m(\theta, \phi) \left( \frac{R_{\oplus}}{r} \right)^{\ell+1}$$

$$B_r(r, \theta, \phi) = -\frac{\partial \chi}{\partial r} = \sum_{\ell, m} (\ell+1) g_{\ell, m} Y_{\ell}^m(\theta, \phi) \left( \frac{R_{\oplus}}{r} \right)^{\ell+2}$$

simplifies w/ increasing r

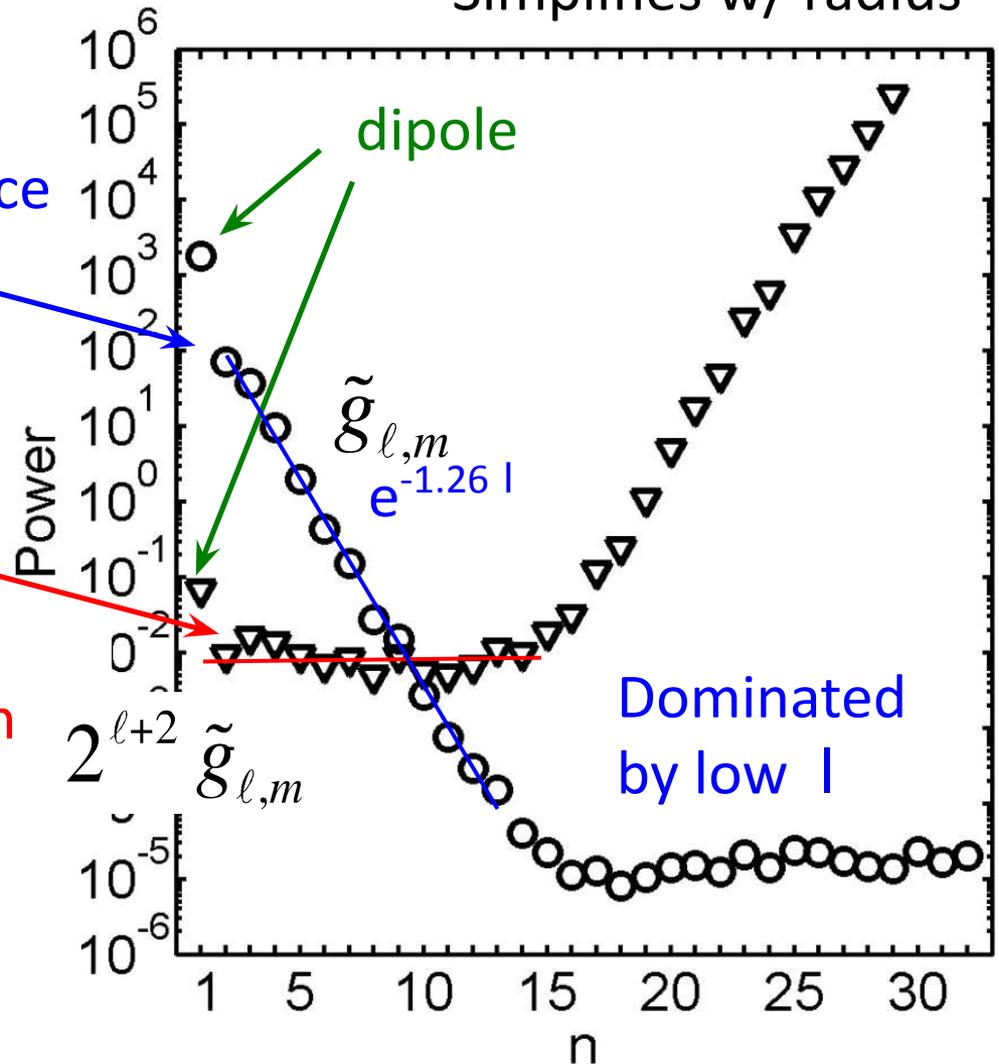
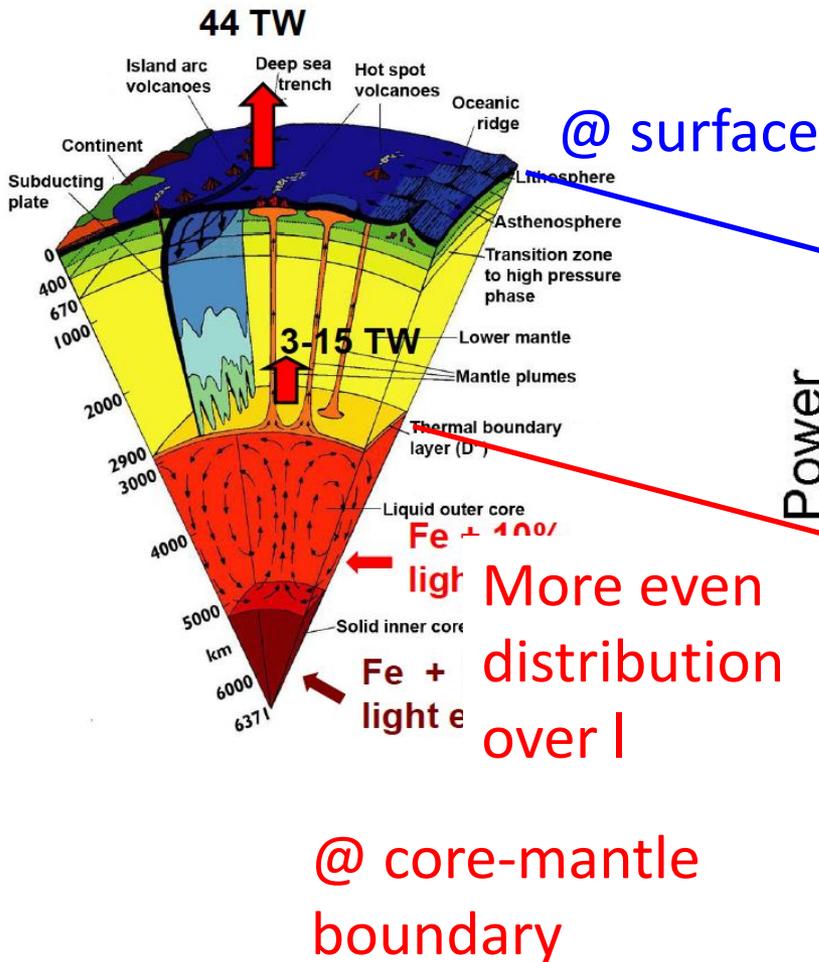
Turbulent conducting fluid:  
**DYNAMO**

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} \neq 0$$

Complex flows – complex field

$$B_r(r, \theta, \phi) = \sum_{\ell, m} (\ell + 1) \tilde{g}_{\ell, m} Y_{\ell}^m(\theta, \phi) \left( \frac{R_{\oplus}}{r} \right)^{\ell+2}$$

Simplifies w/ radius

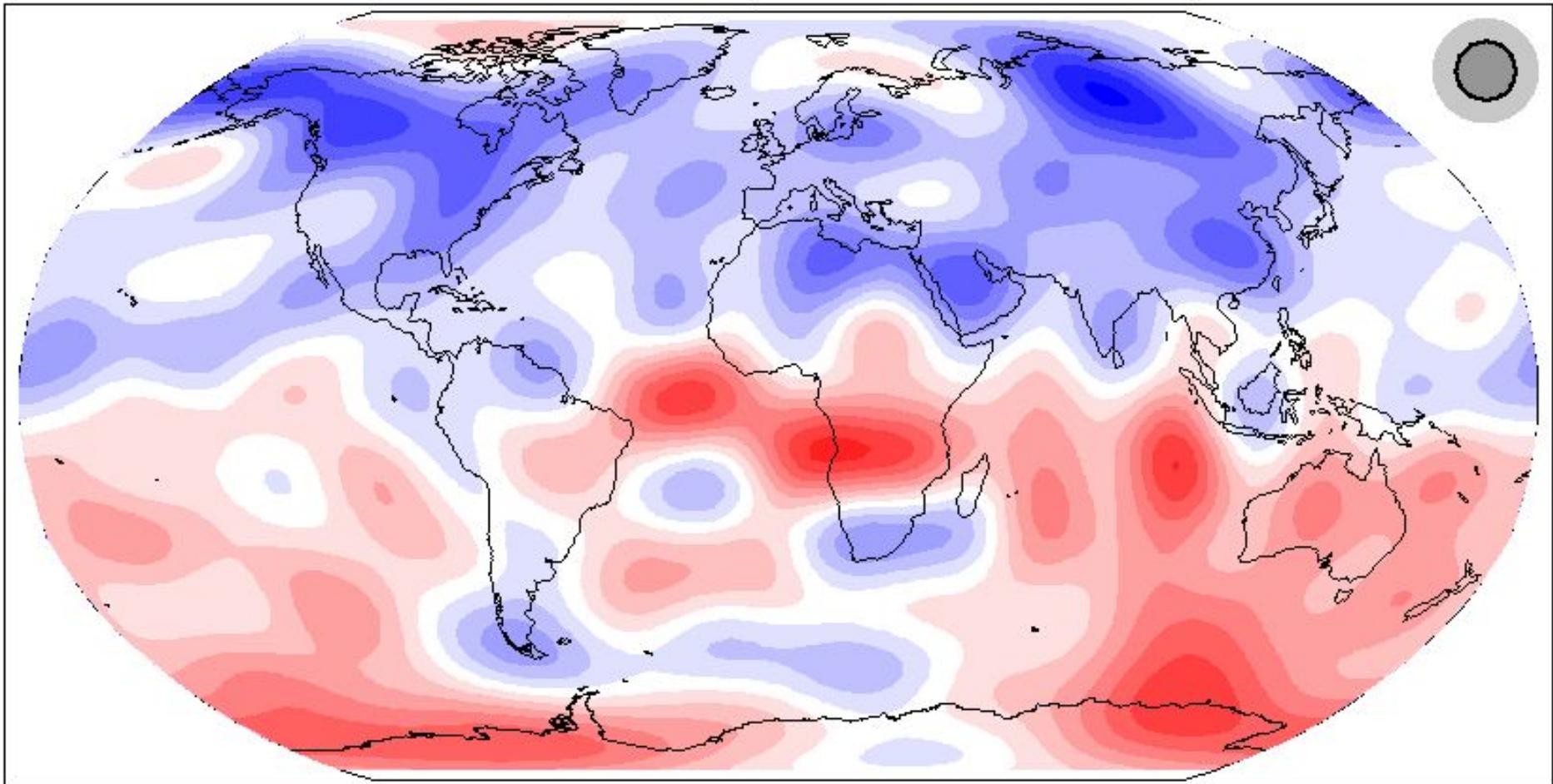


# Observation: what $B_r$ looks like today

@ core-mantle boundary: lower boundary of potential region

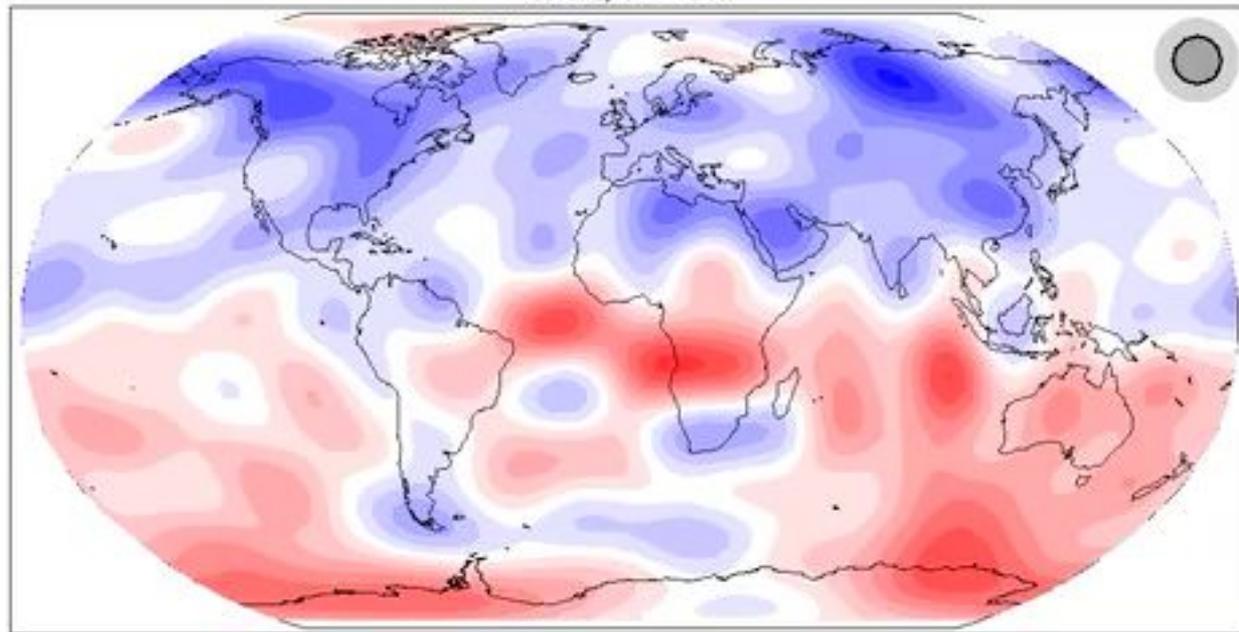
$\Pi < 14$

2010:  $B_r$  @  $r = 0.55$



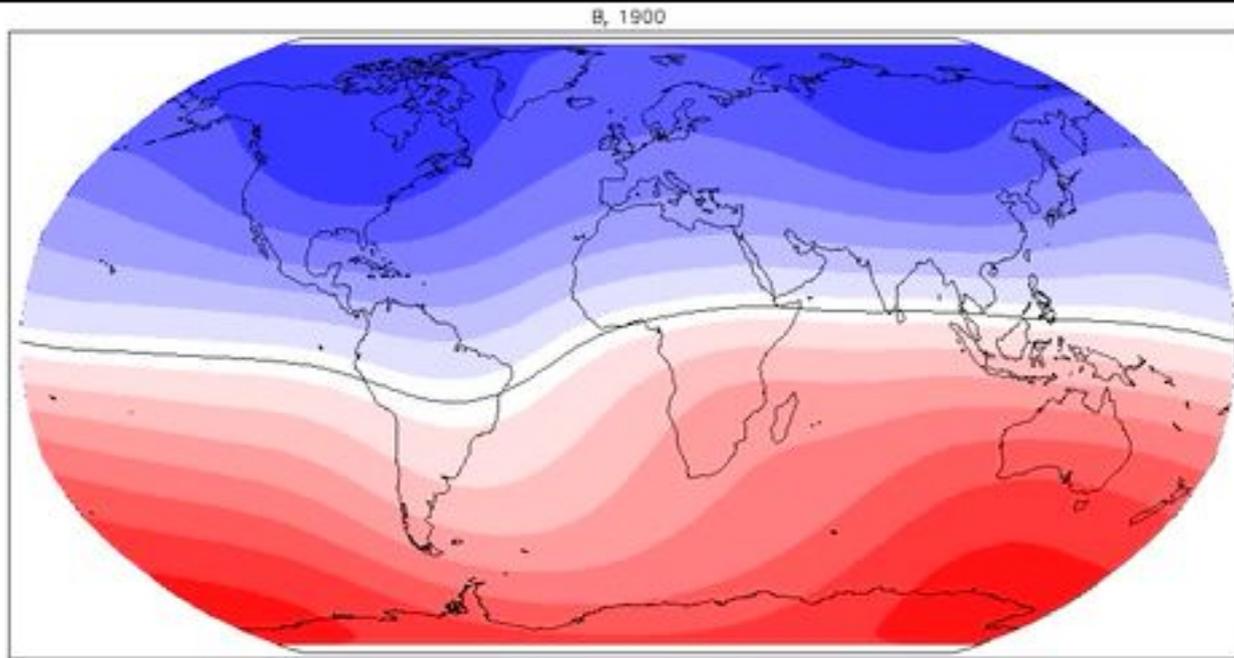
# Simplifies w/ increasing r

$$B_r(r, \theta, \phi) = -\frac{\partial \chi}{\partial r} = \sum_{\ell, m} (\ell + 1) g_{\ell, m} Y_{\ell}^m(\theta, \phi) \left( \frac{R_{\oplus}}{r} \right)^{\ell+2}$$



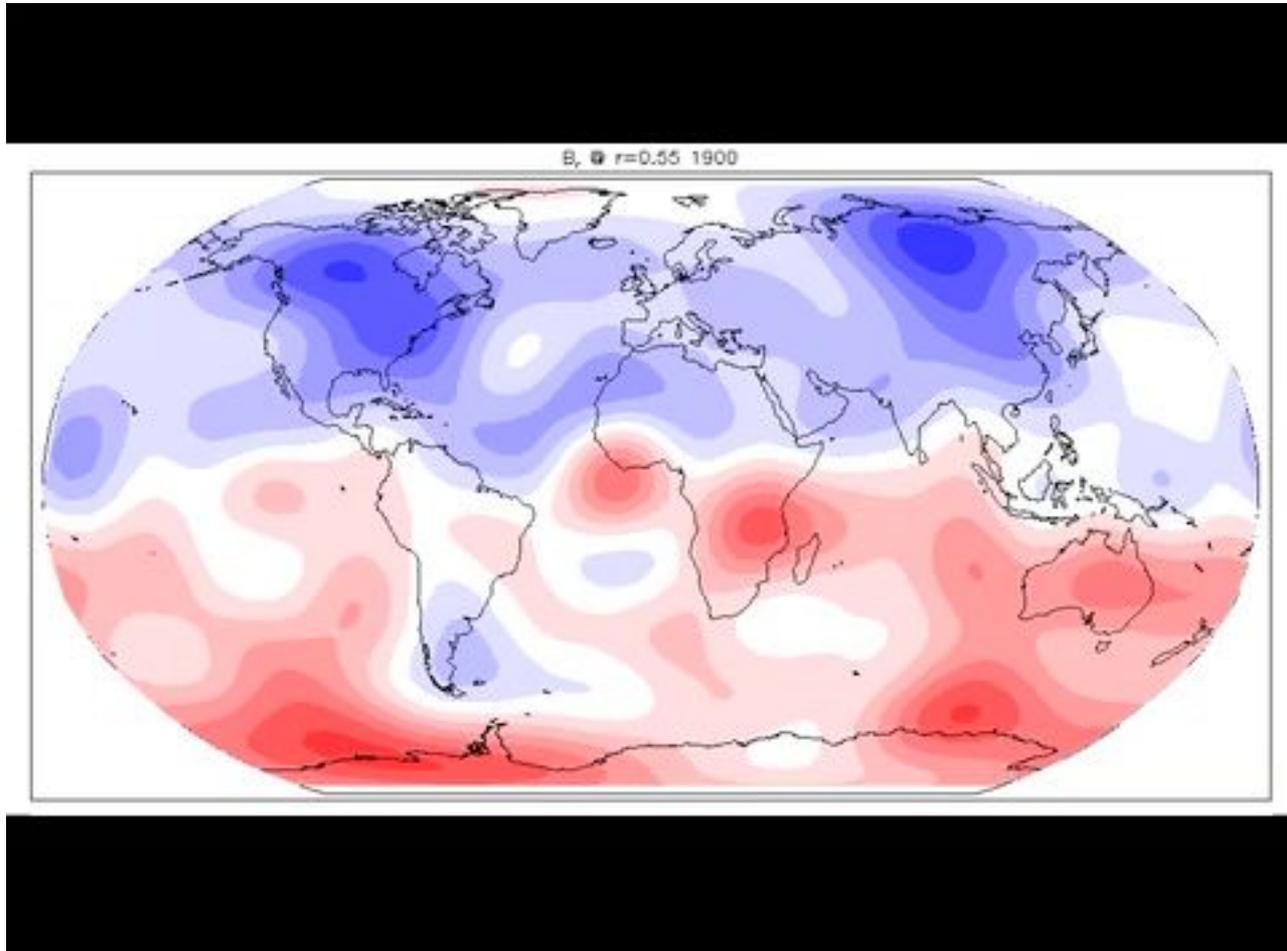
# Evolution of field

@ surface for 100 years



# Evolution of field

@ core-mantle boundary for 100 years



# Use evolution to infer fluid velocity

2010: B, r = 0.55

B, r = 0.55 1910

$$v \sim 3 \times 10^{-4} \text{ m/s}$$

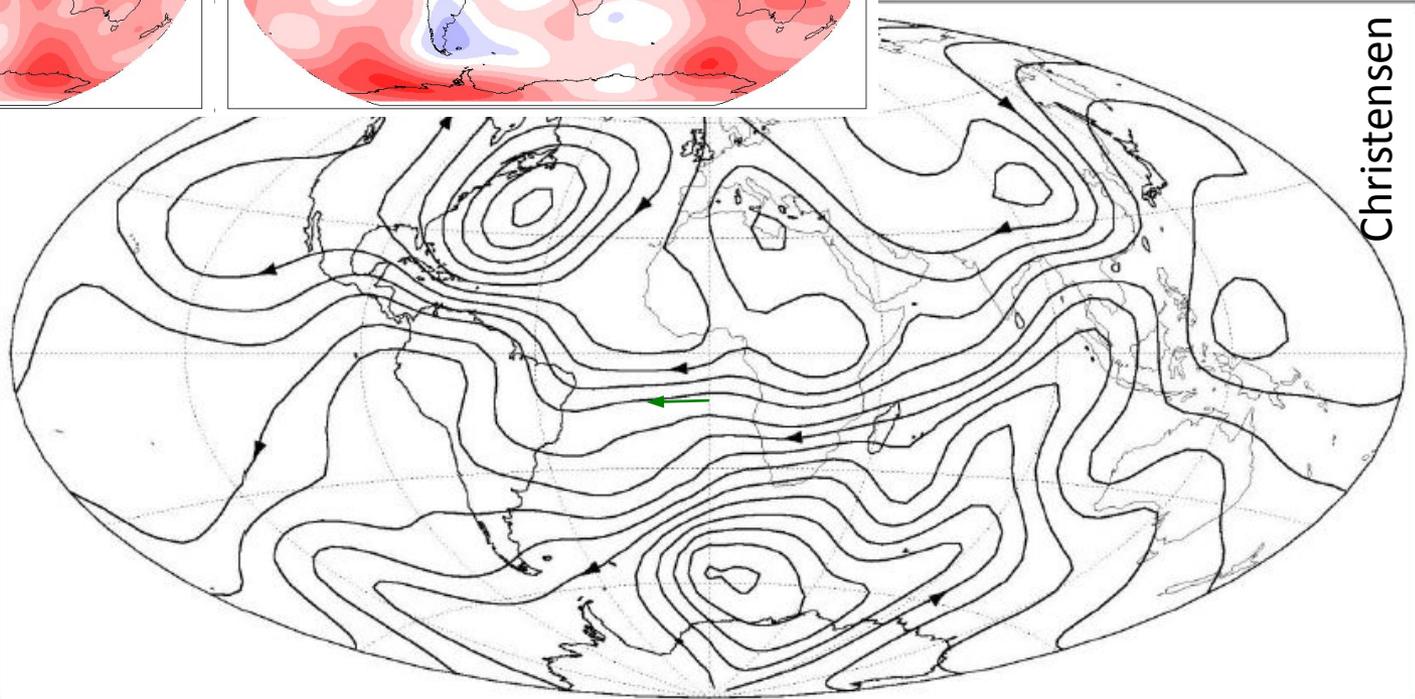
$$\square \Delta x = 1,000 \text{ km}$$

$$\square \text{ in 100 years}$$

Subsonic flow,  
Ignore  
stratification

$$\nabla \cdot \mathbf{v} = 0$$

Ignore diffusion

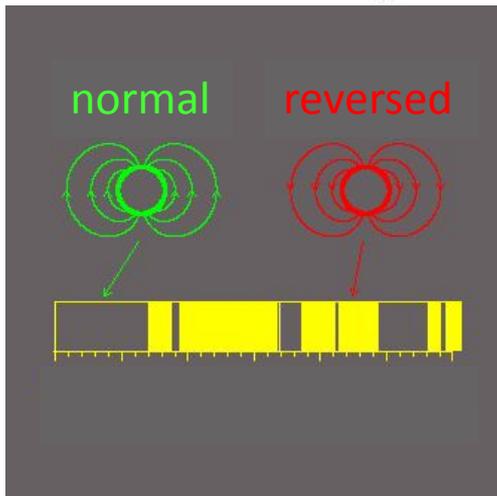
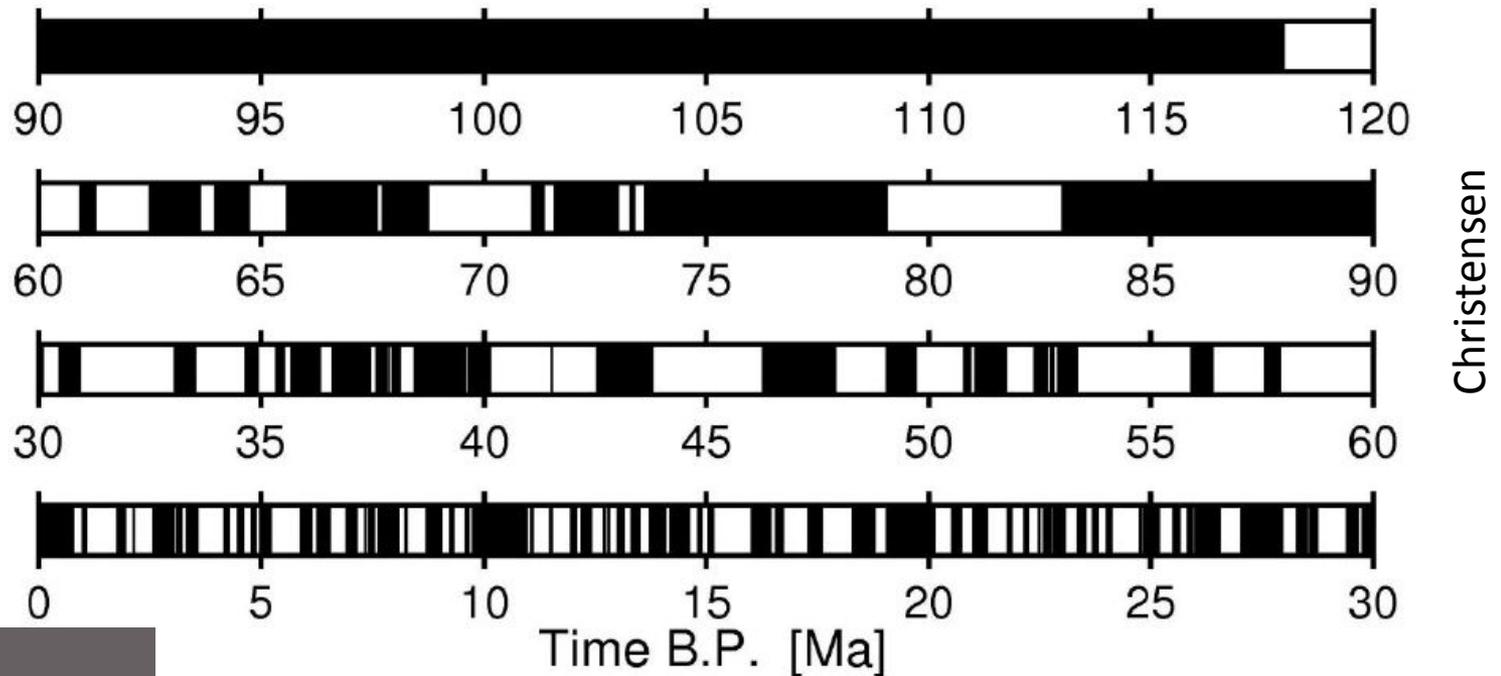


$$\frac{\partial \mathbf{B}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{B} = (\mathbf{B} \cdot \nabla) \mathbf{v} - \mathbf{B}(\nabla \cdot \mathbf{v}) + \eta \nabla^2 \mathbf{B}$$

known

$$\frac{\partial B_r}{\partial t} + \nabla \cdot (\mathbf{v} B_r) = 0$$

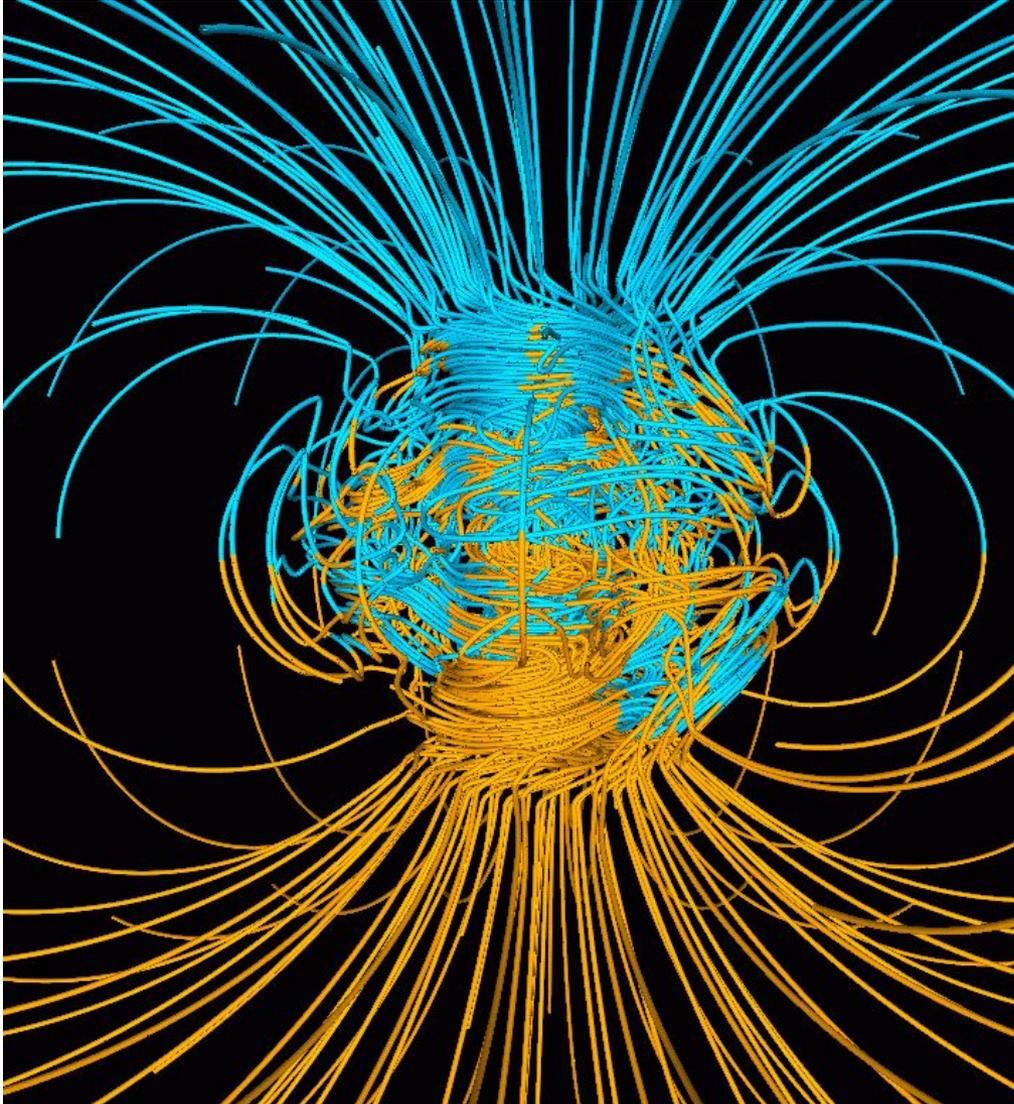
# Longer-term evolution



Like toy dynamo, Earth works in 2 modes. Flips between them seemingly at random

# Model geodynamo

<http://www.es.ucsc.edu/~glatz/geodynamo.html>

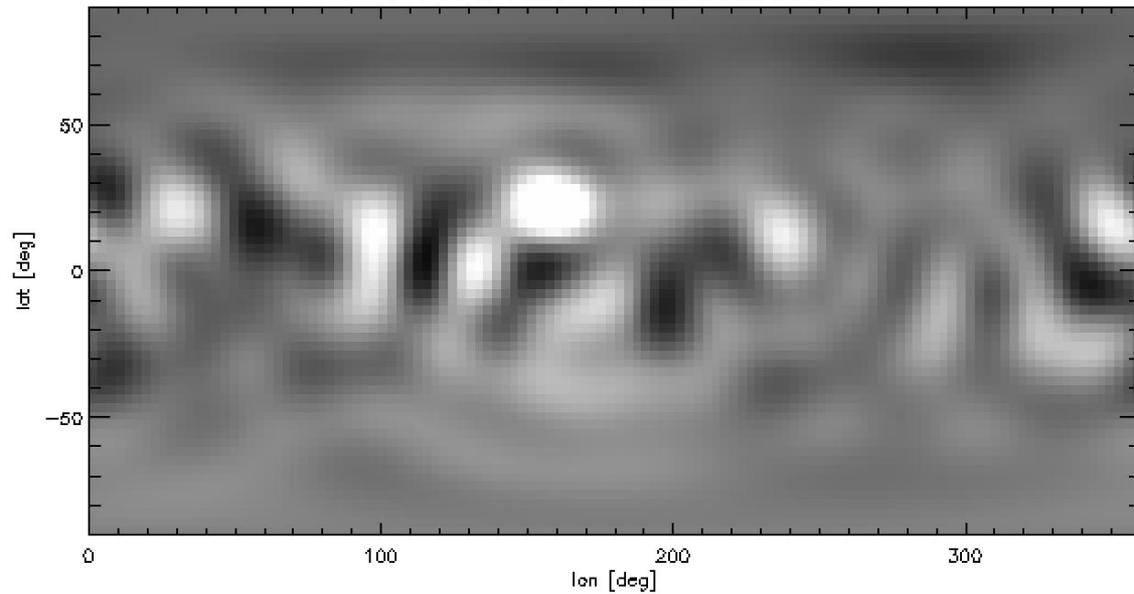


- Glatzmaier & Roberts 1995
- Numerical solution of MHD
- Toroidal structure inside convecting core

Activity 57: Summarize the contrast between the dynamo of a terrestrial planet with that of stars as discussed in this Chapter 4: consider, among others, flow speed, rotation period, stratification, differential rotation, and meridional advection.

	$\eta$ [m <sup>2</sup> /s]	$\nu$ [m <sup>2</sup> /s]	L [m]	v [m/s]	$\Omega$ [rad/s]	Rm	Re	Ro
Sun (CZ)	1	10 <sup>-2</sup>	10 <sup>8</sup>	1	10 <sup>-6</sup>	10 <sup>8</sup>	10 <sup>10</sup>	10 <sup>-2</sup>
Earth (core)	1	10 <sup>-5</sup>	10 <sup>6</sup>	10 <sup>-4</sup>	10 <sup>-4</sup>	10 <sup>2</sup>	10 <sup>7</sup>	10 <sup>-6</sup>

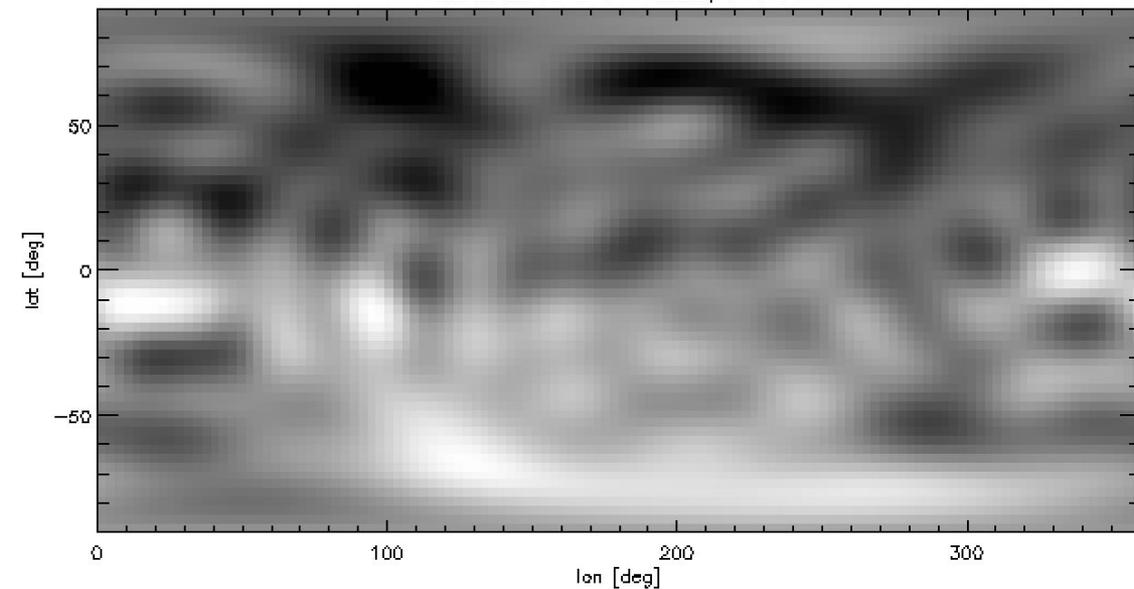
$l < 14$  @ 2001-05-19T20:26:15.000Z



$$Rm = 10^8$$
$$Ro = 10^{-2}$$

Dynamo  
comparison:  
Sun vs. Earth

Earth 2010 @  $r = 0.55$ ;  $l < 14$

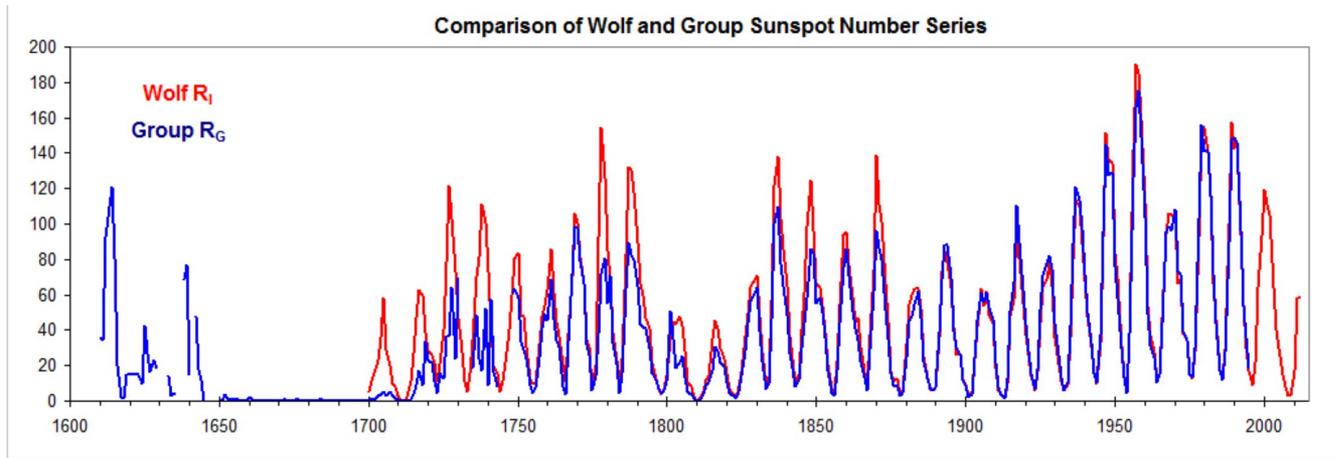


$$Rm = 10^2$$
$$Ro = 10^{-6}$$

$$D \sim Ro^{-2}$$

$$Ro = 10^{-2}$$

$$D \sim 10^4 \gtrsim D_{\text{crit}}$$

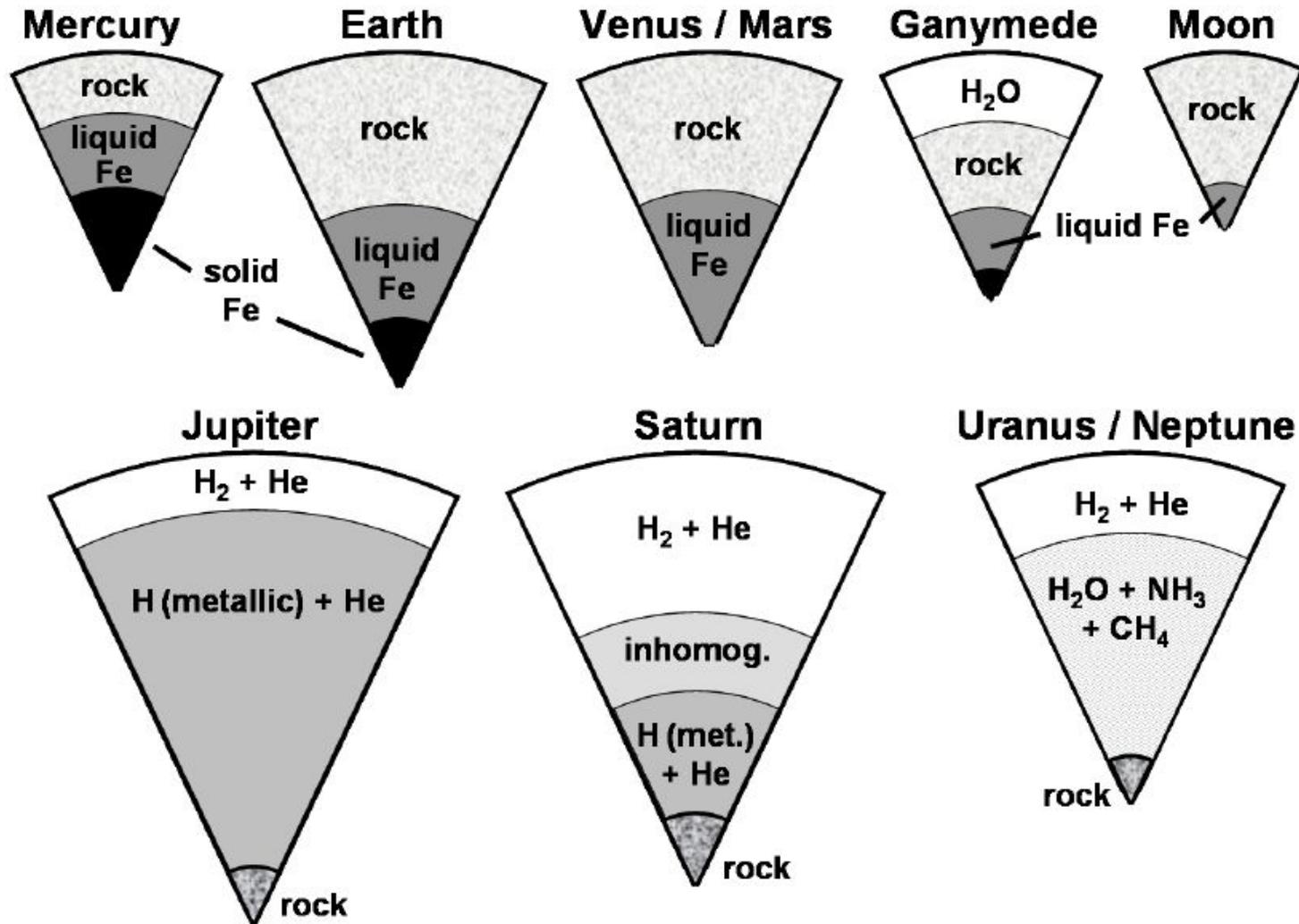


$$Ro = 10^{-6}$$

$$D \sim 10^{12} \gg D_{\text{crit}}$$



# Other planets

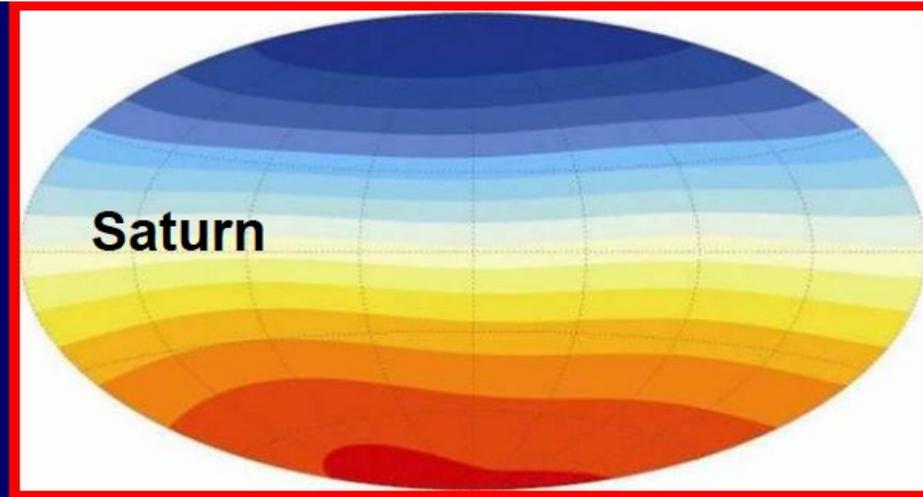
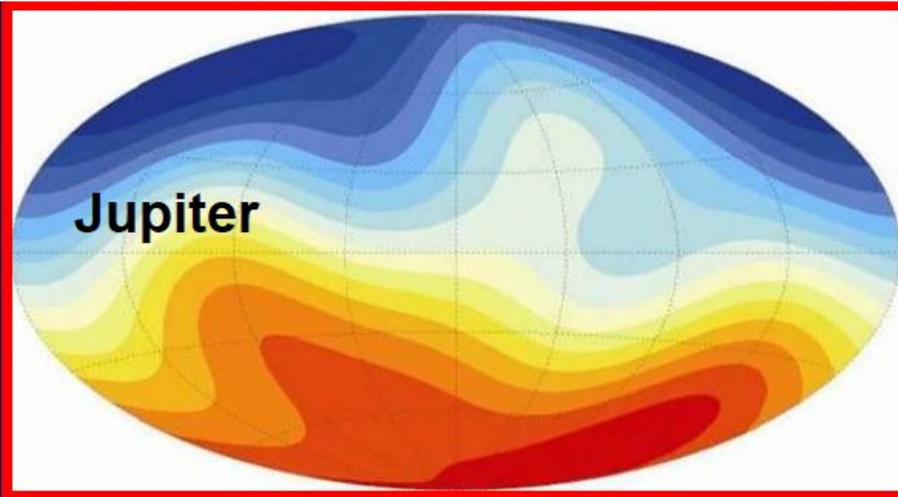


# Other planets

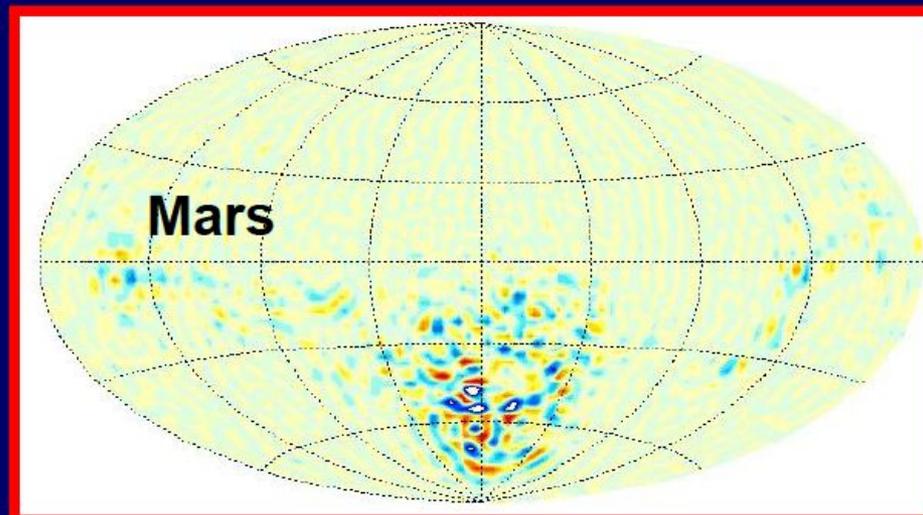
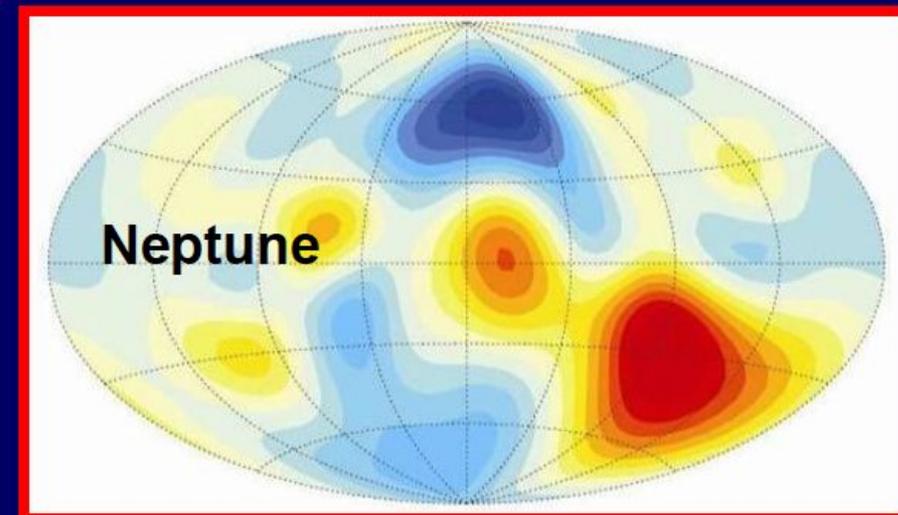
Christensen

Planet	Dynamo	$R_c/R_p$	$B_s$ [ $\mu T$ ]	Dip. tilt	<u>Quadr</u> Dipole
Mercury	Yes (?)	0.75	0.35	$<5^\circ$ ?	0.1-0.5
Venus	No	0.55			
Earth	Yes	0.55	44	$10.4^\circ$	0.04
Moon	No	0.2 ?			
Mars	No, but in past	0.5			
Jupiter	Yes	0.84	640	$9.4^\circ$	0.10
Saturn	Yes	0.6	31	$0^\circ$	0.02
Uranus	Yes	0.75	48	$59^\circ$	1.3
Neptune	Yes	0.75	47	$45^\circ$	2.7
Ganymede	Yes	0.3 ?	1.0	$< 5^\circ$ ?	?

# What that means

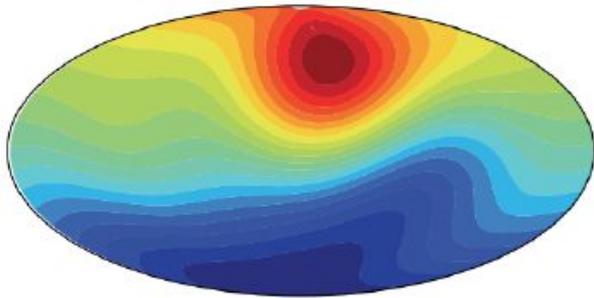


Christensen

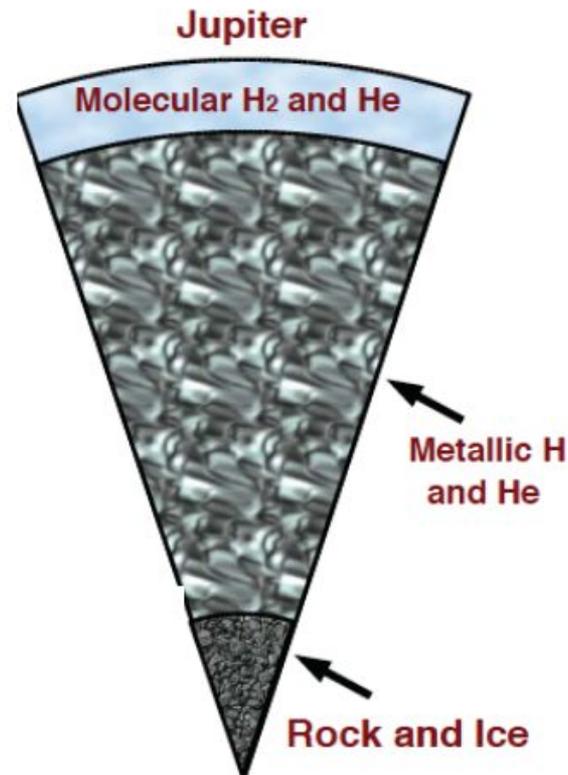
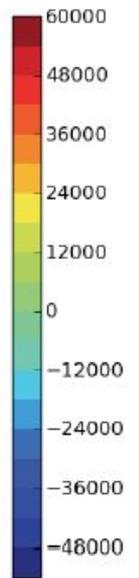
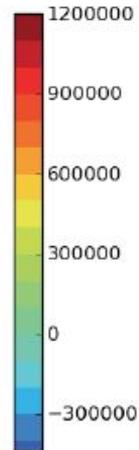
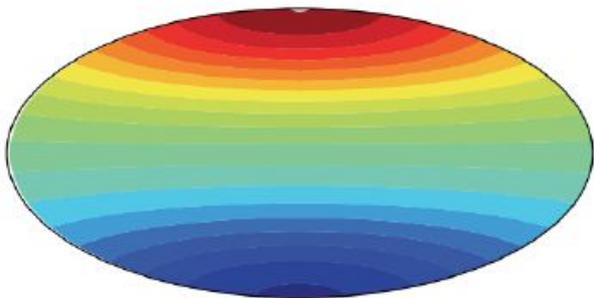


# GAS GIANTS

Jupiter



Saturn

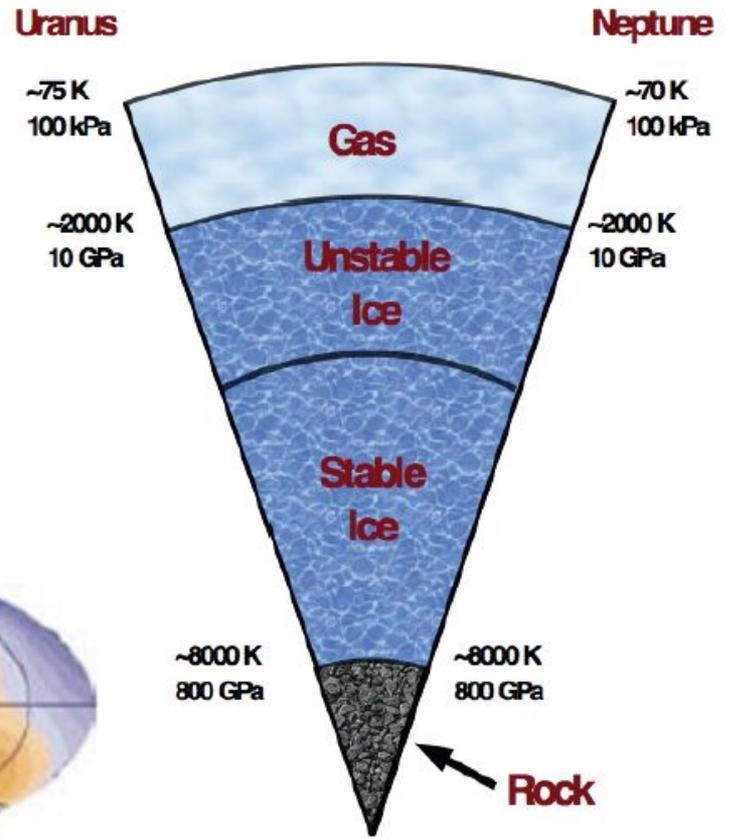
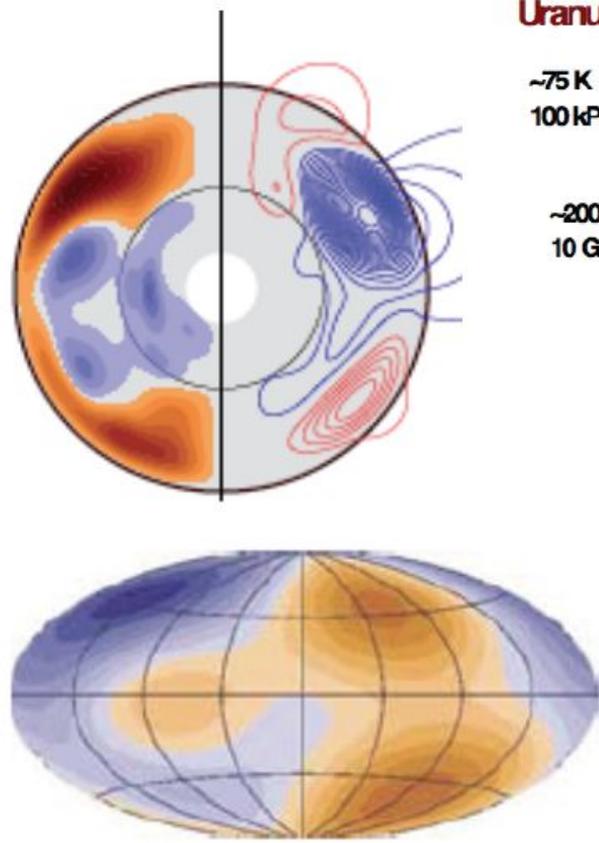
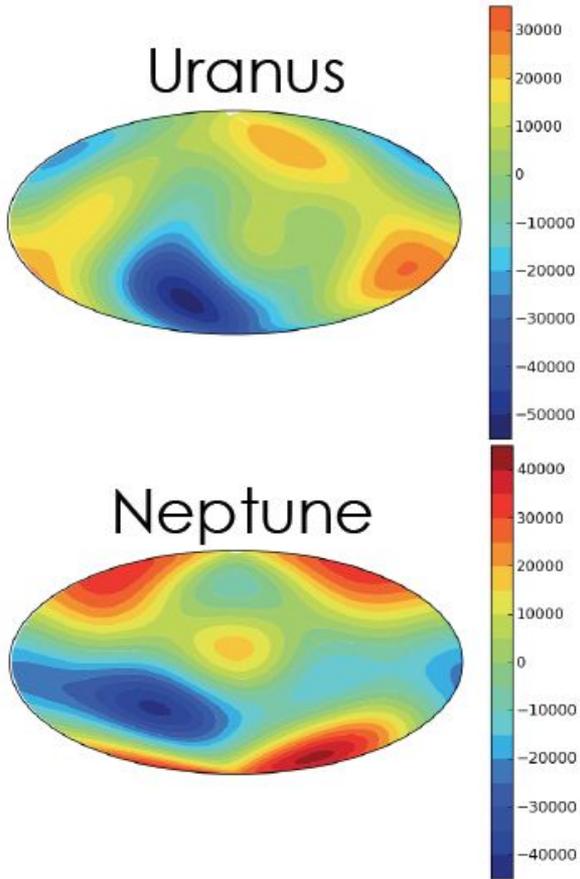


- large pressure range → radially variable properties can be important

# ICE GIANTS

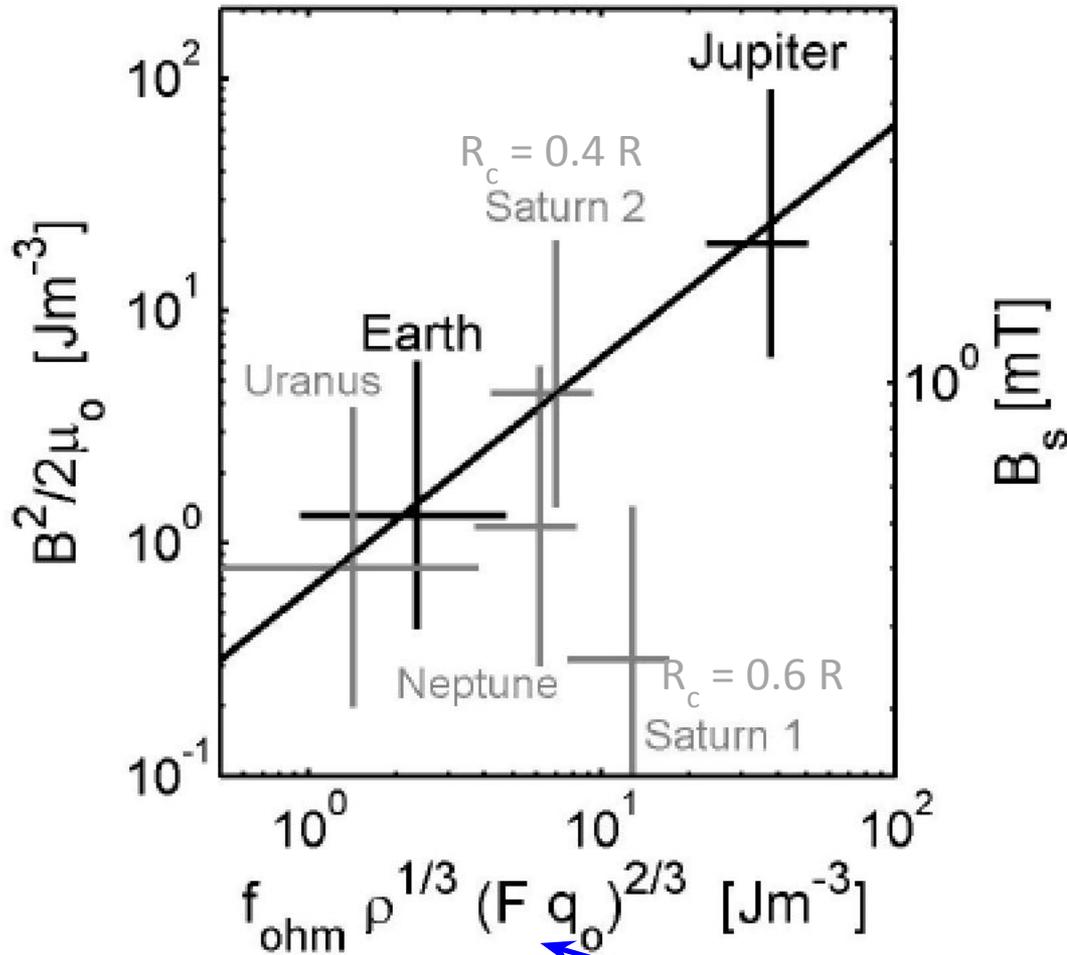
Stanley & Bloxham 2004, 2006

- work in a geometry suggested by low heat flow observations



# Level of saturation

from Christensen



B saturates  
(exp growth ends)  
when driving  
power – thermal  
conduction  $q_0$  –  
balances Ohmic  
dissipation

dimensionless factors

# Summary

- Magnetic fields – all from dynamos
  - Conducting fluid
  - Complex motions w/ enough *umph*
- Create complex fields
- Fields evolve in time – reverse occasionally
- Differences from different parameters:  
Rm, Re, Ro